

# RECENT DEVELOPMENTS IN RAYLEIGH-BÉNARD CONVECTION

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**ABSTRACT:** This review summarizes results for Rayleigh-Bénard convection which have been obtained over the past decade or so. It concentrates on convection in compressed gases and gas mixtures with Prandtl numbers near one and smaller. In addition to the classical problem of a horizontal stationary fluid layer heated from below, it also covers briefly convection in such a layer with rotation about a vertical axis, with inclination, and with modulation of the vertical acceleration.

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INTRODUCTION . . . . . 1

**The whole review (70 pages) can be obtained as a pdf-file upon request ([werner.pesch@uni-bayreuth.de](mailto:werner.pesch@uni-bayreuth.de))**

## 1 INTRODUCTION

Fluid motion driven by thermal gradients (thermal convection) is a common and important phenomenon in nature. Convection is a major feature of the dynamics of the oceans, the atmosphere, and the interior of stars and planets (Busse, 1978, 1989; Getling, 1998). It is also important in numerous industrial processes. For many years, the quest for the understanding

of convective flows has motivated numerous experimental and theoretical studies.

In spatially extended systems convection usually occurs when a sufficiently steep temperature gradient is applied across a fluid layer. The spatial variation of a convection structure often is referred to as a *pattern*. The nature of such convection patterns is at the center of this review. Pattern formation is determined by nonlinear aspects of the system under study. For this reason the elucidation of pattern formation is a challenging problem in condensed-matter physics as well as in fluid mechanics. Pattern formation is common also in many other spatially-extended nonlinear nonequilibrium systems in physics, chemistry, and biology (Manneville, 1990; Cross & Hohenberg, 1993). Patterns observed in diverse systems are often strikingly similar, and their understanding in terms of general, unifying concepts has long been a main direction of research (Cross & Hohenberg, 1993; Newell et al, 1993).

Many fundamental aspects of patterns and their instabilities have been studied intensively over the past three decades in the context of Rayleigh-Bénard convection (RBC). In a traditional RBC experiment a horizontal fluid layer of height  $d$  is confined between two thermally well conducting, parallel plates. When the difference  $\Delta T = T_b - T_t$  between the bottom-plate temperature  $T_b$  and the top-plate temperature  $T_t$  exceeds a critical value  $\Delta T_c$ , the conductive motionless state is unstable and convection sets in. The simplest pattern which can occur is that of straight, parallel convection rolls with a horizontal wavelength  $\lambda \approx 2d$  (wave number  $q \approx \pi/d$ ). Such rolls can be found near onset; however, as the dimensionless distance  $\epsilon \equiv \Delta T/\Delta T_c - 1$  increases, the patterns often become progressively more complicated, and thus also more interesting.

Rayleigh-Bénard convection is perhaps the most thoroughly investigated and understood pattern-forming system. The experimental setup is simple in principle and the basic physical mechanism (buoyancy vs. dissipation) well understood. For the standard description in terms of the Oberbeck-Boussinesq

equations, Eqs. ?? and ?? below, only two nondimensionalized control parameters are sufficient (Busse, 1978, 1989). The first is the Rayleigh number  $R \equiv \alpha g \Delta T d^3 / (\nu \kappa)$  with  $\alpha$  the thermal expansion coefficient,  $\kappa$  the thermal diffusivity,  $\nu$  the kinematic viscosity, and  $g$  the acceleration of gravity. Convection starts (under ideal conditions) at the critical value  $R_c = 1708$ . The second parameter is the Prandtl number  $\sigma \equiv \nu / \kappa$ , which can be viewed as the ratio of the vertical thermal diffusion time  $t_v = d^2 / \kappa$  to the vertical viscous relaxation time  $t_\nu = d^2 / \nu$ . It measures the relative importance of the nonlinear terms in the Boussinesq equations, namely those terms describing temperature and momentum advection.

There are several recent reviews of Rayleigh-Bénard convection<sup>1</sup>. In the present one we will focus on new developments during the last decade or so. Let us, however, briefly outline some of the seminal earlier results which were of major importance to the later work. Quite early it was established theoretically that the stable pattern in an infinitely extended layer of a Boussinesq fluid close to onset (see Sect. ?? below) consists of straight, parallel rolls of wavenumber  $q$  (Schlüter et al, 1965). Further above onset, the stability regimes of these rolls in the  $R - q$  space as functions of  $\sigma$  (the “Busse balloon”) are well understood owing to the impressive work by Busse and coworkers (Busse, 1978, 1989). The value of  $\sigma$  varies widely for different experimental fluids, from  $\mathcal{O}(10^{-2})$  for liquid metals to values near one for gases and for liquid helium, to the range from 2 to 12 for water, and into the 1000’s for silicone oil (see de Bruyn et al, 1996, and references therein). Although  $R_c$ , the critical wavevector  $q_c = 3.117$ , and the patterns in the close vicinity of onset ( $R \simeq R_c$ ) are independent of  $\sigma$ , this does not apply to the subsequent bifurcations which occur with increasing  $R$ . For example one finds an oscillatory secondary instability at medium and small Prandtl numbers, in contrast to a stationary bimodal (“knot”) bifurcation at large  $\sigma$  (Busse, 1978, 1989). Not too far

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<sup>1</sup>Busse (1989), Croquette (1989a,b), Cross & Hohenberg (1993), de Bruyn et al (1996), and Getling (1991, 1998).

from threshold the Busse balloon was found to agree well with the experiments for large  $\sigma$  (water) (Busse & Whitehead, 1971) and reasonably well for gases with  $\sigma \simeq 1$  (Croquette, 1989a).

Although ideal periodic patterns can be created in experiments, “natural” convective patterns, particularly when they form in the presence of lateral walls, typically are disordered and develop persistent spatio-temporal dynamics as  $\epsilon$  increases. Snapshots of such patterns are characterized by local roll patches with grain boundaries and point defects (dislocations)<sup>2</sup>. These spatio-temporal chaotic patterns are irregular in time and in the horizontal plane; but they maintain a relatively simple structure in the vertical direction. They should be contrasted with fully developed turbulence (Frisch, 1995), at very large  $R$ , which is disordered in three spatial dimensions and not the topic of this review.

For the description of nonuniform patterns not too far from threshold various reductions of the original hydrodynamical equations have proven to be useful (Manneville, 1990; Cross & Hohenberg, 1993; Newell et al, 1993). One particularly important theoretical result, which motivated much of the work during the last decade, was that for a fluid of  $\sigma \simeq 1$  roll curvature induces slowly-varying long-range pressure gradients (Siggia & Zippelius, 1981) that drive a so called “mean flow” which in turn couples back to the roll curvature. In subsequent experimental work<sup>3</sup> these ideas have found their convincing confirmation. Model equations which generalize the so called Swift-Hohenberg equations (SH-equations) (Swift & Hohenberg, 1977) were developed. They allow the isotropic description of the pattern-formation processes in the presence of mean flow (Manneville, 1983; Greenside & Coughran, 1984). Although the SH-equations can not be derived systematically from the Boussinesq equations, they capture much of the observed physical behavior and have now become a general tool

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<sup>2</sup>see Croquette (1989b), Manneville (1990), Busse et al (1992), Newell et al (1993), Cross & Hohenberg (1994), Xi et al (1997), and Busse & Clever (1998).

<sup>3</sup>See, for instance, Croquette et al (1986b); Croquette (1989a); Daviaud & Pocheau (1989); Hu et al (1995a); and Pocheau & Daviaud (1997)

to investigate not only RBC, but also other pattern-forming systems (Cross & Hohenberg, 1993).

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