

Lithospheric processes: From oceanic lithosphere accretion to hydrothermal cooling

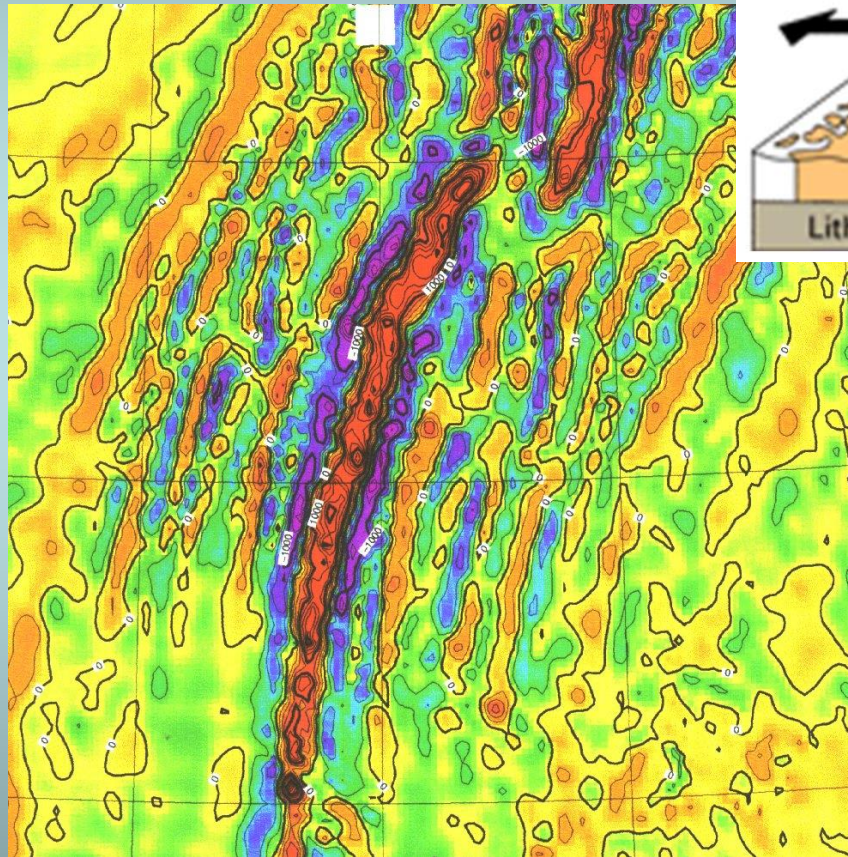
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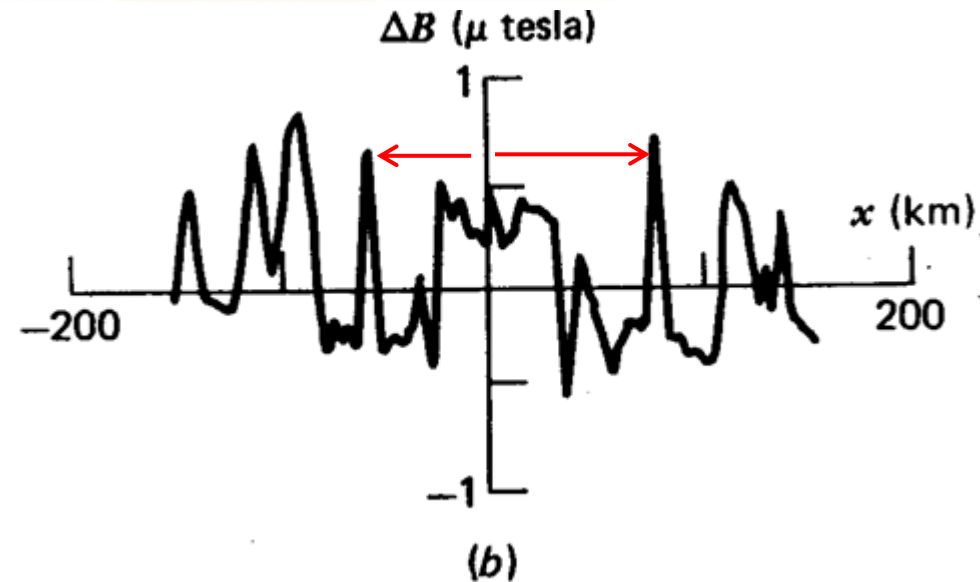
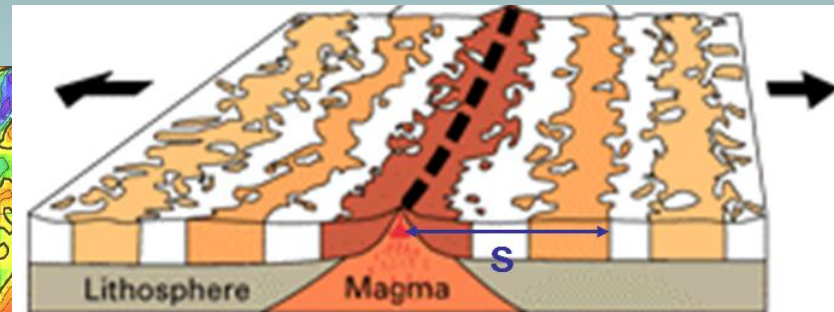
Outline

- Why do MORs spread symmetrically and how is this related to ridge migration?
 - Relation MOR symmetry – ridge migration
 - Modelling migrating ridges → symmetric spreading
 - Symmetric even when overriding plumes?
 - Comparison to lava lakes (thermal boundary condition)
- Square root cooling away from the MOR
 - Deviations due to hydrothermal convection
 - Fitting to observations
- Modelling crust generation at MOR
 - Focussing melt towards the ridge (Katz 2008)
 - Asymmetric mantle potential temperature (Katz 2010)

Examples symmetric and asymmetric spreading

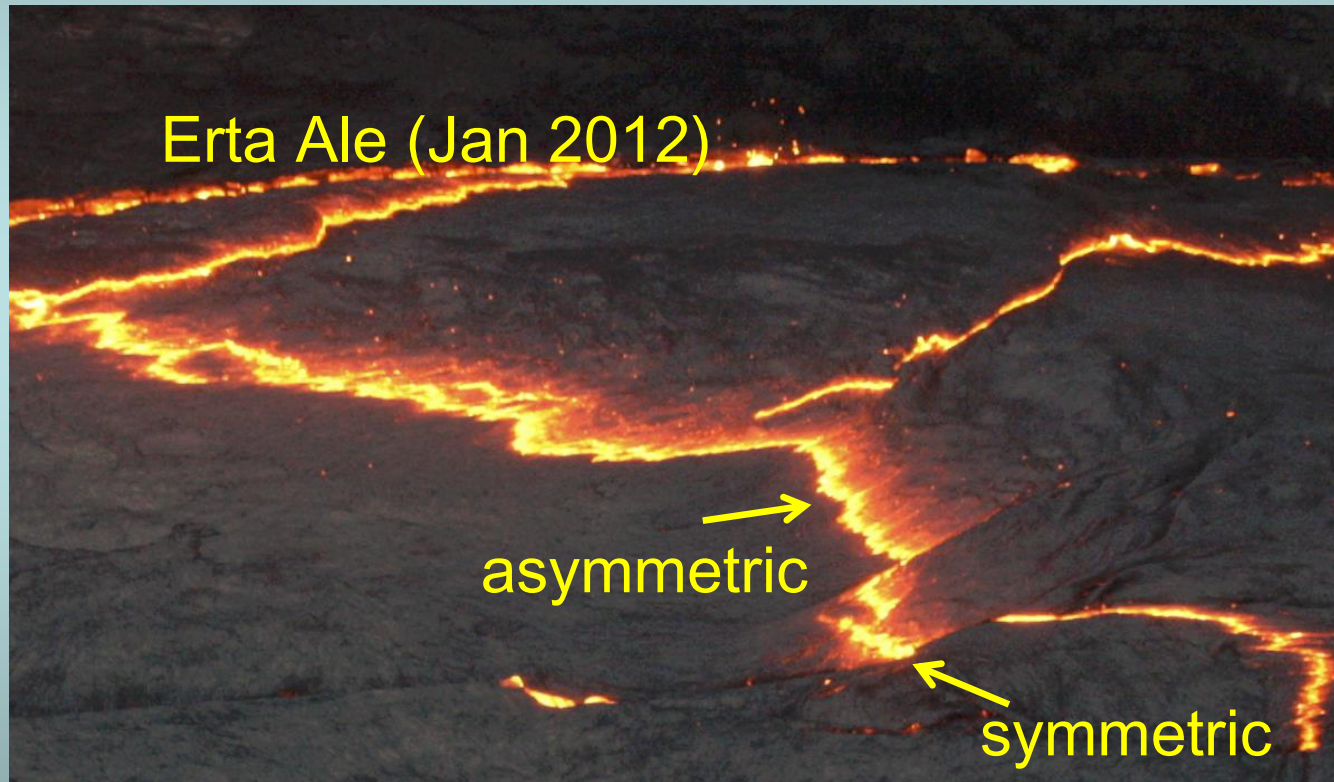


Magnetic anomalies,
Kolbeinsey ridge (symmetric)



South East Indian Rise
(asymmetric)

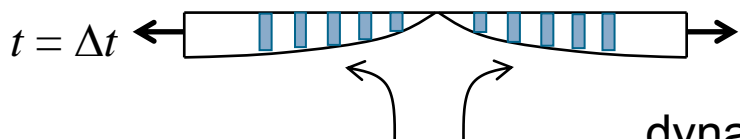
Examples symmetric and asymmetric spreading



Some kinematics

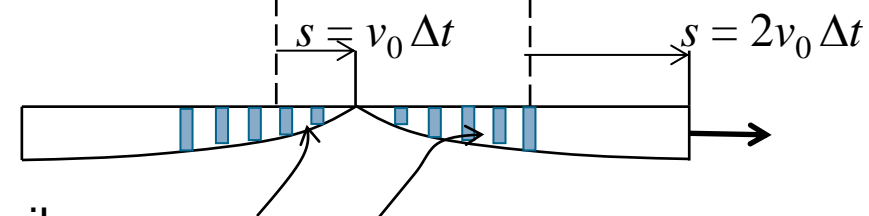
Symmetric spreading (Assume: passive mantle)

Coordinate system fixed at ridge



No ridge migration, $v_{ridge} = 0$

Coordinate system fixed at one plate

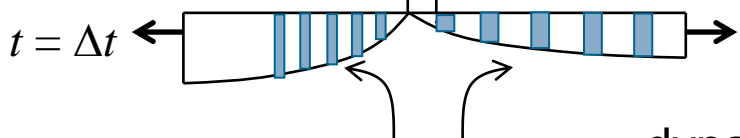


$v_{ridge} = v_0$

=
dynamically similar

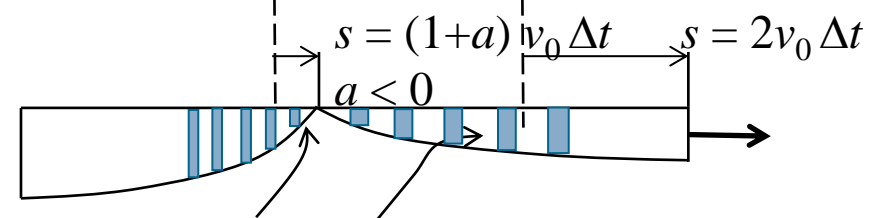
Asymmetric spreading (Assume: passive mantle)

Coordinate system fixed at $(v_{left} + v_{right})/2$



Ridge migration, $v_{ridge} = a v_0$

Coordinate system fixed at v_{left}



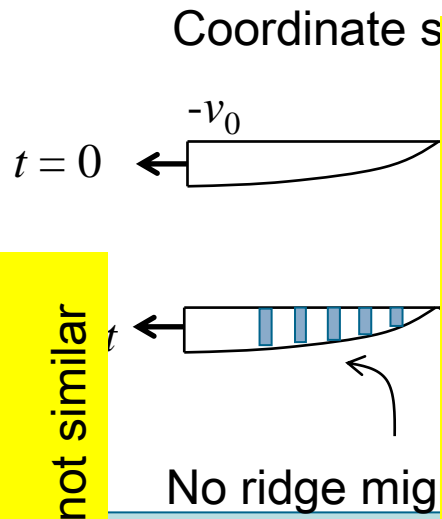
$v_{ridge} = (1+a) v_0$

=
dynamically similar

Some kinematics

Symmetric spreading (Assume: passive mantle)

Coordinate system fixed at one plate



Asymmetry parameters

Asymmetric boundary condition parameter

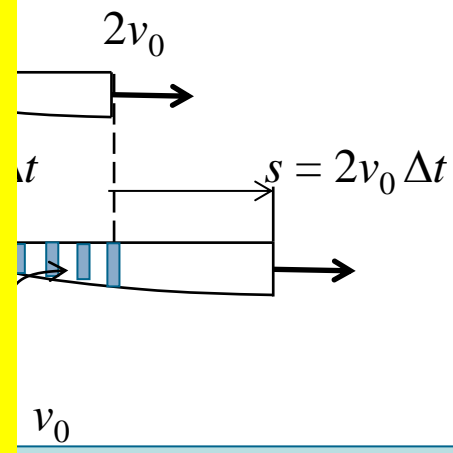
$$c = (v_{left} + v_{right}) / (2v_0)$$

0 → symmetric boundary conditions
up to ±1 → asymmetric b.c.'s

Asymmetric spreading parameter

$$a = (v_{ridge} - v_{plates-mean}) / v_0$$

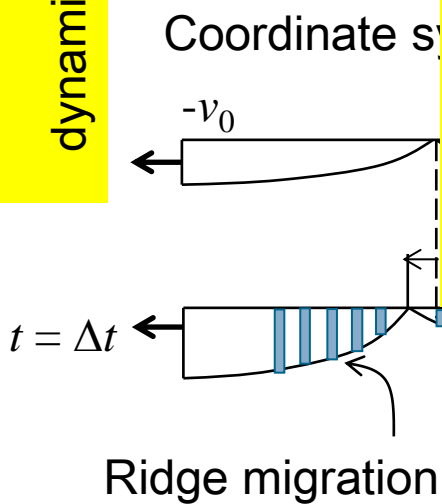
0 → symmetric
up to ±1 → asymmetric



dynamically not similar

(Assume: passive mantle)

Coordinate system fixed at v_{left}

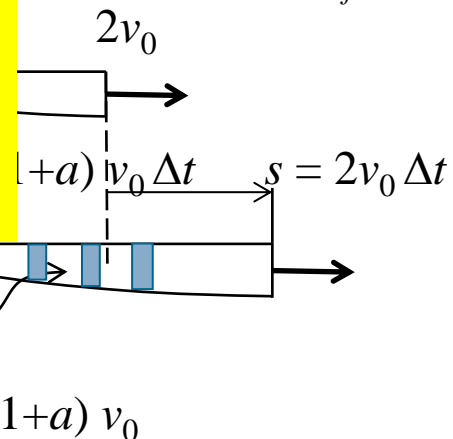


Asymmetric spreading parameter

$$a = (v_{ridge} - v_{plates-mean}) / v_0$$

0 → symmetric
up to ±1 → asymmetric

dynamically similar



=

Governing equations

$$-\vec{\nabla} P + \frac{\partial}{\partial x_j} \left[\eta \left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right) \right] - \rho g \vec{e}_3 = 0$$

Navier Stokes

$$\dot{e}_{ij} = \frac{1}{2} \left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right),$$

$$\tau_{ij} = 2\eta \dot{e}_{ij}$$

$$\vec{\nabla} \cdot \vec{v} = 0$$

Mass conservation

$$\rho c_p \left(\frac{\partial T}{\partial t} + \vec{v} \cdot \vec{\nabla} T \right) = \vec{\nabla} \cdot (k \vec{\nabla} T) + \rho H$$

Energy conservation

"Pore pressure" factor

Mohr Coulomb (Byerlee) plasticity:

$$\dot{e}_{ij}^{MC} = \begin{cases} 0 & \tau_{II} < \tau_{\max} \\ \text{arbitrary} & \tau_{II} = \tau_{\max} \end{cases} \quad \dot{e}_{ij}^{MC} = \frac{1}{2\eta^{MC}} \tau_{ij} \quad \text{with} \quad \eta^{MC} = \frac{\tau_{\max}}{2\dot{e}_{II}} \quad \tau_{\max} = (a_{By} z + b_{By}) \lambda_p$$

Dislocation creep

$$\dot{e}_{ij}^{duc} = B \exp\left(\frac{-E}{RT}\right) \tau_{II}^{n-1} \tau_{ij}$$

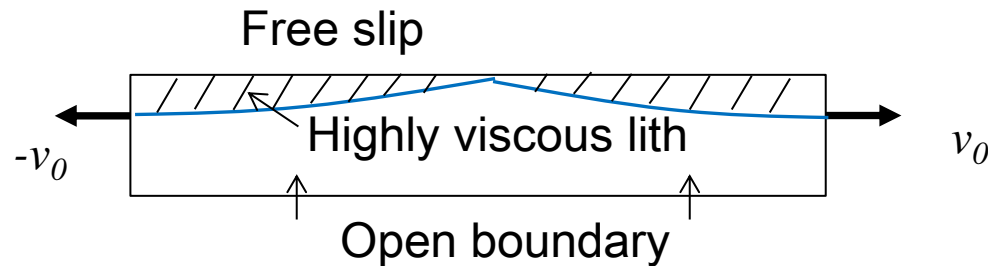
$$\eta^{duc} = \frac{1}{2B^{1/n}} \exp\left(\frac{E}{nRT}\right) \dot{e}_{II}^{duc \left(\frac{1}{n}-1\right)}$$

Composite viscosity

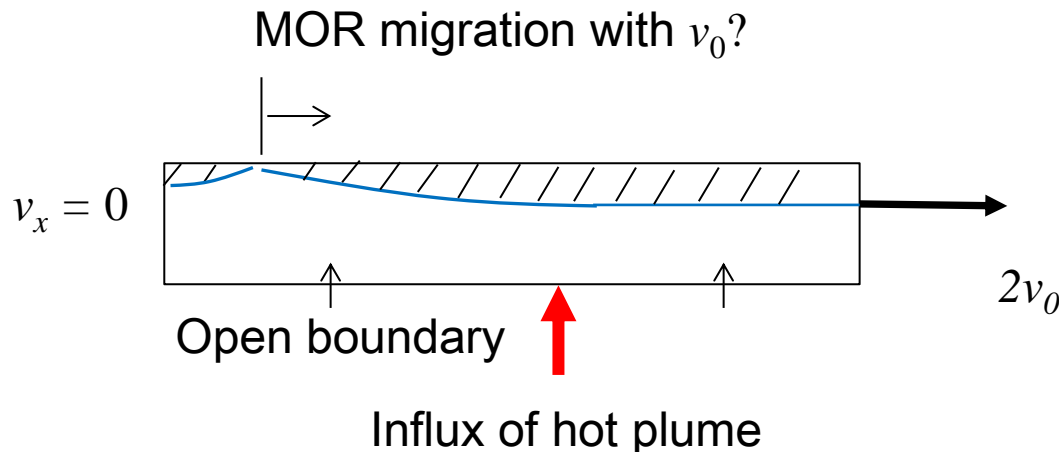
$$\dot{e}_{ij} = \dot{e}_{ij}^{MC} + \dot{e}_{ij}^{duc} = \left(\frac{1}{2\eta^{MC}} + \frac{1}{2\eta^{duc}} \right) \tau_{ij}$$

$$\frac{1}{\eta_{eff}} = \frac{1}{\eta^{MC}} + \frac{1}{\eta^{duc}}$$

Model setup



or:

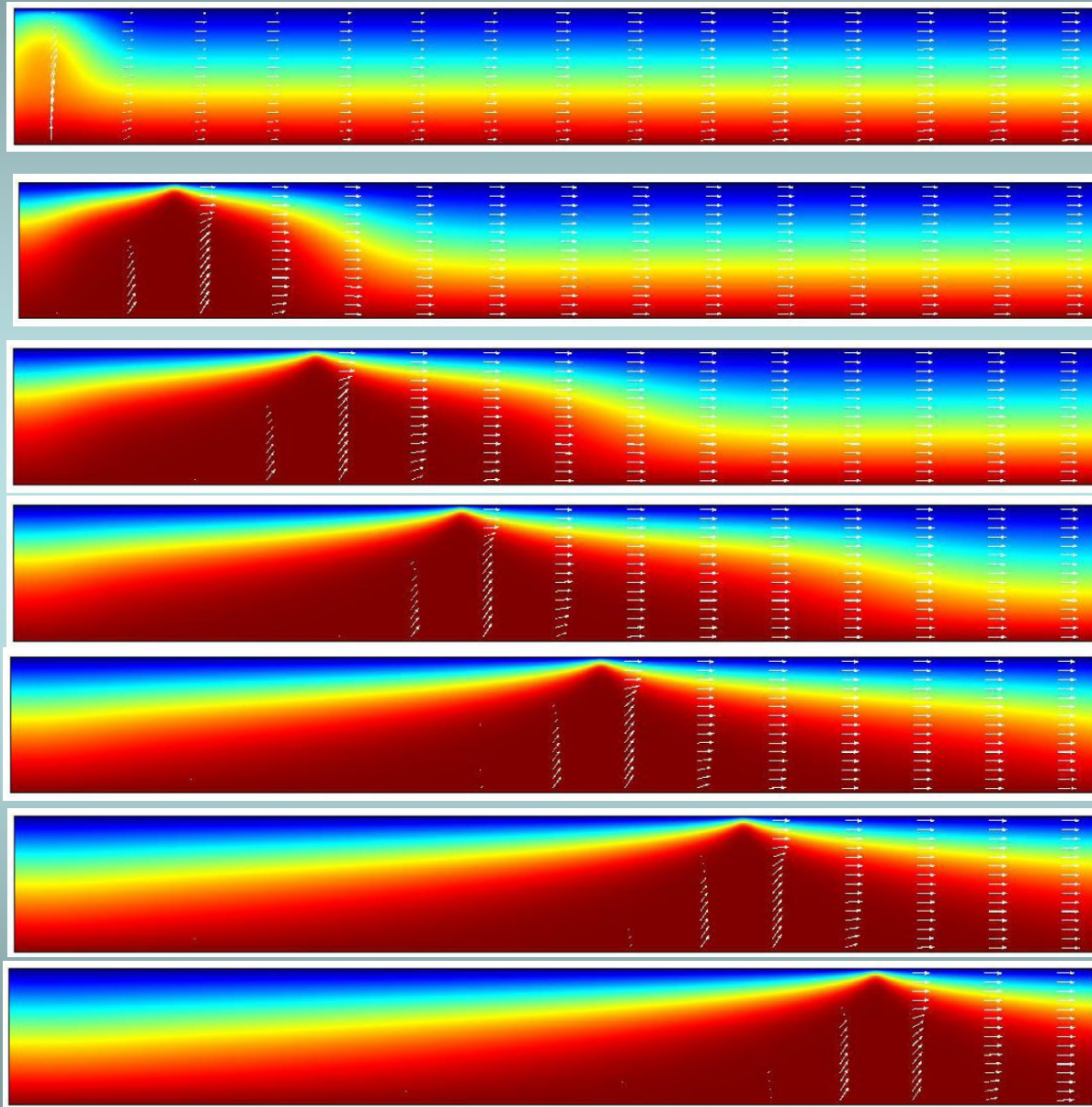


Variation:

- Spreading velocity
- "Pore pressure factor λ_p
- Plume temperature

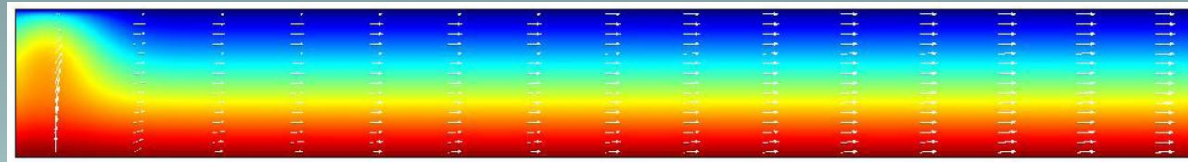
Asymmetry spreading
parameter a ?

Fixed
Free slip

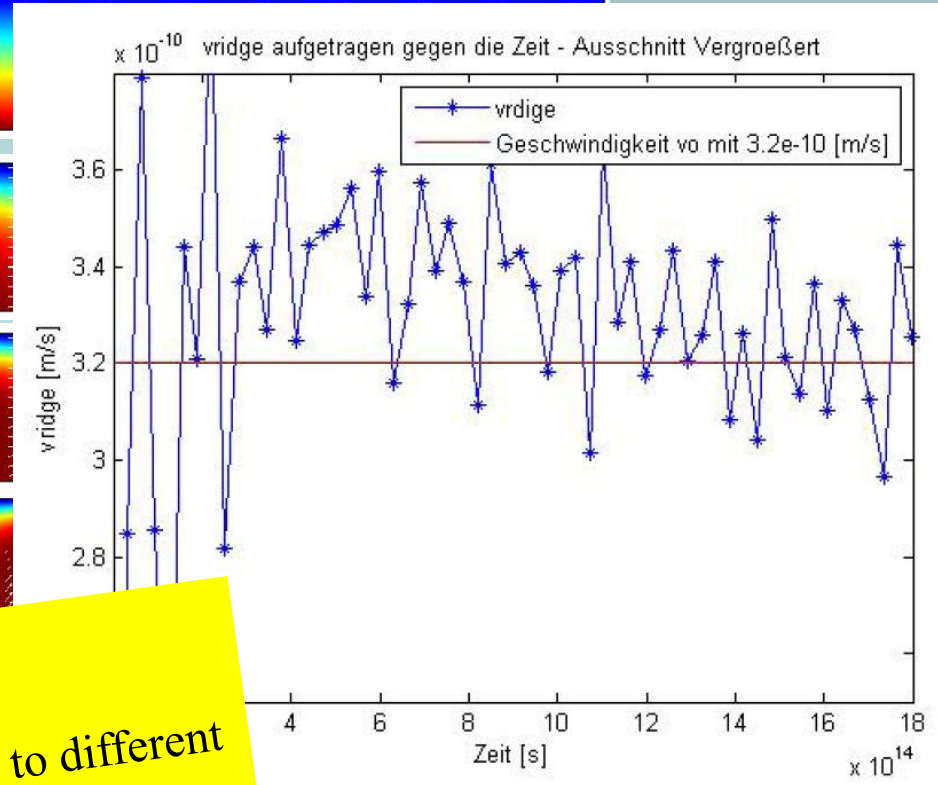
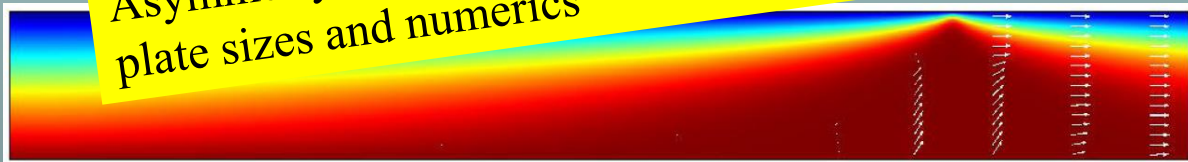
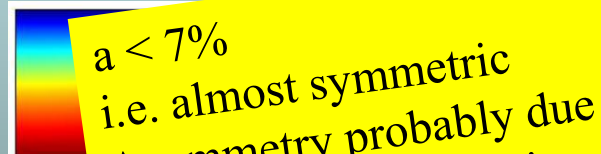
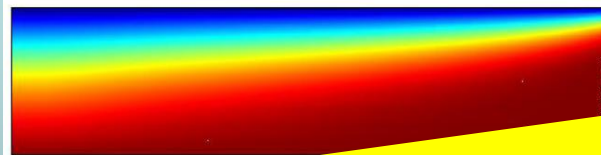
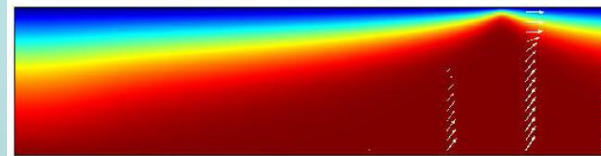
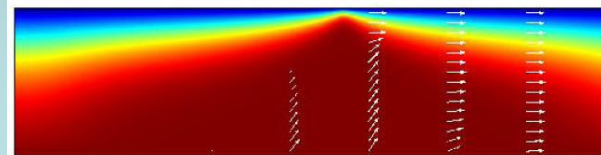
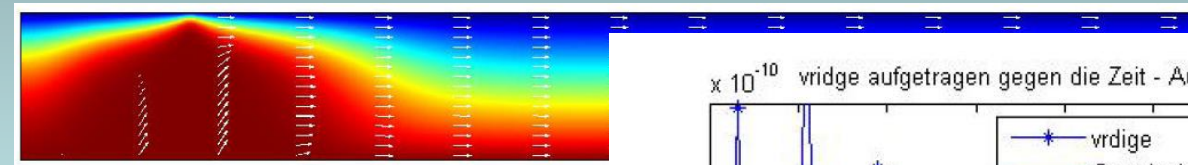


→ $2v_0$

Fixed
Free slip



→ $2v_0$



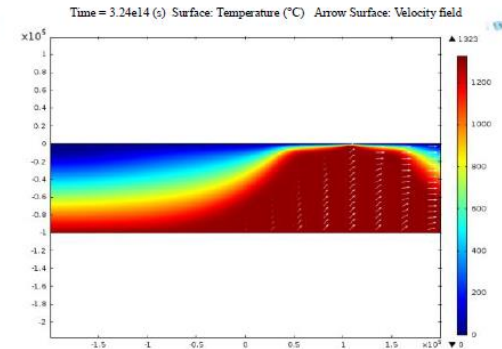
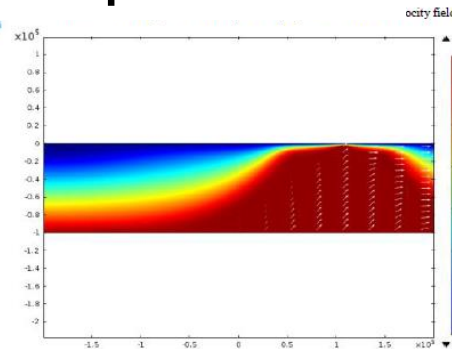
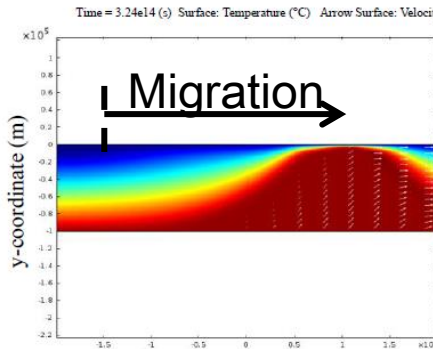
$a < 7\%$
i.e. almost symmetric
Asymmetry probably due to different
plate sizes and numerics

Variation of spreading velocity and "pore pressure" parameter

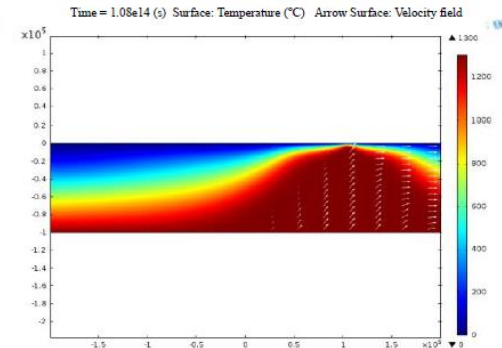
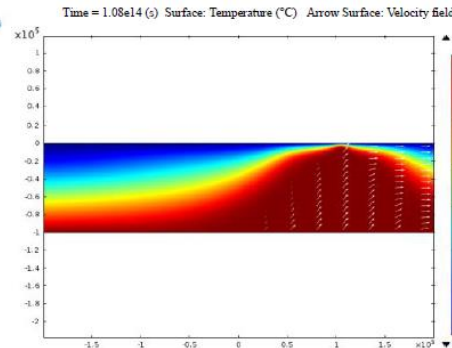
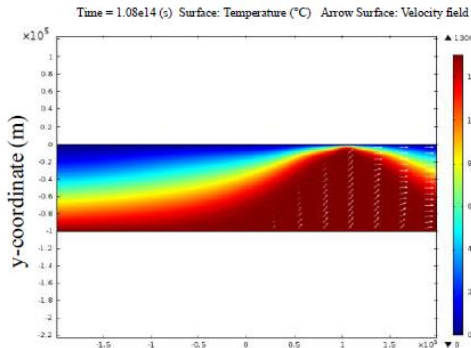
Temperature

v_0
[mm/a]

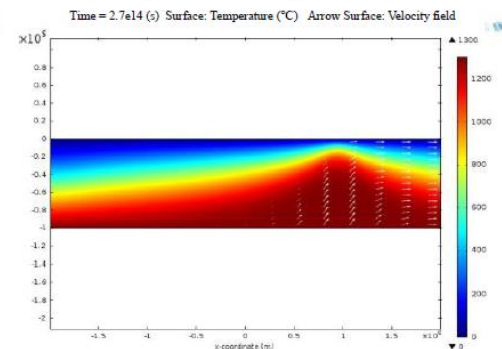
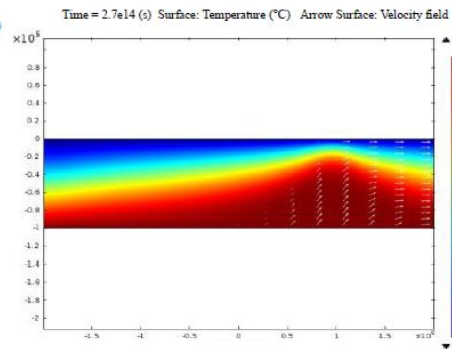
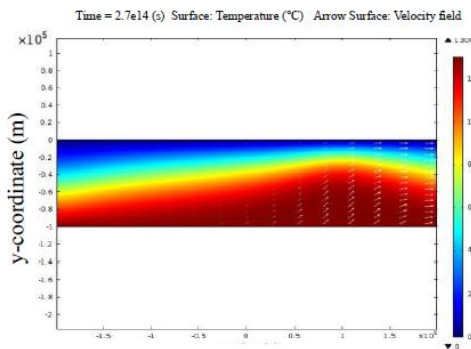
100



30



10



x-coordinate (m)

x-coordinate (m)

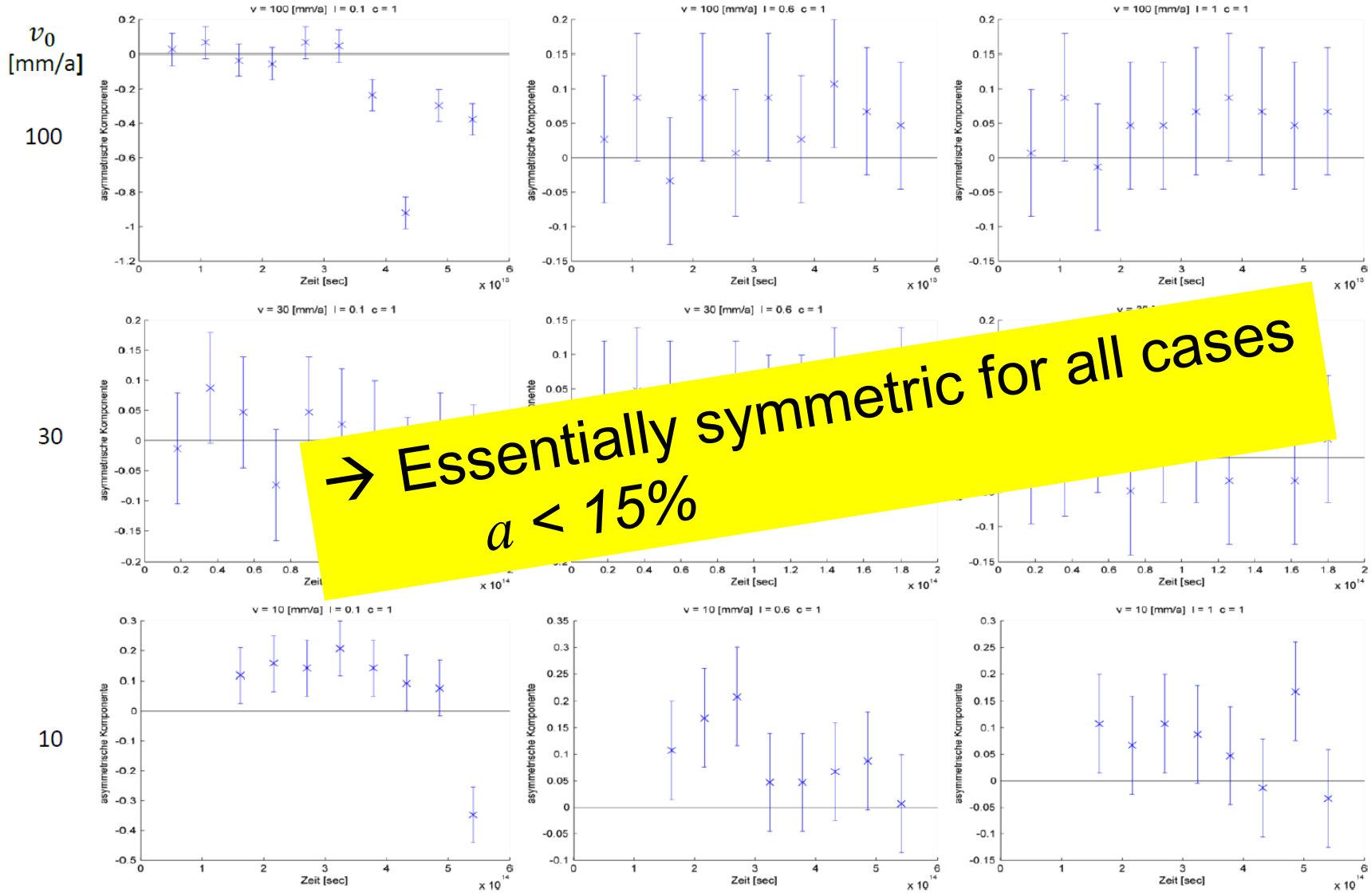
x-coordinate (m)

λ_p 0,1

0,6

1

Asymmetry spreading parameter a (=0 if symmetric)



**→ Essentially symmetric for all cases
 $a < 15\%$**

λ_p

0,1

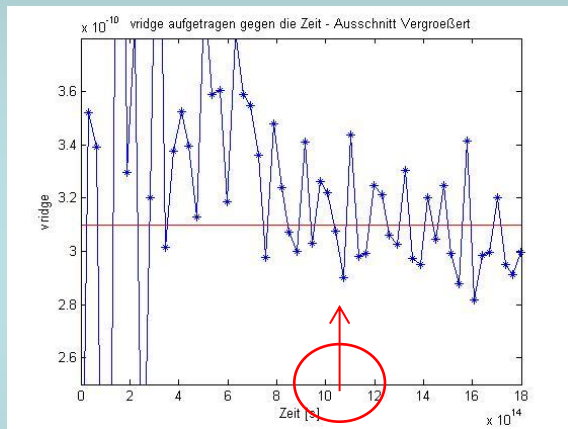
0,6

1

MOR overriding a plume

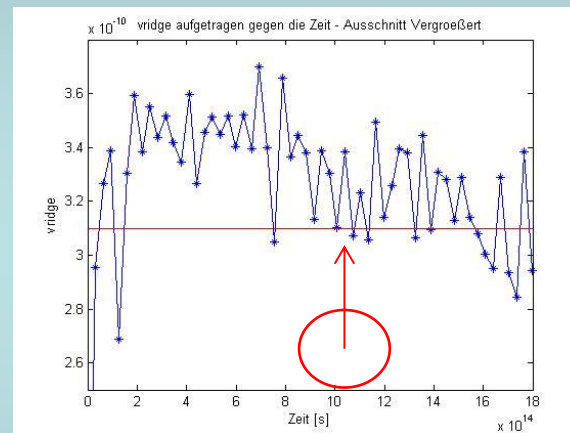
MOR migration velocity for different plume excess temperatures

150 K



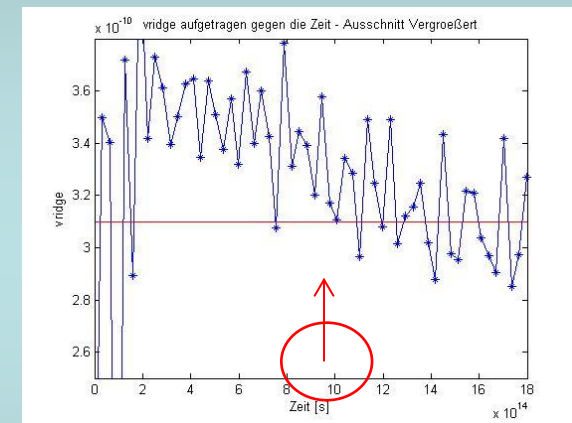
260420 dof's

250 K



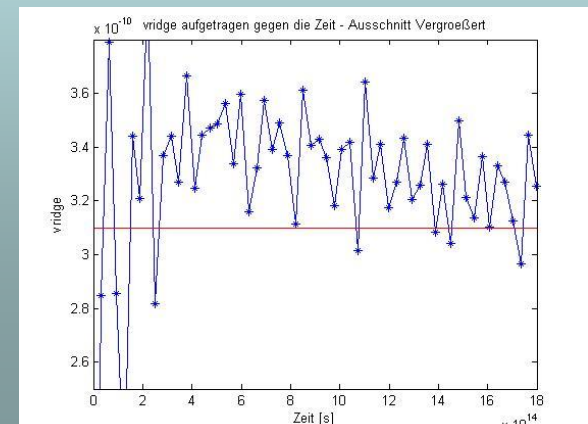
260420 dof's

350 K



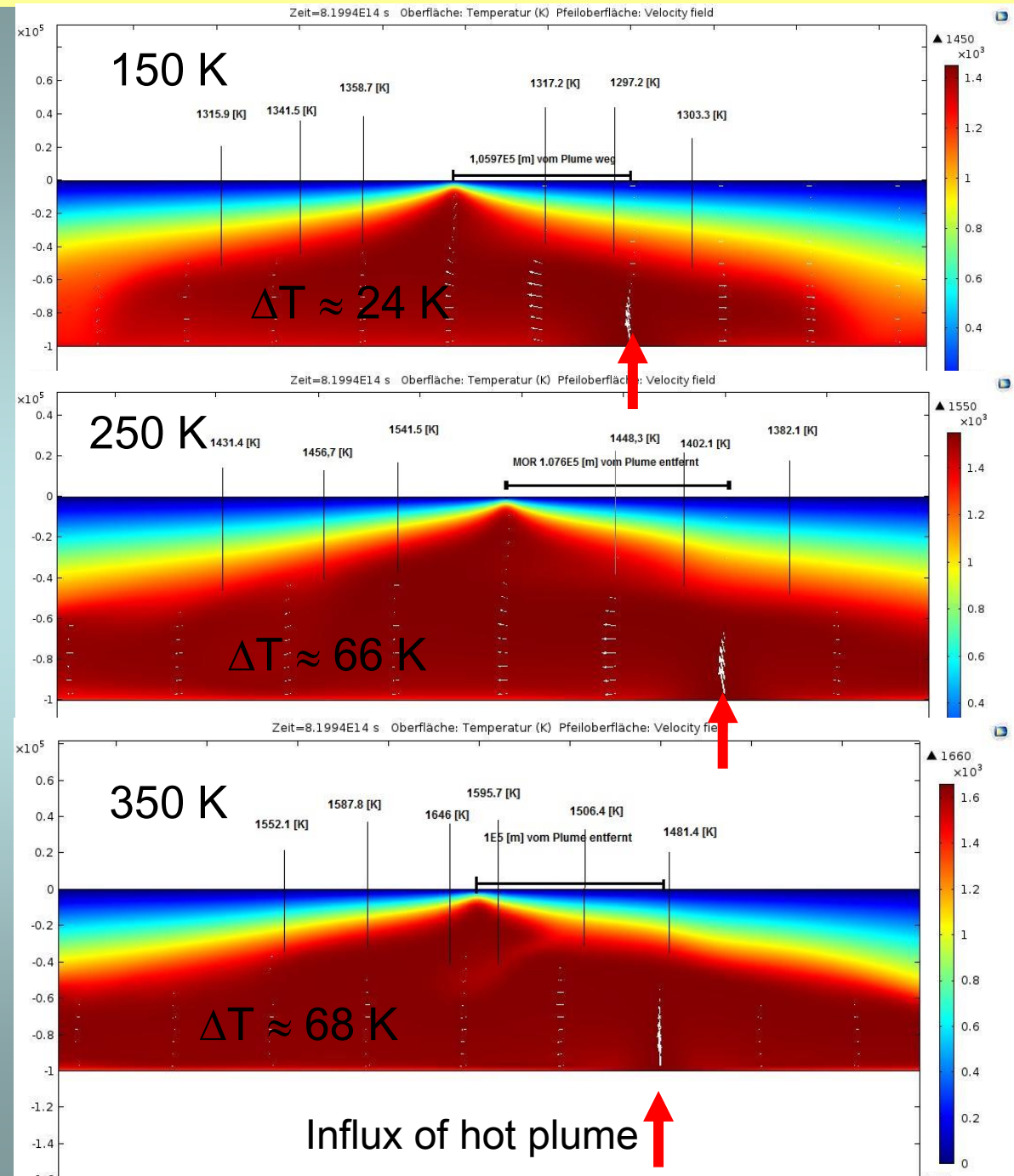
34420 dof's

- MOR slows down when approaching and overriding the plume
- For comparison, without plume →

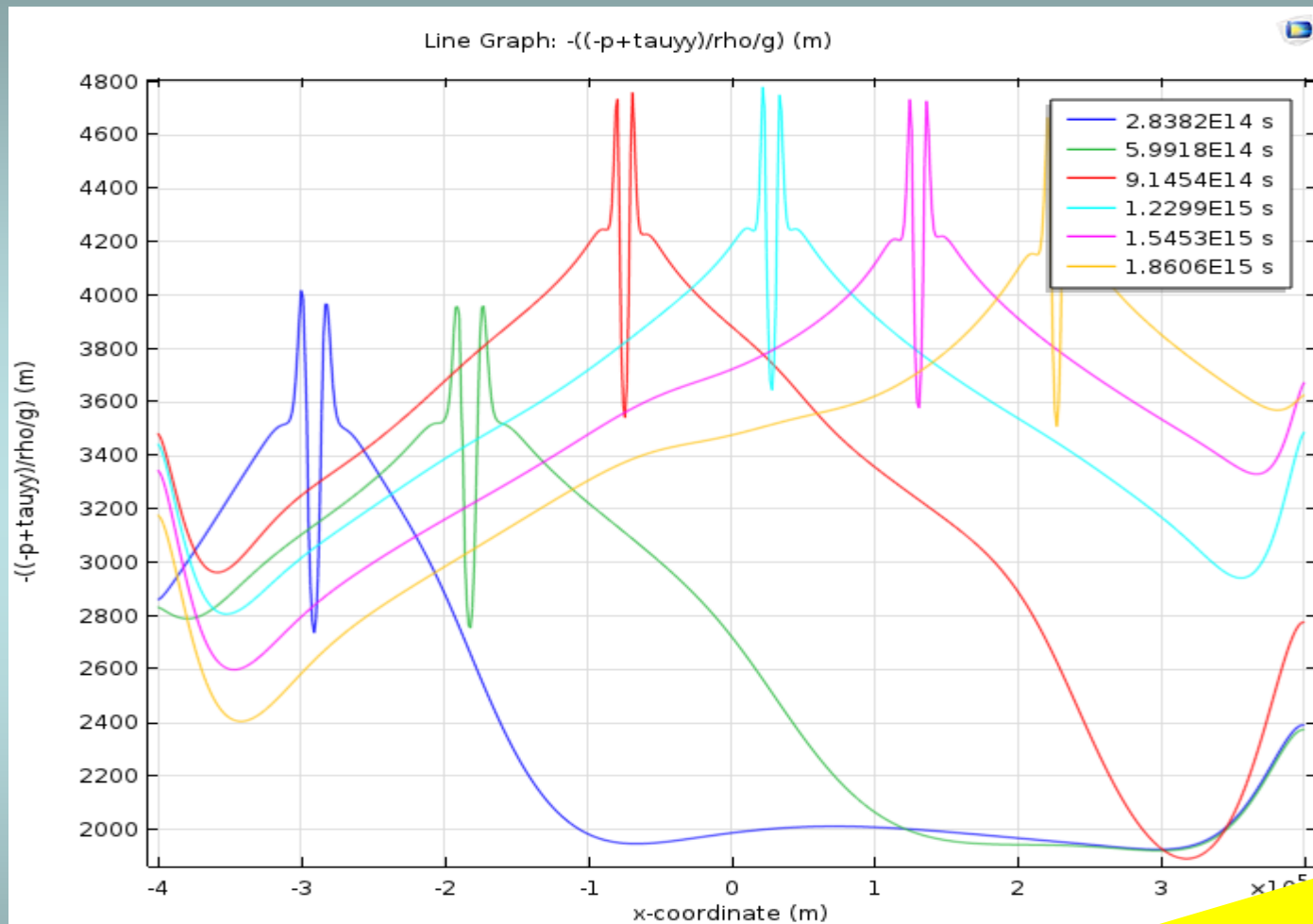


Snapshots at times before MOR will override plume

- LAB cooler (!) on plume side
- Weak, cooler part of thermal boundary layer is swept towards MOR
- Cooler accretion more effective than hot accretion (?) → MOR migration slightly slows down



Plume effect on topography



- Topography remains essentially symmetric
- Hot plume material lifts up the whole ridge

Conclusion:
Both observed and modelled
asymmetries < 15%

What is the difference to lava lake plate tectonics?

Rheology

$$\eta^{duc} = \eta_0 \exp\left(\frac{A_v}{T - T_v}\right) \quad \text{Vogel-Fulcher-Tammann equation, empirical}$$

$$\dot{\epsilon}_{ij}^{MC} = \frac{1}{2\eta^{MC}} \tau_{ij} \quad \text{with} \quad \eta^{MC} = \frac{\tau_{\max}}{2\dot{\epsilon}_{II}}$$

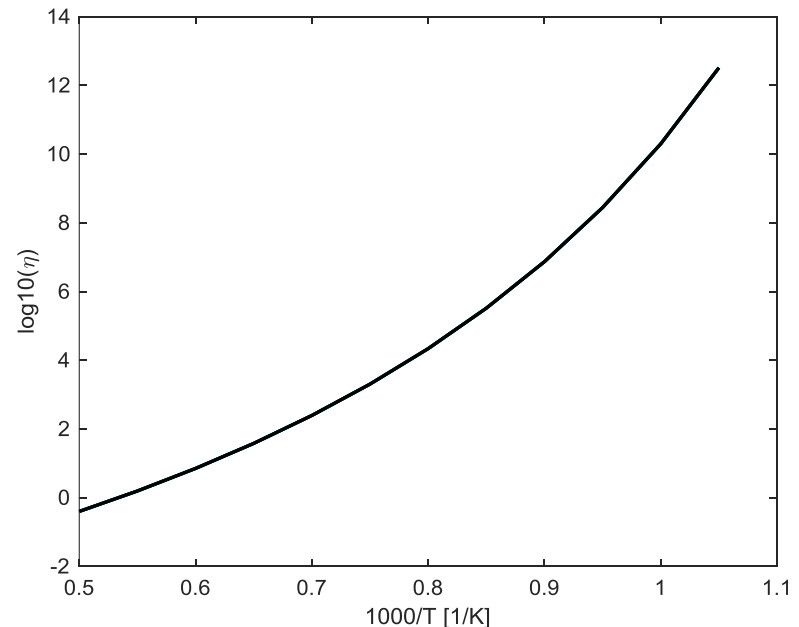
- Stronger T-dependence (more than Arrhenius)
- Linear

Thermal boundary condition:
Radiation

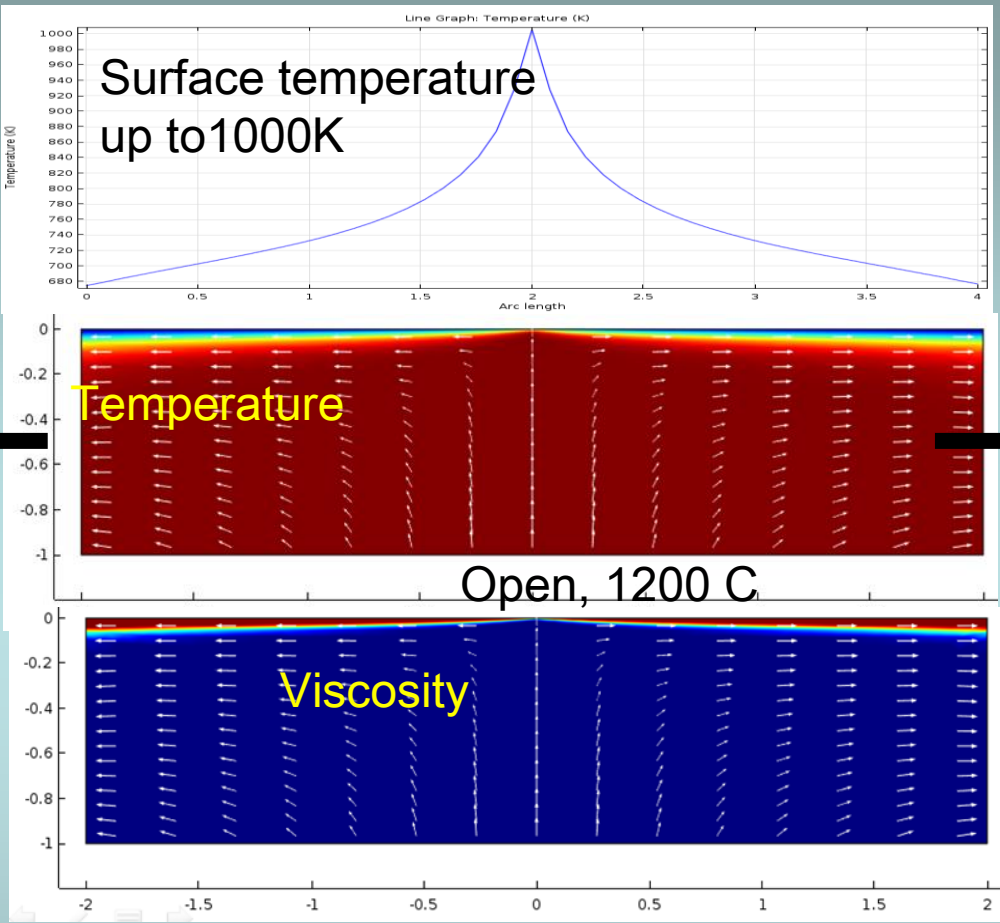
$$q = \varepsilon\sigma(T^4 - T_{amb}^4)$$

ε - surface emissivity (0.9)

σ - Stefan-Boltzmann constant



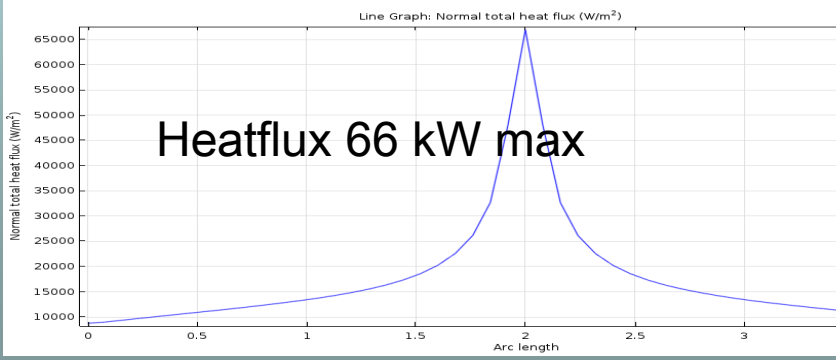
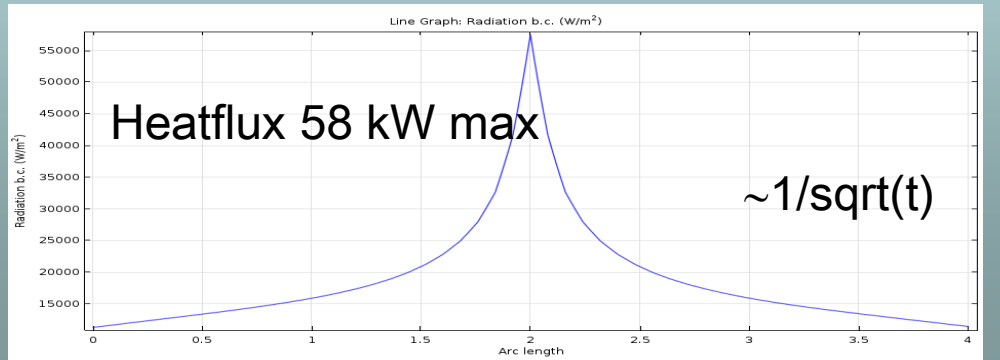
Spreading on lava lake



Results:

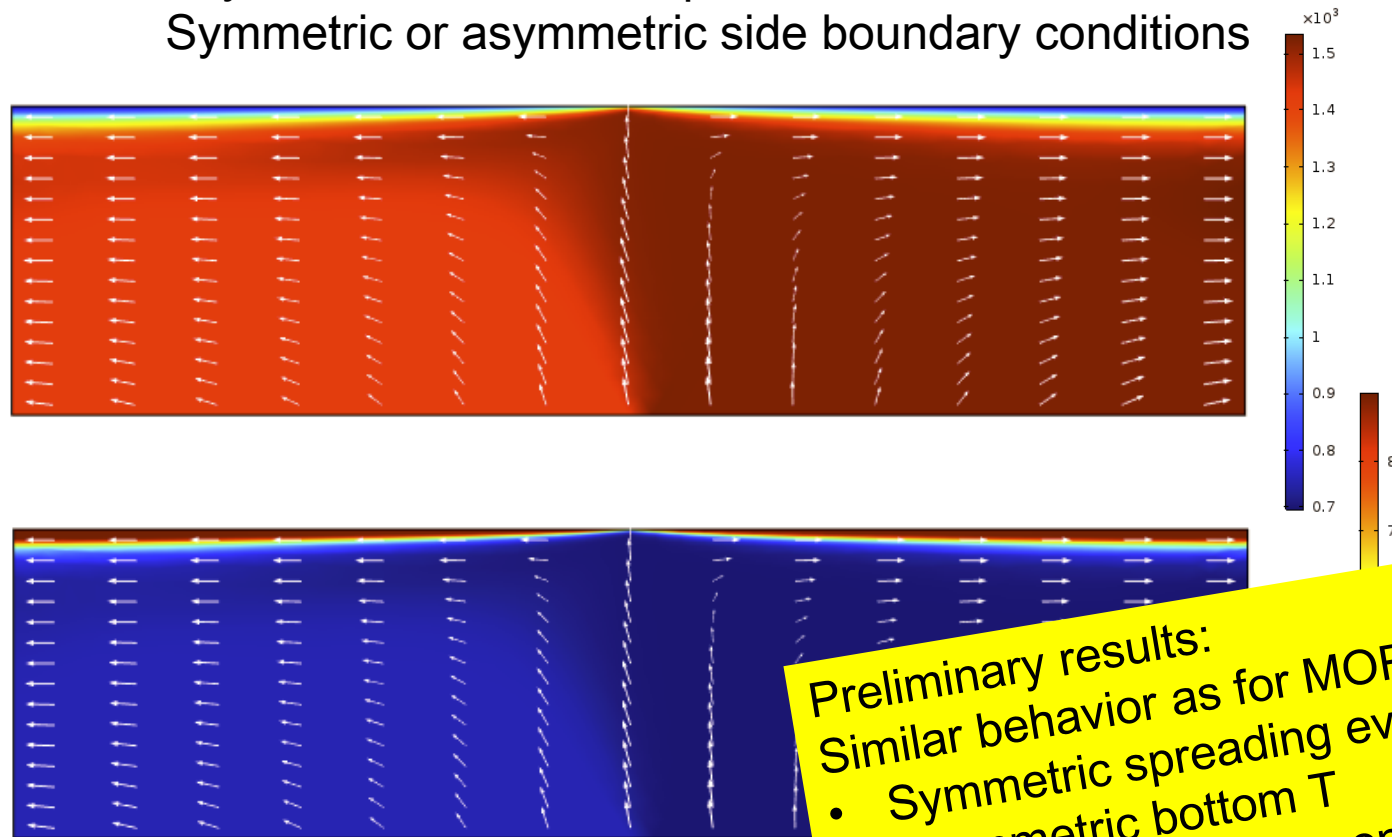
- Similar behavior as with const T boundary condition
- Symmetric spreading

Compare: constant surface temperature 773 K



Attempts to obtain asymmetric spreading....

Asymmetric bottom temperature
Symmetric or asymmetric side boundary conditions

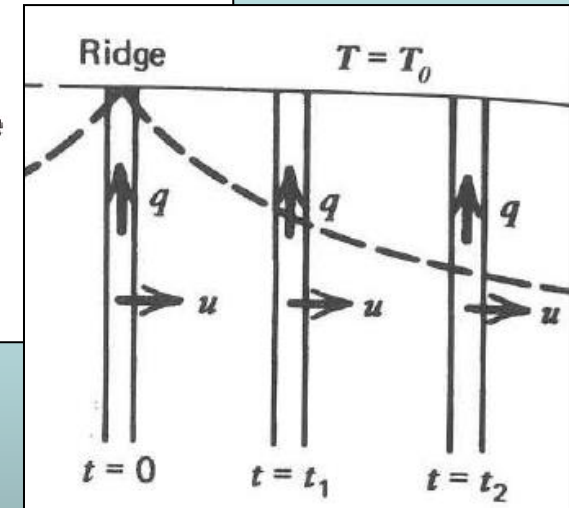
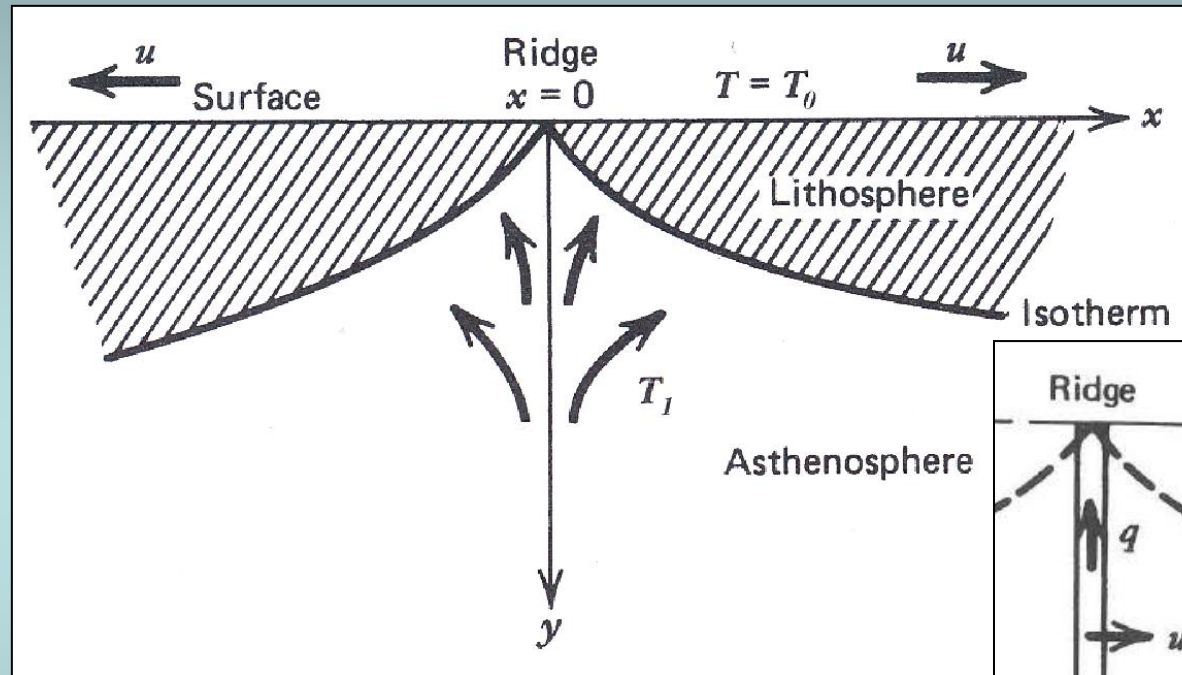


- Preliminary results:**
- Similar behavior as for MOR:
 - Symmetric spreading even for asymmetric bottom T
 - Radiative boundary condition similar to const T b.c.
 - Observed asymmetries: to be explored

Cooling of oceanic lithosphere with hydrothermal convection

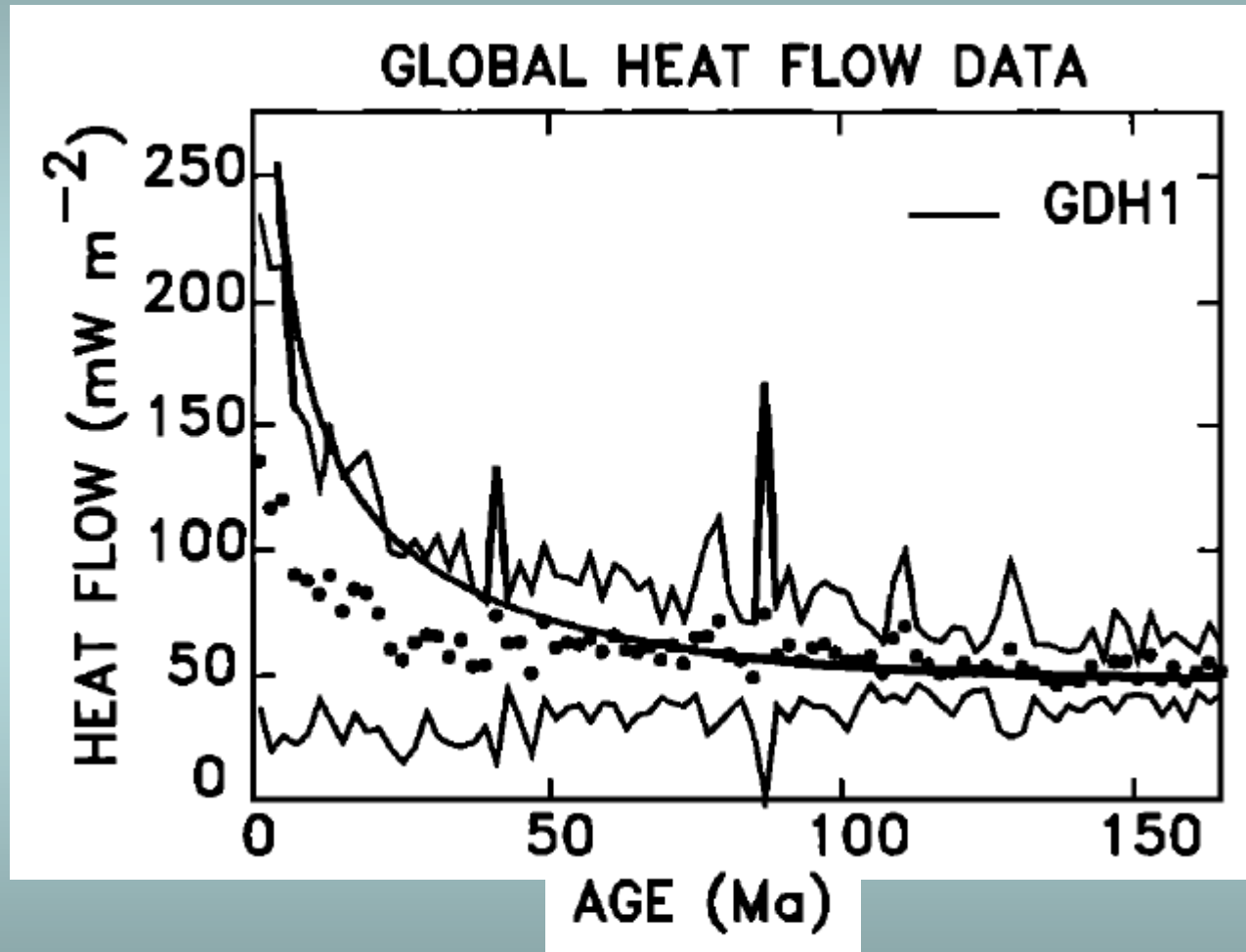
Inversion of observed heatflow and bathymetry

Cooling lithosphere, \sqrt{t} - law



Turcotte & Schubert:

$$q_0(t) = \lambda \cdot \left. \frac{dT}{dz} \right|_{z=0} = - \frac{\lambda \cdot \Delta T}{\sqrt{\pi} \sqrt{\kappa \cdot t}}$$

Cooling lithosphere, \sqrt{t} - law

Cooling lithosphere, \sqrt{t} - law

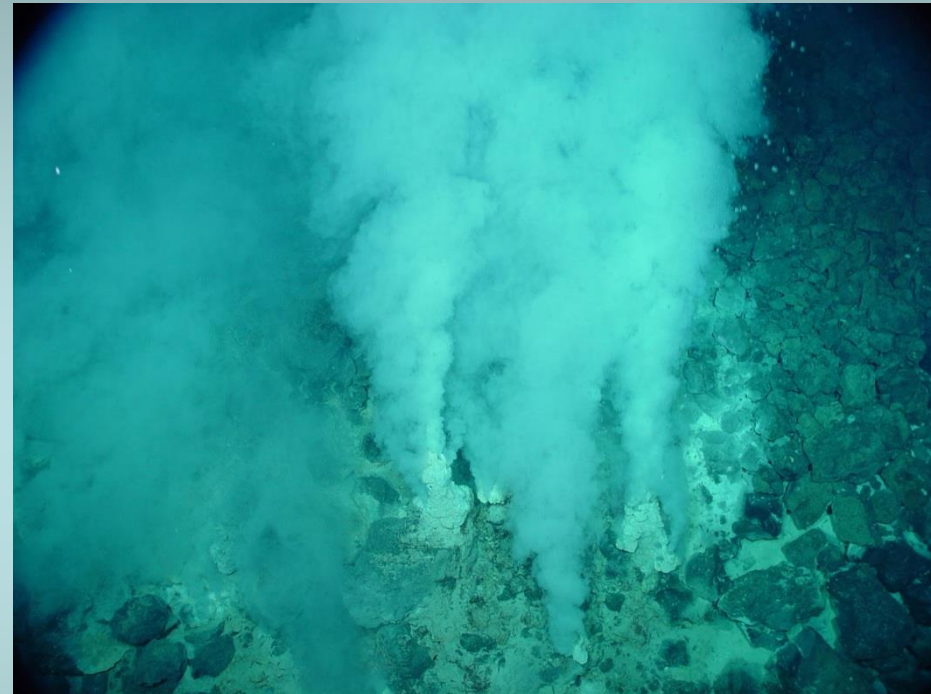
Discrepancy: hydrothermal cooling

Fraction of hydrothermally removed heat: **20 - 40%** of total heat flow of the earth (Sclater et al., 1980; Stein and Stein, 1992; Lowel et al., 2008; Spinelli and Harris, 2011)

But... no cooling plate model exists which consistently includes hydrothermal convection

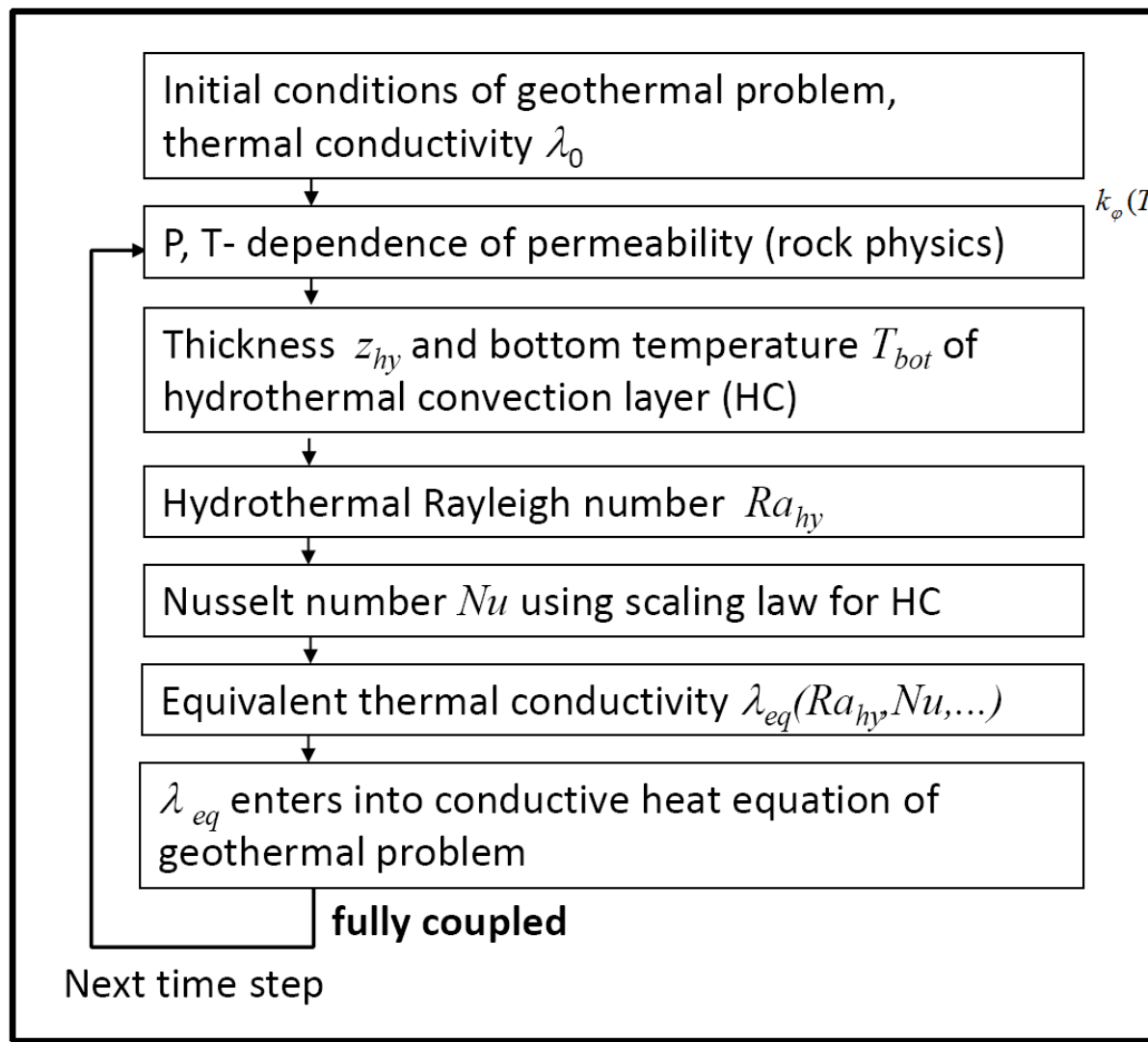
Effect on \sqrt{t} - law?

Fitting observations?



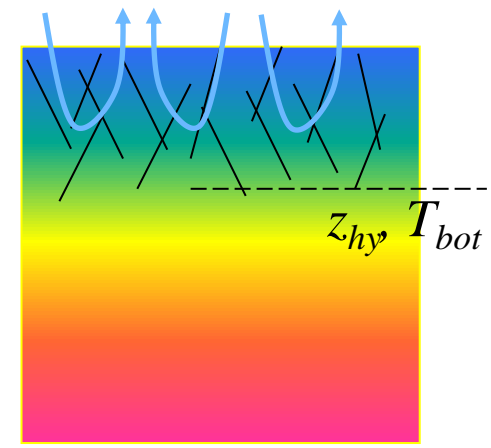
Our approach: Parameterized hydrothermal convection and equivalent conductivity

Schmeling, H, and G. Marquart, 2014: A scaling law for approximating porous hydrothermal convection by an equivalent thermal conductivity: theory and application to the cooling oceanic lithosphere. Geophys. J. Int., 197 (2): 645-664



$$k_{\varphi}(T, z) = k_{\varphi 0} \cdot [1 - \gamma \cdot (T - T_0)]^n \cdot H(1 - \gamma \cdot (T - T_0)) \cdot$$

$$H(T_{duc} - T) \cdot \exp\left(-n \cdot \frac{z}{z_{ch}}\right)$$

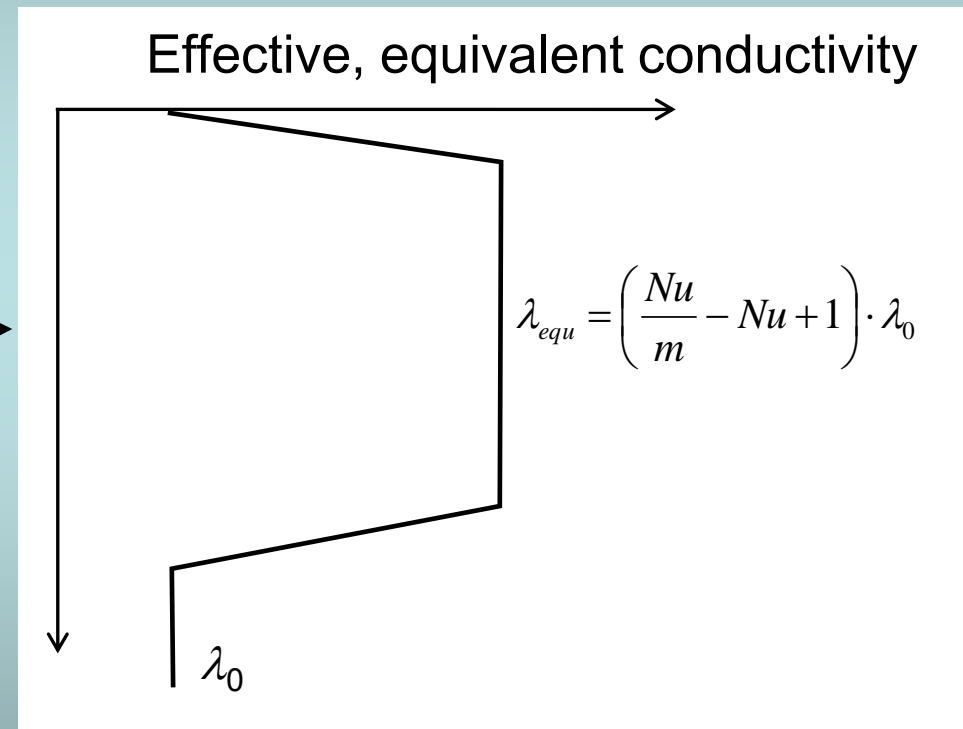
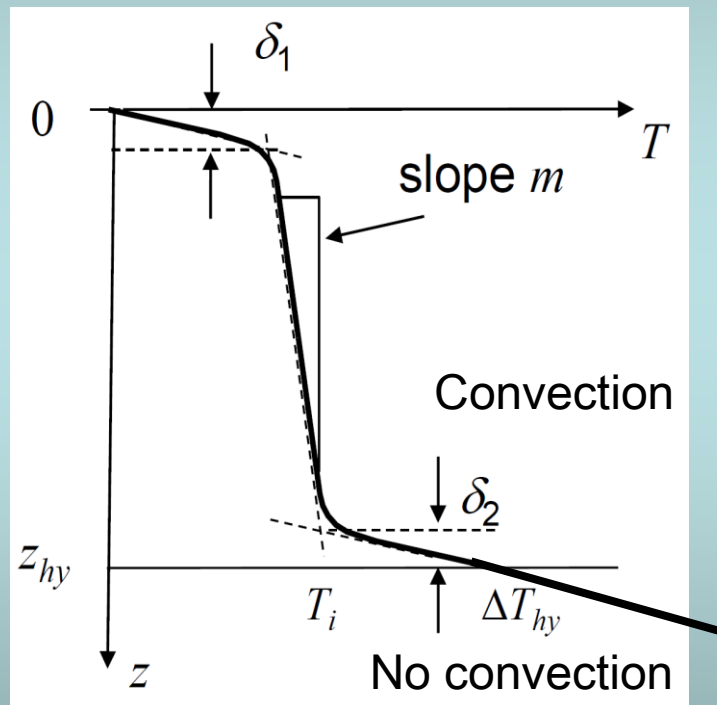


$$Ra_{hy} = \frac{\alpha_f g \rho_f^2 c_{pf} k_{\varphi 0} \Delta T_{hy} z_{hy}}{\lambda_m}$$

$$Nu = \frac{q_{top} z_{hy}}{\lambda_0 \Delta T_{hy}}$$

Equivalent thermal conductivity

Approximate convective layer by conductive layer with an effective, higher thermal conductivity



Solve 1D equation for cooling plate with simulated hydrothermal convection

Conductive 1D heat equation with $\lambda_{eq}(z,t)$ based on Nu, z_{hy} etc from parameterized hydrothermal convection:

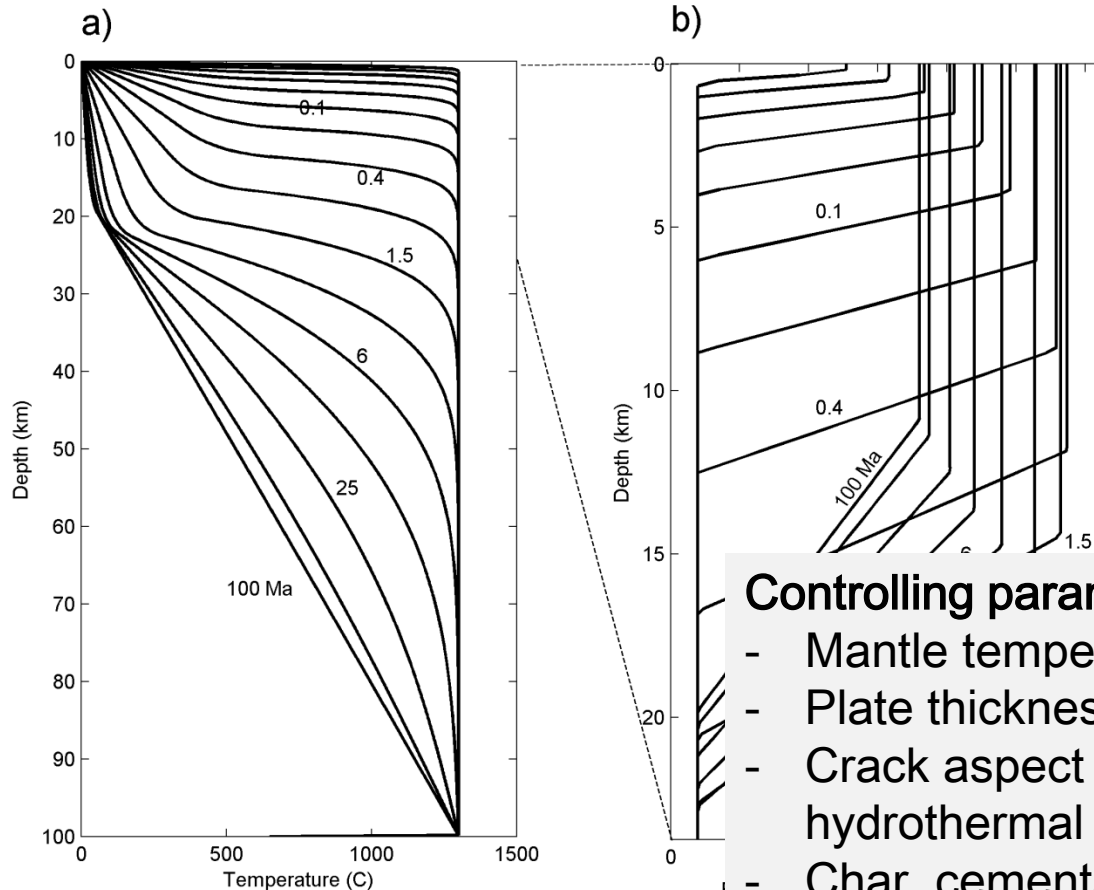
$$\frac{\partial T}{\partial t} = \frac{1}{\rho c_p(T)} \frac{\partial}{\partial z} \left(\lambda_{eq}(Nu, T, z) \frac{\partial T}{\partial z} \right) - \frac{T}{c_p(T)} \frac{\partial c_p(T)}{\partial t}$$

Hydrothermal convection

Time-dependent c_p
 → Few iterations
 per time step

- Solved with 4th order FD-scheme, forward in time
- Reduced diffusive time steps due to increased λ_{eq}

Example of cooling plate with simulated hydrothermal convection

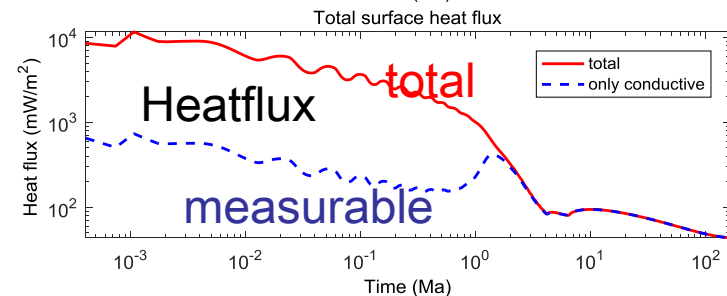
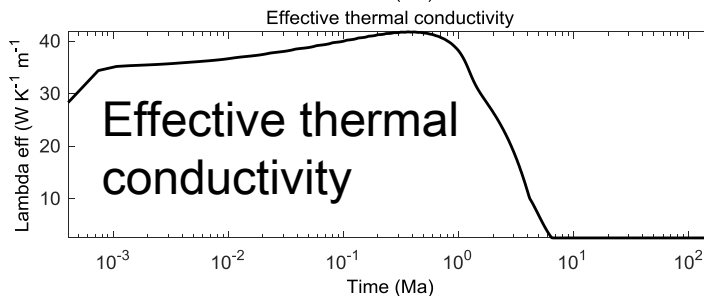
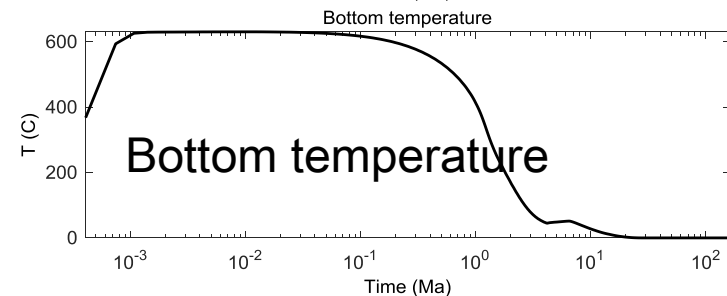
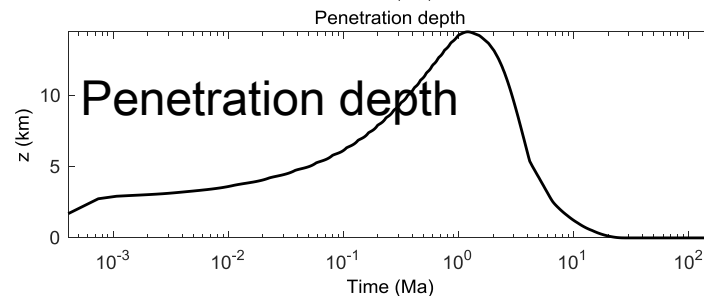
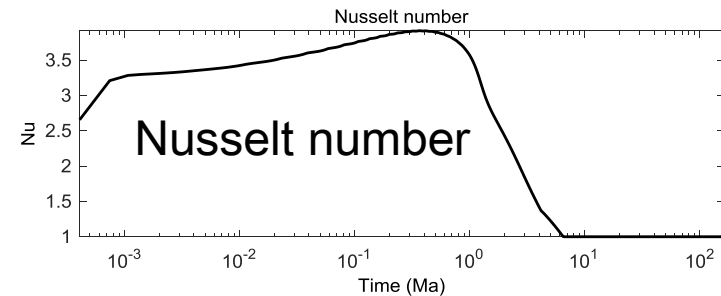
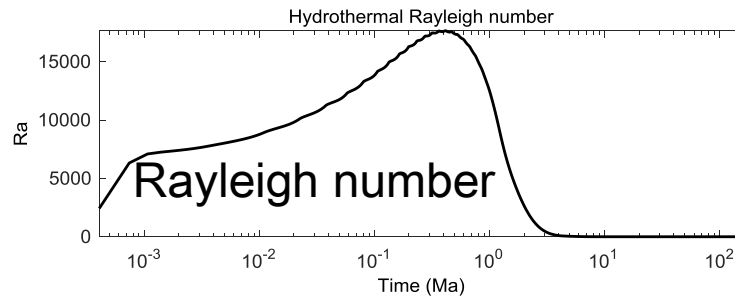


Solving the conductive 1D heat equation with $\lambda_{eff}(z,t)$ based on Nu , z_{hy} etc from parameterized hydrothermal convection

Controlling parameters:

- Mantle temperature
- Plate thickness
- Crack aspect ratio (controls permeability → hydrothermal convection penetration)
- Char. cementation time τ_{decay}
- Sealing time (open → covered convection)

Time evolution of a typical model



3 Phases:

0 – 1 Ma increasing vigor of hydrothermal convection

1 – 10 Ma declining convective vigor due to cooling and sealing

> 10 Ma no convection, only conductive cooling and cementation

Optimizing parameters by downhill simplex inversion

Parameters to be optimized:

With hydrothermal convection:

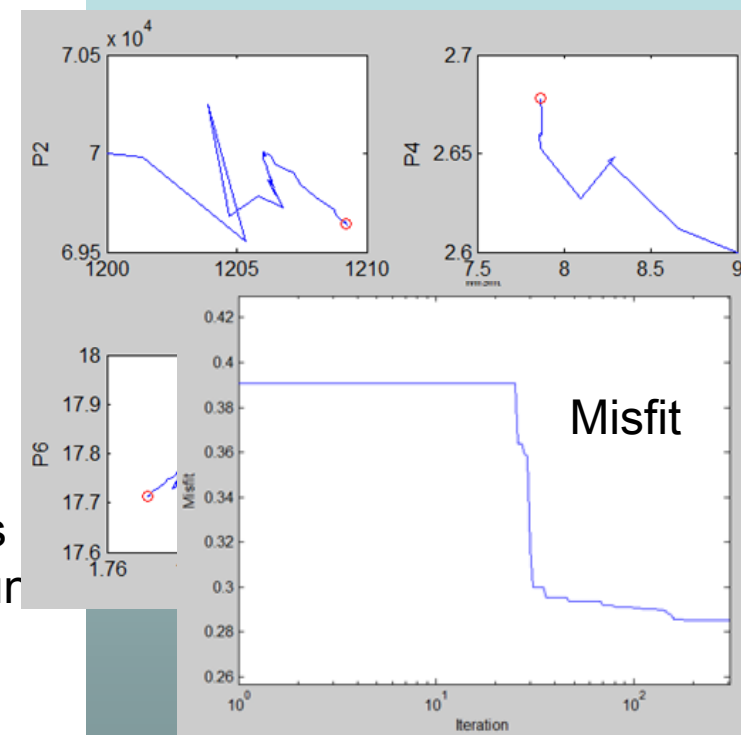
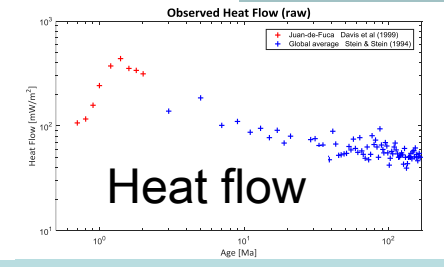
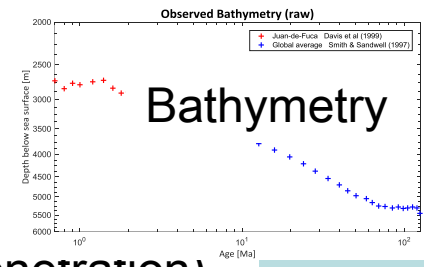
- Mantle temperature
- Plate thickness
- Crack aspect ratio (controls permeability → hydrothermal convection penetration)
- Char. cementation time τ_{decay}
- Sealing time (open → covered convection)
- Prefactor thermal expansivity of lithosphere

Without hydrothermal convection:

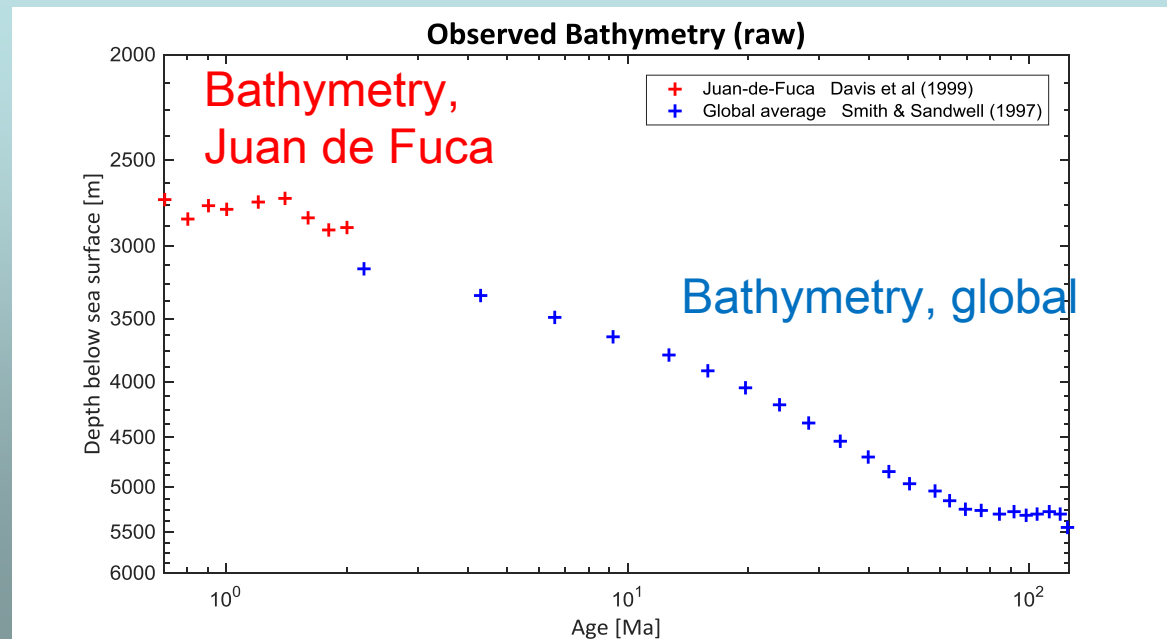
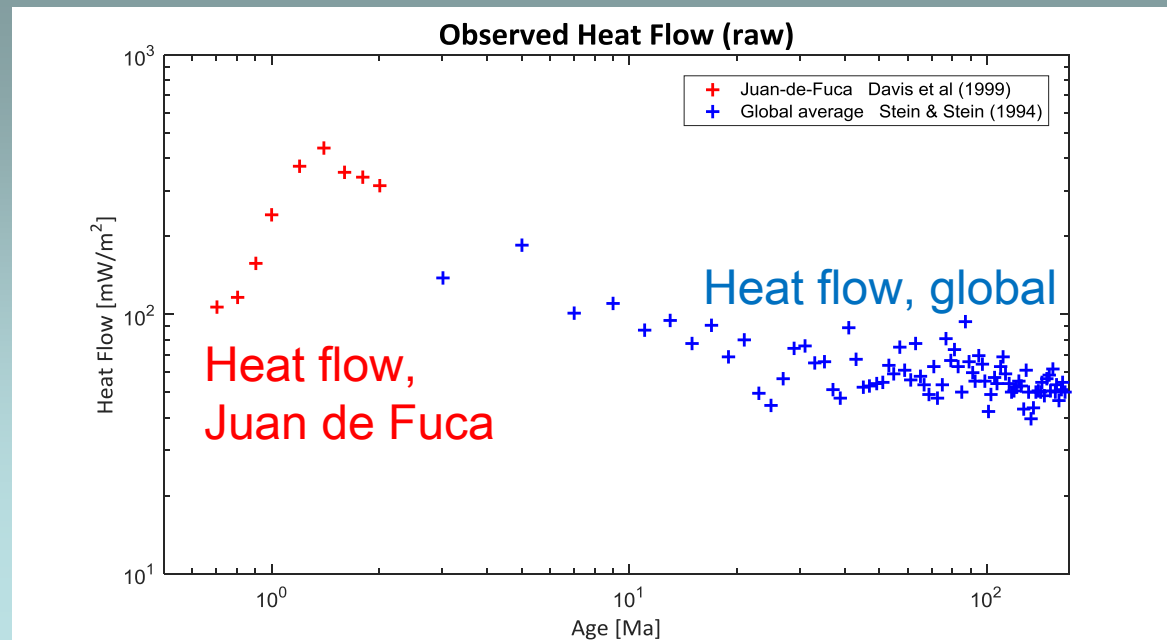
- Mantle temperature
- Plate thickness
- Thermal expansivity of lith
- Downhill simplex inversion run: typically 400 models
- Depends on choice of starting parameters → 200 run → O(80 000) cooling models

Stronger weight (factor 5) for young lithosph

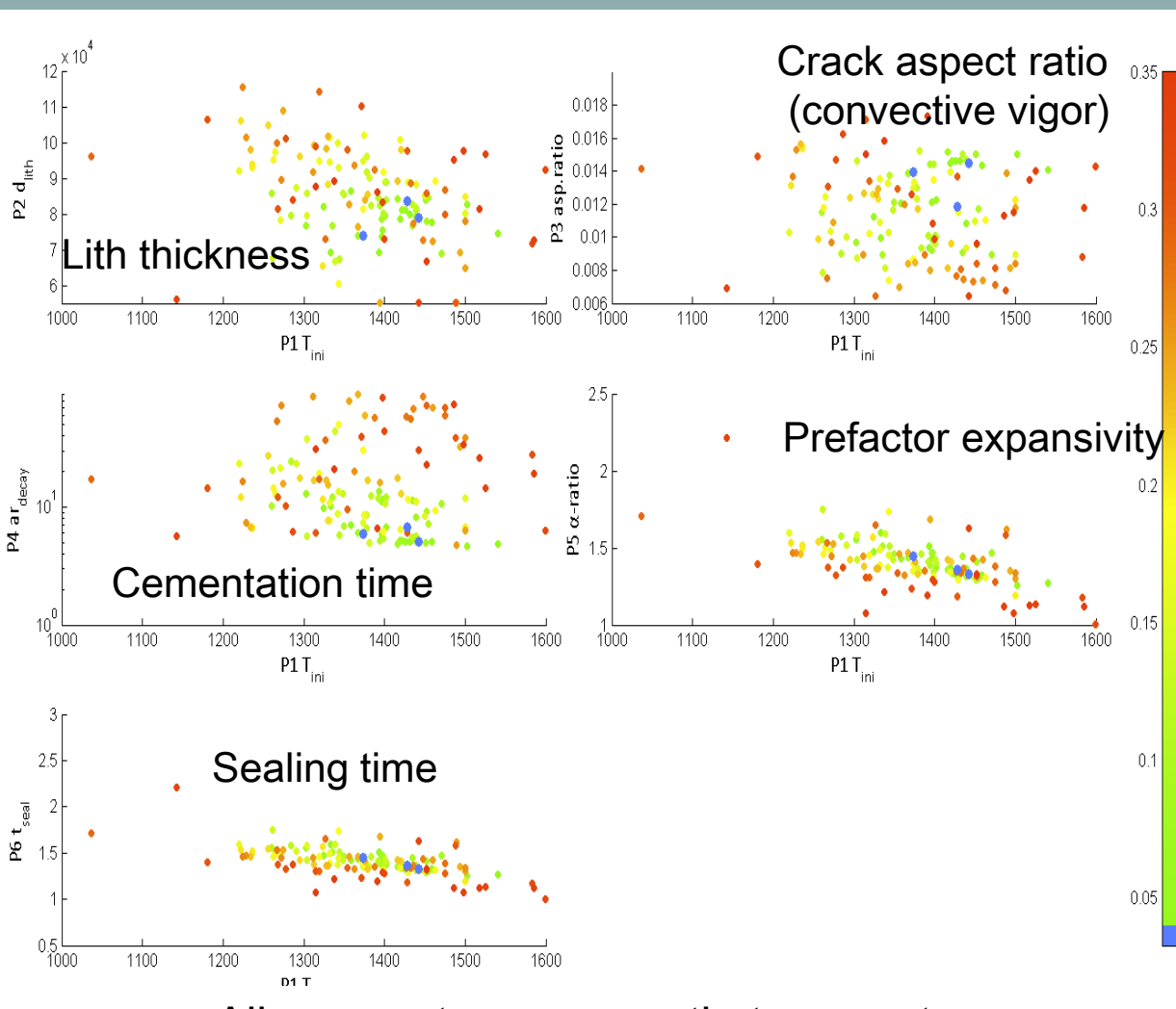
Data to be inverted



Data to be inverted



Different projections of best models within parameter space



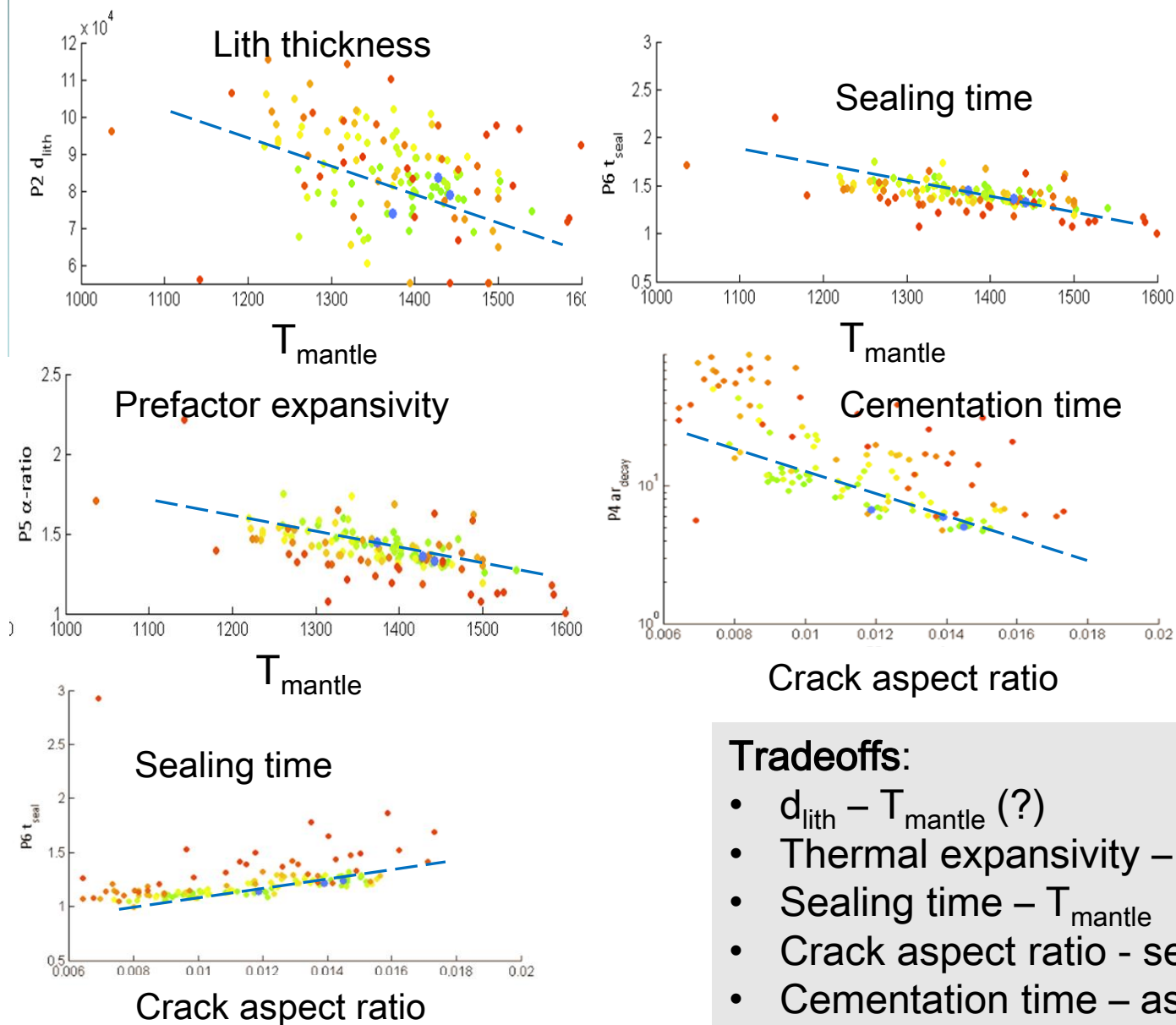
→ Many local minima in 6D parameter space

Best parameters

- 1350 – 1450 °C mantle
- 70 – 90 km lithosphere
- 0.01 – 0.015 aspect ratio, i.e. a narrow range of convective vigor
- 5 – 10 Ma cementation time
- Narrow range sealing time 1.3 – 1.6 Ma
- 1.3 – 1.5 prefactor lith expansivity

All parameters vs. mantle temperature

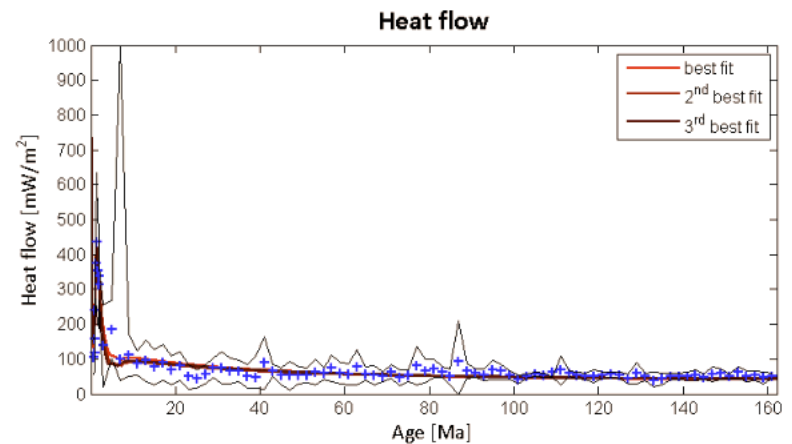
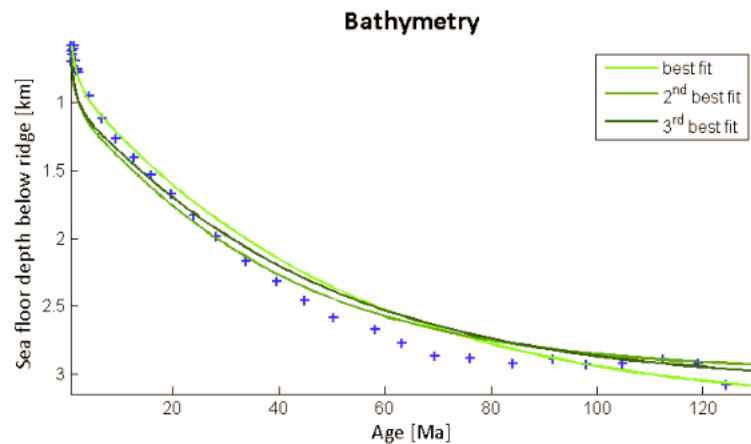
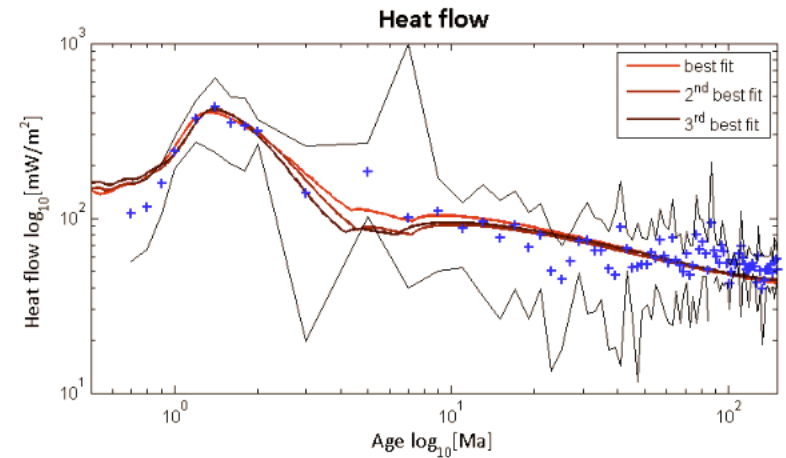
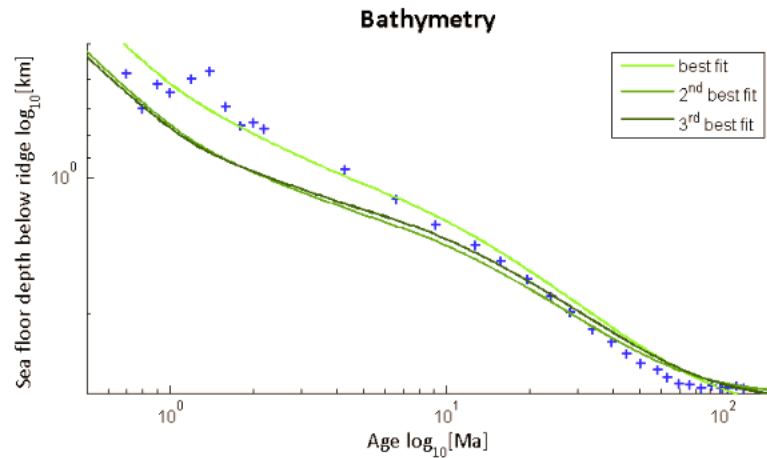
Tradeoffs for different projections



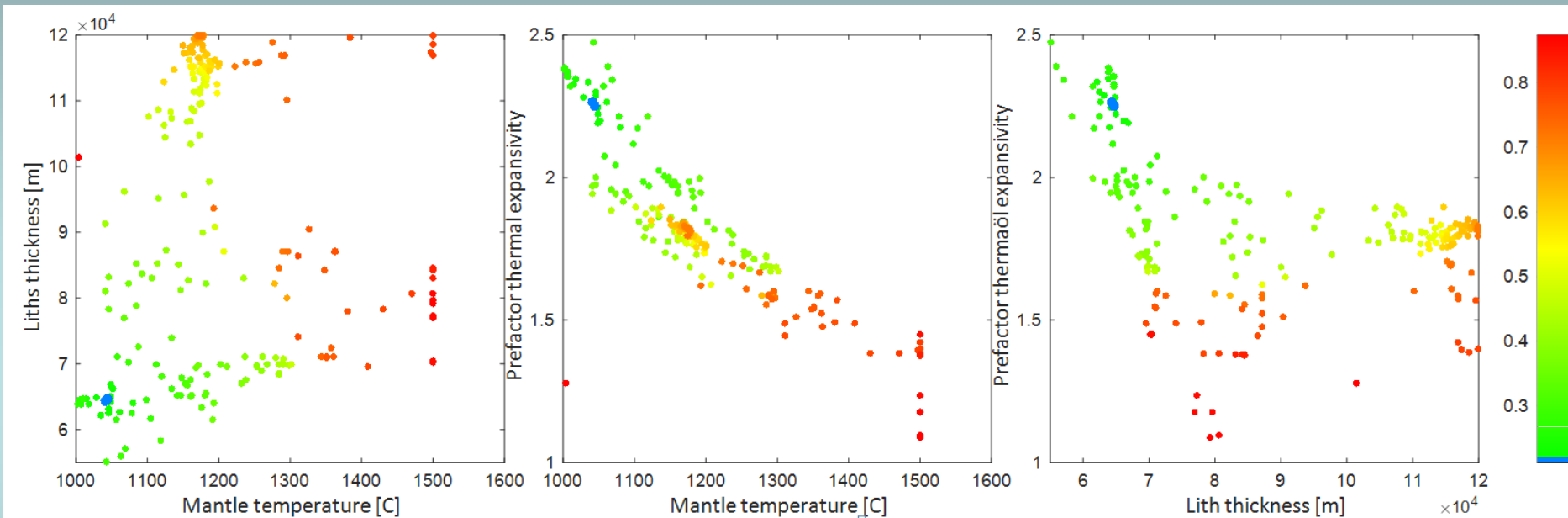
Tradeoffs:

- $d_{\text{lith}} - T_{\text{mantle}}$ (?)
- Thermal expansivity – T_{mantle}
- Sealing time – T_{mantle}
- Crack aspect ratio - sealing time
- Cementation time – aspect ratio

3 best models and the observations



Inversion without allowing for hydrothermal convection



Without hydrothermal convection

- T_{mantle} 1000 – 1200C
- Lith thickness 60 – 85 km
- Prefactor expansivity 1.6 – 2.4
- Misfit never below 0.2

With hydrothermal convection

- T_{mantle} 1350 – 1450 C
- Lith thickness 70 – 90 km
- Prefactor expansivity 1.3 – 1.5
- **Misfit down to 0.03**

Conclusions

- Self-consistent modelling: **Mostly symmetric** spreading with and without ridge migration for many spreading velocities, plasticity factors, even asymmetric temperatures
- Both MOR and models show asymmetric spreading **up to 15%**
- MOR slightly **slows down** when approaching and overriding the plume
- Physics of observed asymmetric spreading (lava lakes) **not clear**

- Lithospheric cooling with hydrothermal convection → significantly **better fit** of bathymetry and heatflux for young plates
- T_{mantle} 1350 – 1450 °C, Lith thickness 70 – 90 km
- Prefactor expansivity 1.3 – 1.5 (i.e. higher than from lab experiments)
- Sealing time 1.3 – 1.6 Ma, cementation time 5 – 10 Ma

Thank you for your attention



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