

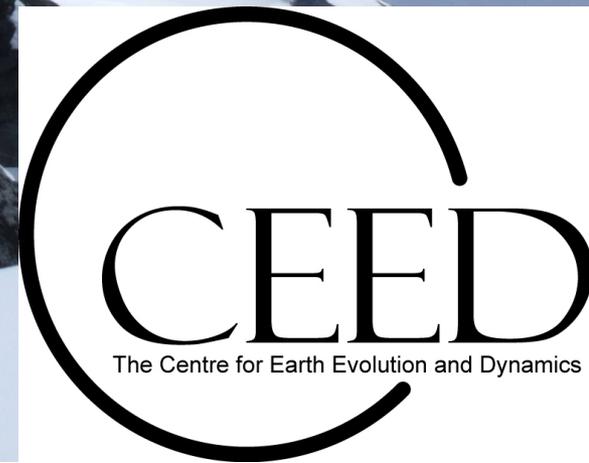
Modelling dynamic topographies of planetary bodies

Bernhard Steinberger

Deutsches GeoForschungsZentrum, Potsdam

and

Centre for Earth Evolution and Dynamics, Univ. Oslo



Outline

→ What is dynamic topography and why is it important?

→ Numerical models of dynamic topography on Earth and comparison to observations

- Instantaneous flow based on tomography

- (a) only radial viscosity variations

- (b) coupling with more realistic lithosphere

- Time-dependent flow

- (a) backward advection based on tomography

- (b) forward models based on subduction history

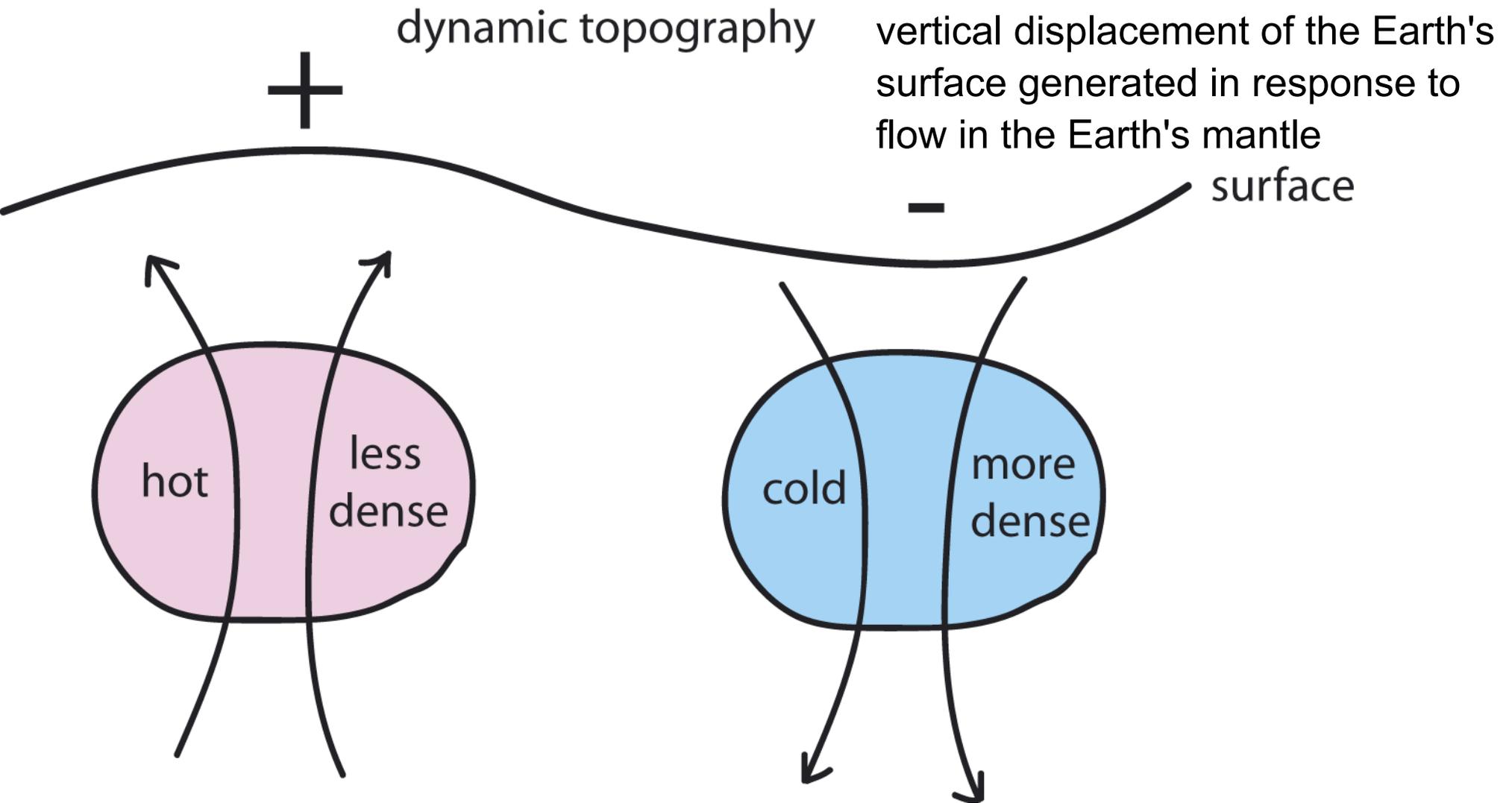
- (c) adjoint models

→ Implications

→ Other planets: equipotential surface and topography

- Instantaneous models (indirectly based on tomography)

- Forward convection models



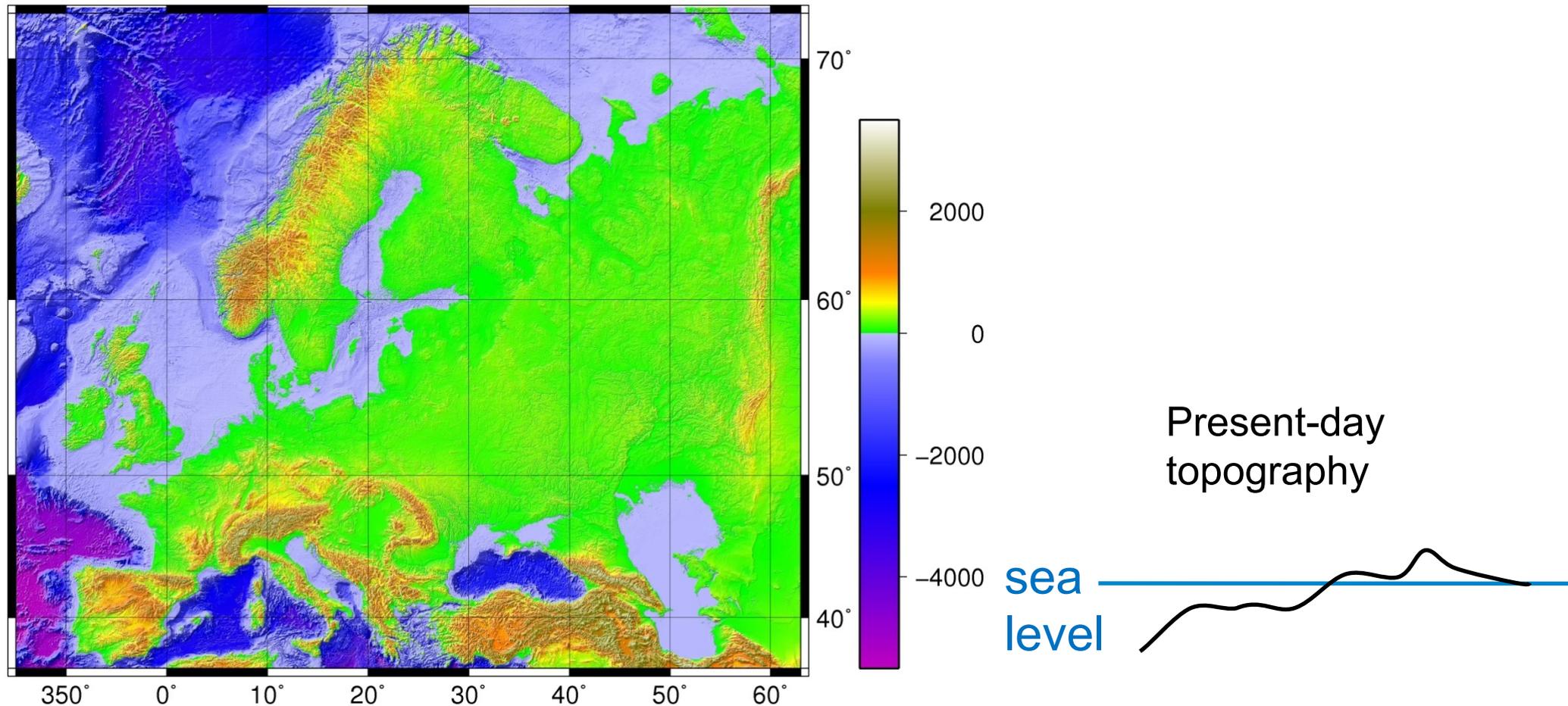
Lower mantle heterogeneity, dynamic topography and the geoid

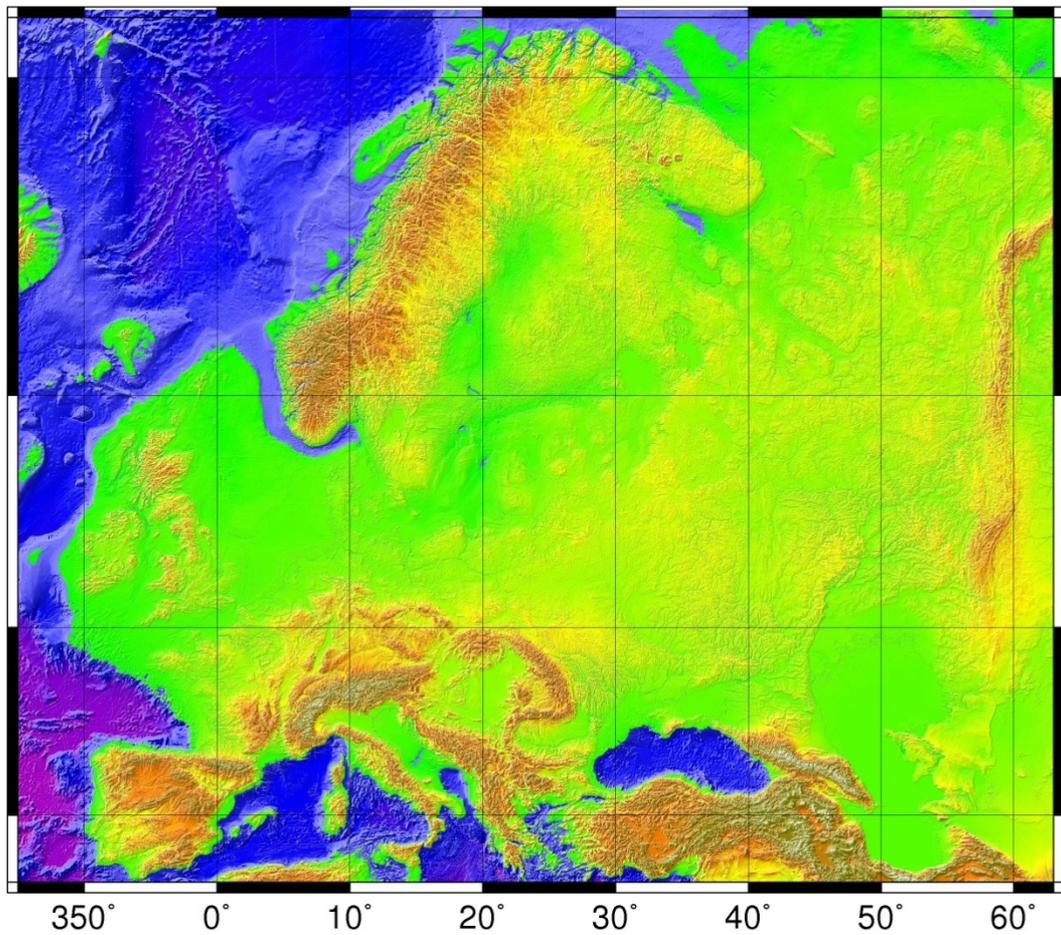
Bradford H. Hager, Robert W. Clayton, Mark A. Richards, Robert P. Comer*
& Adam M. Dziewonski*

Why is it important to know dynamic topography?

Many areas on Earth within few hundred meters above or below sea level (bright green / light blue on map)

Dynamic topography expected to reach a few hundred meters and hence may influence when and where sediments and natural resources may form





70°

2000

0

-2000

-4000

350° 0° 10° 20° 30° 40° 50° 60°

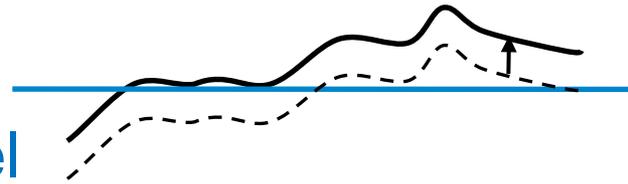
60°

50°

40°

Present-day
topography + 200 m

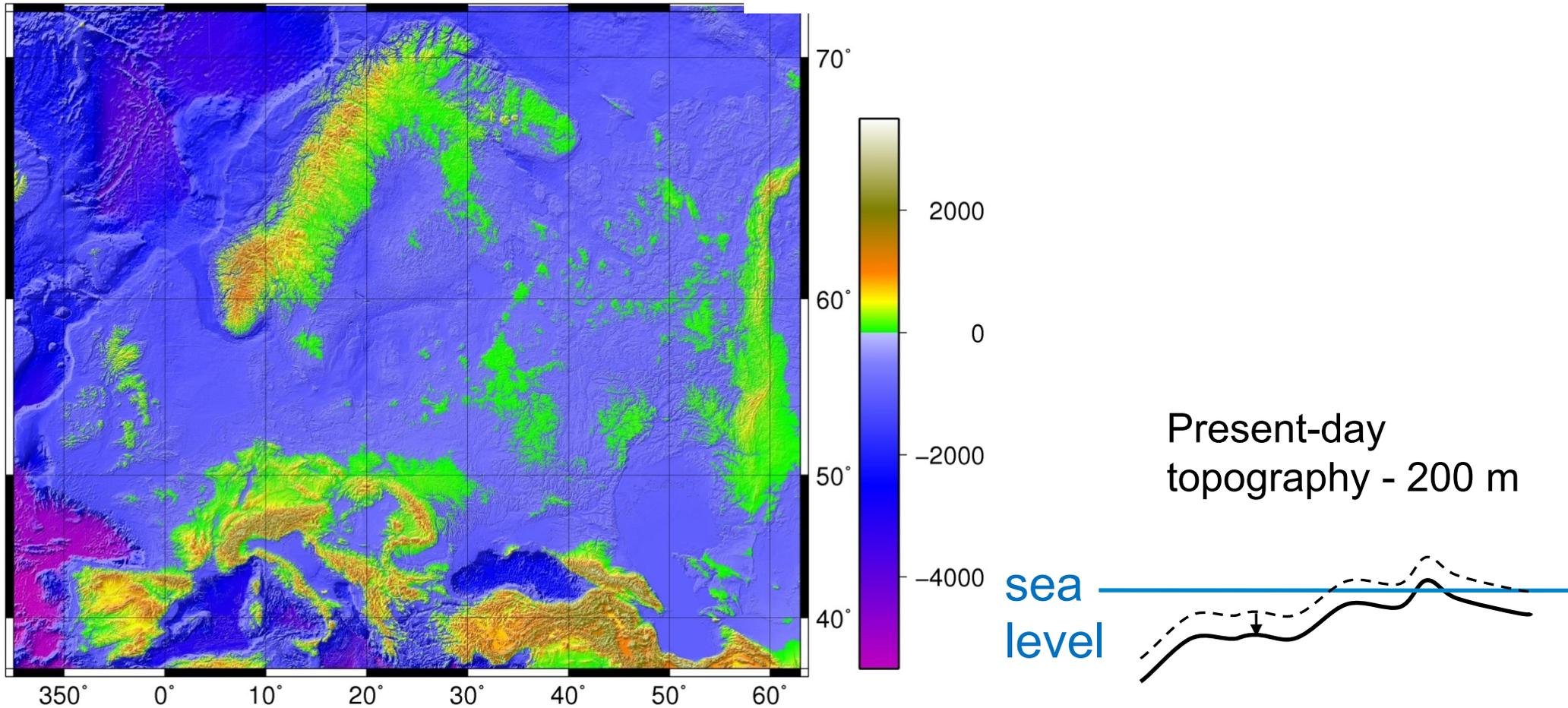
sea
level



Bounds on global dynamic topography from Phanerozoic flooding of continental platforms

Michael Gurnis

Department of Geological Sciences, University of Michigan, Ann Arbor, Michigan 48109-1063, USA



Dynamic topography changes ocean basin volume and hence sea level
Figure from Conrad and Husson
(Lithosphere, 2009)

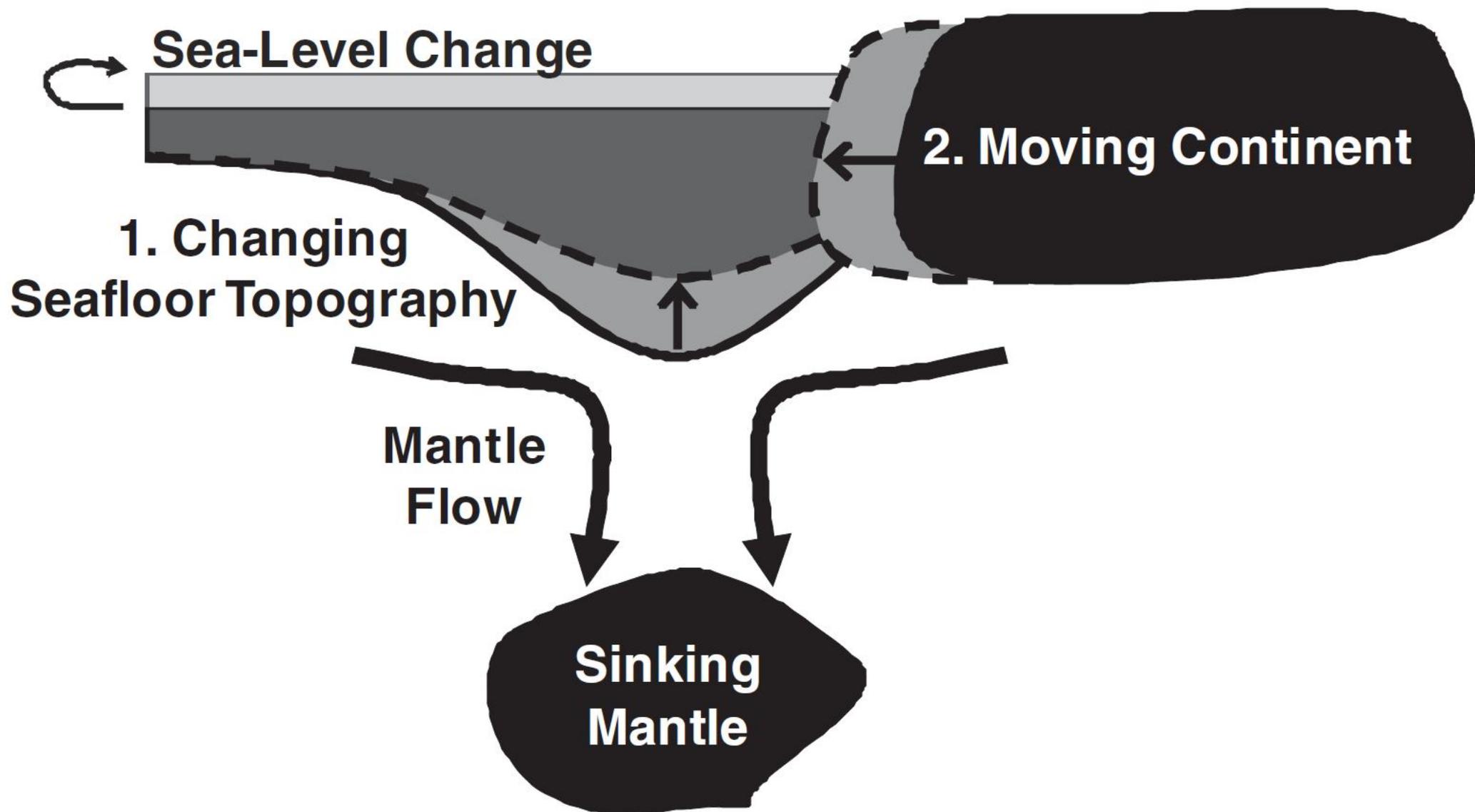
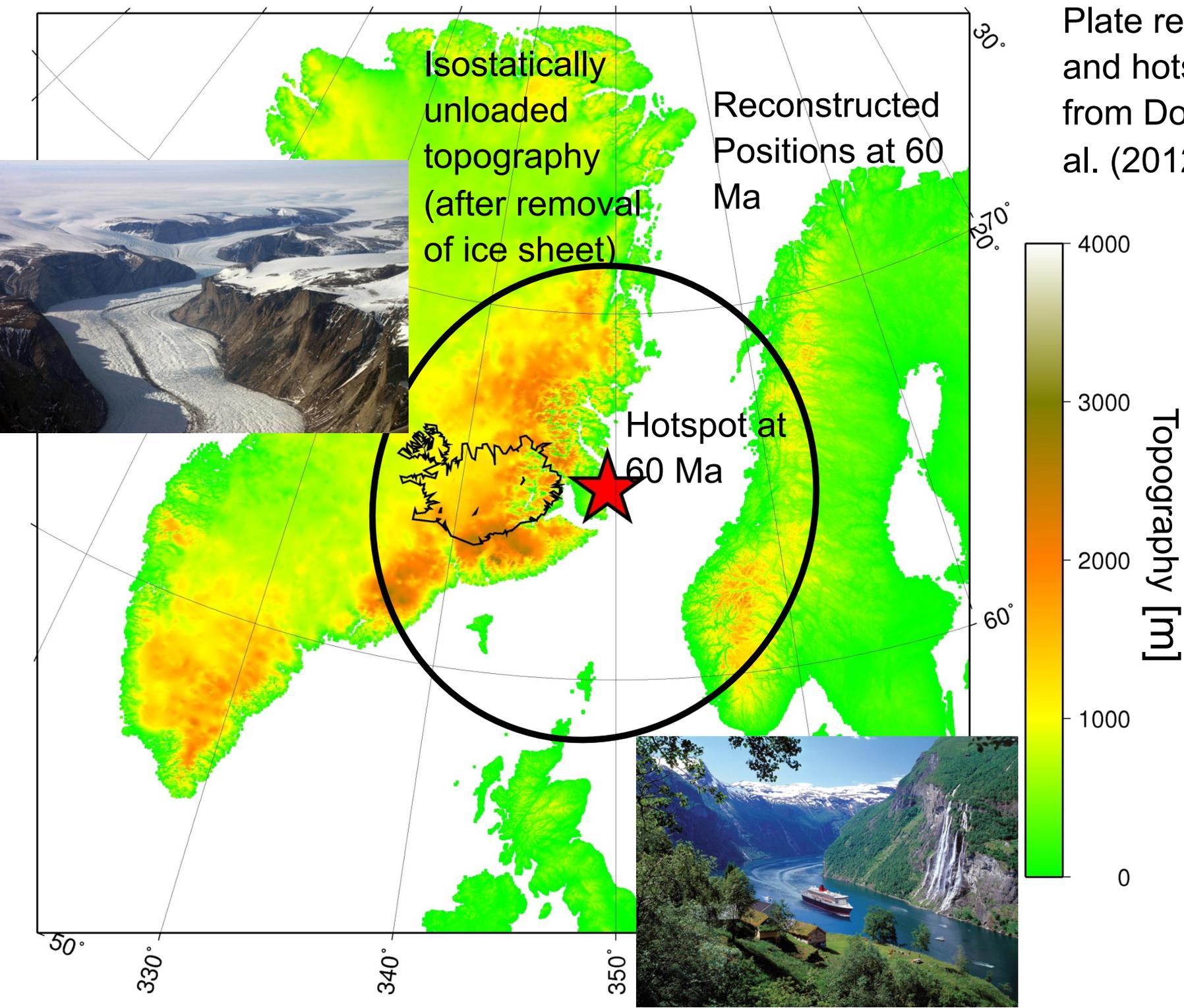


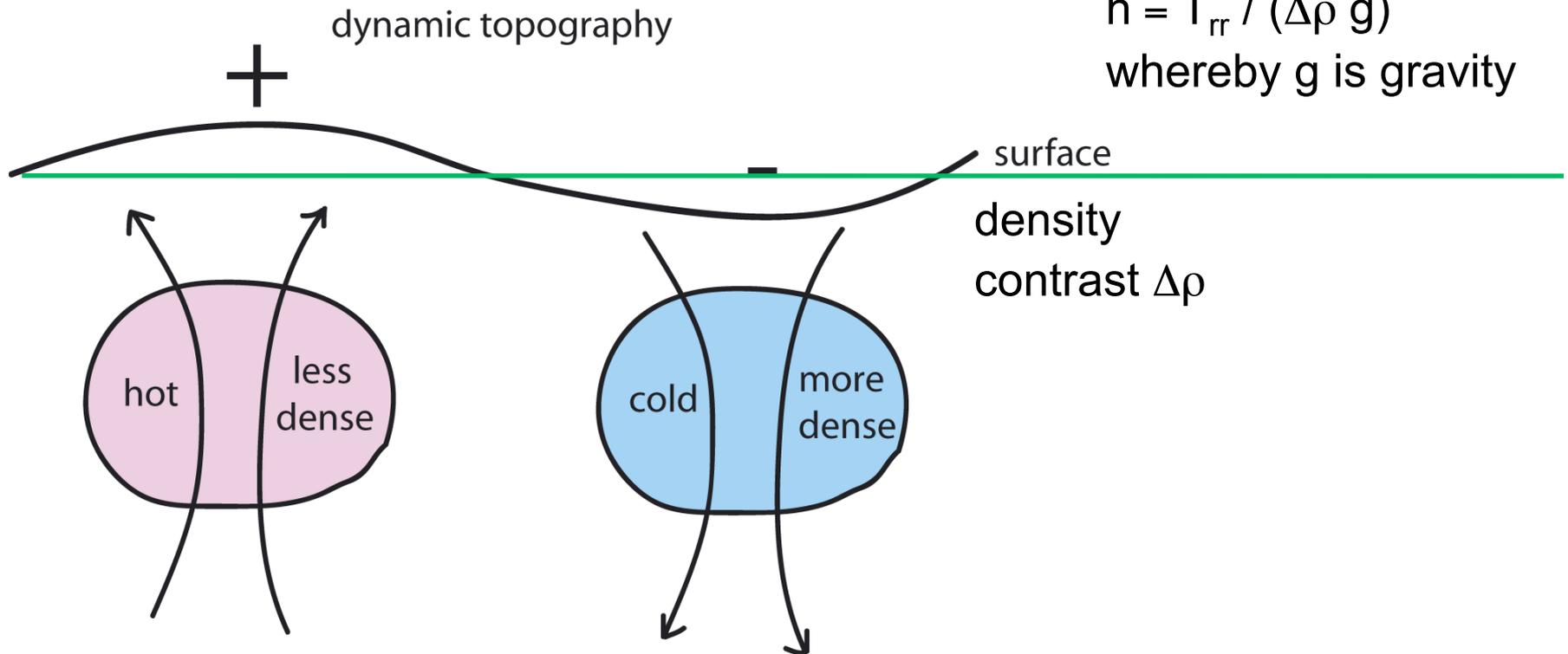
Plate reconstruction and hotspot motion from Doubrovine et al. (2012)



Computing dynamic topography

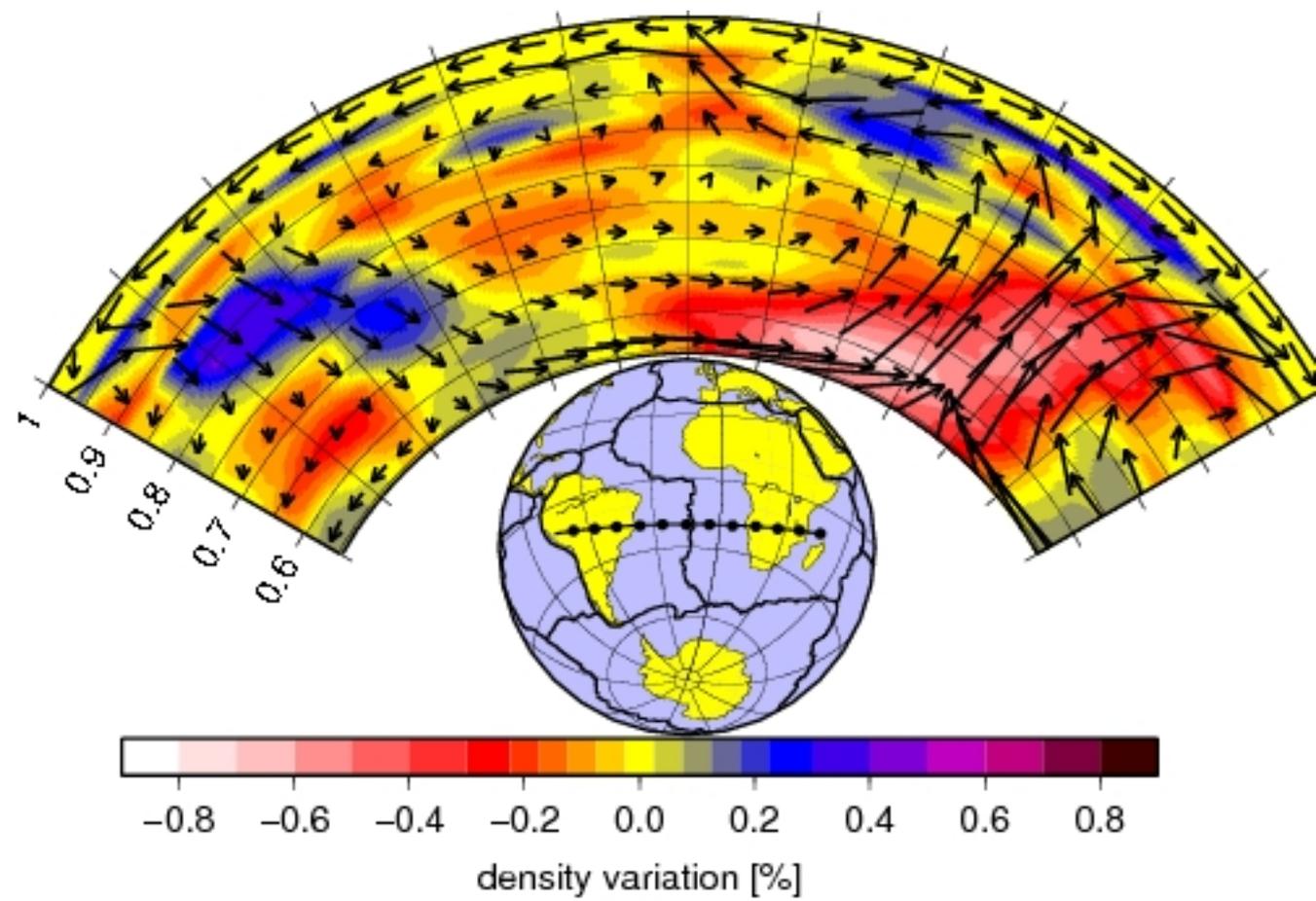
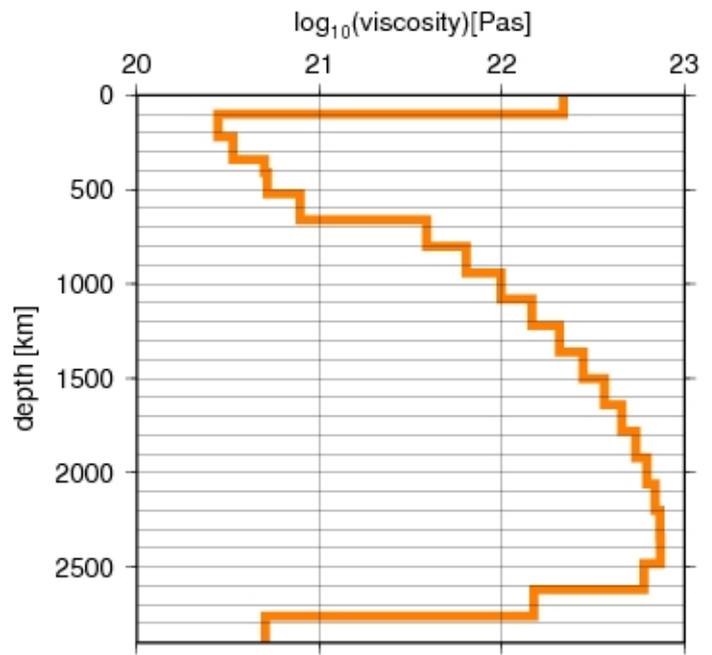
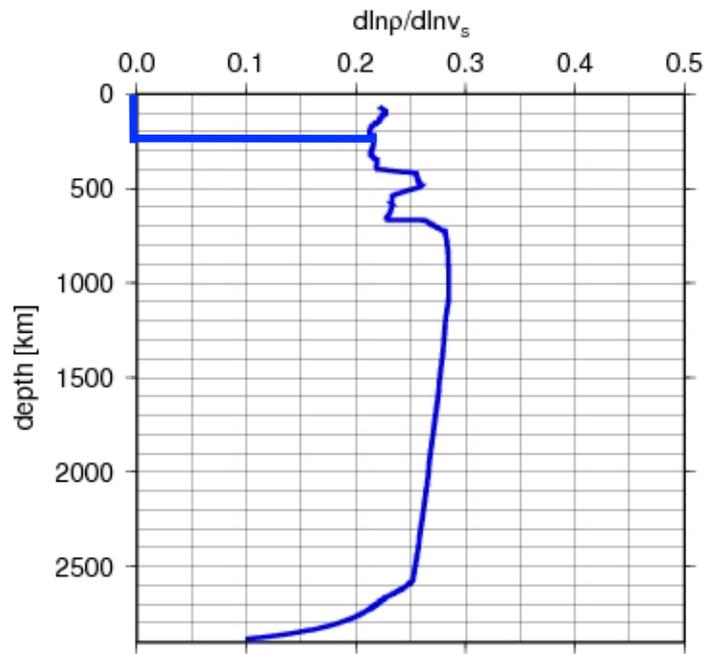
Numerical model:
Often fixed computational domain
No boundary displacement
Instead convert computed vertical stresses T_{rr} to topography h :
 $h = T_{rr} / (\Delta\rho g)$
whereby g is gravity

Reality: Boundary displacement

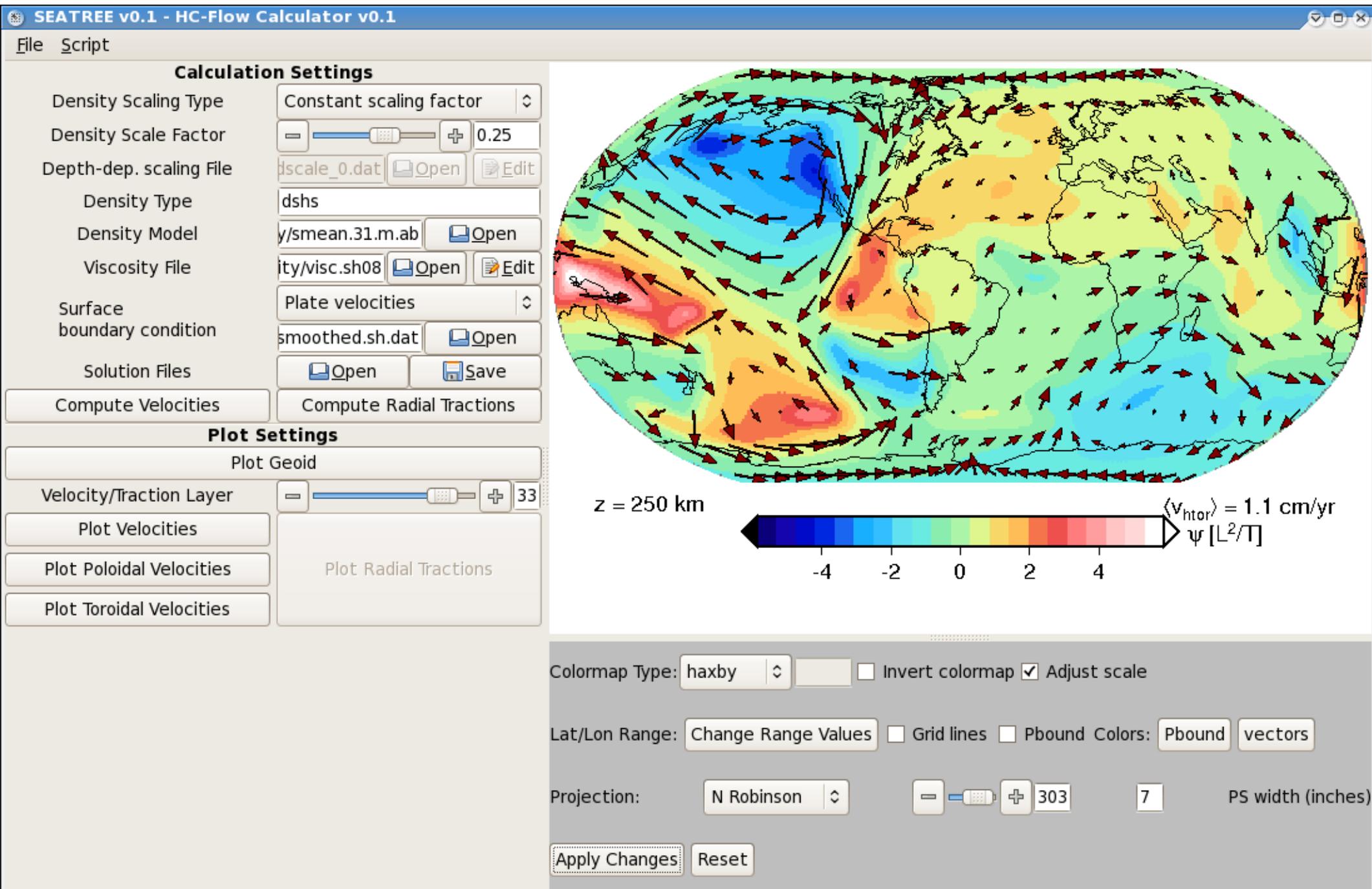


Instantaneous mantle flow computation

- Density model based on tomography (here: Simmons, Forte, Grand, 2006)
- Here, all density anomalies above 220 km are removed
- thermal velocity-density scaling based on mineral physics
- radial viscosity structure based on mineral physics and optimizing fit to geoid etc. (Steinberger and Calderwood, 2006)
- Spectral method (Hager and O'Connell, 1979, 1981)



Download flow code «HC» with user interface from <http://geodynamics.org/cig/software/hc>; courtesy of Thorsten Becker



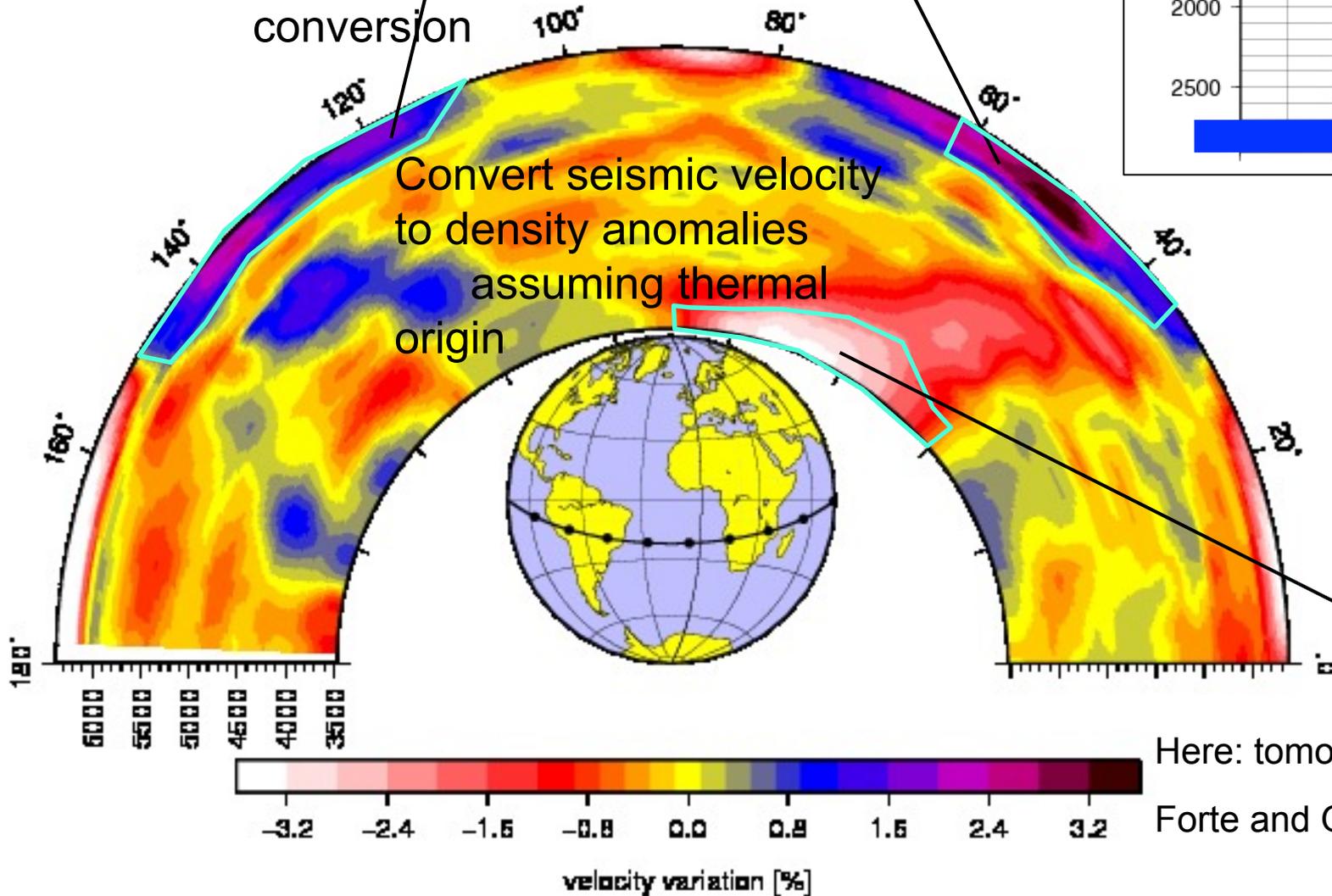
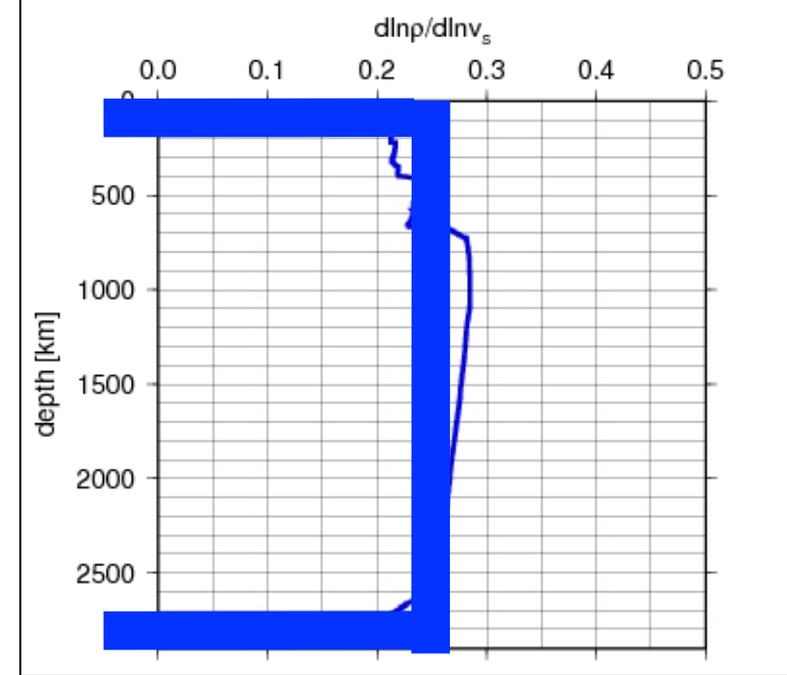
Computation for instantaneous flow based on seismic tomography

Within continental lithosphere, probably cannot use "thermal" conversion

Convert seismic velocity to density anomalies assuming thermal origin

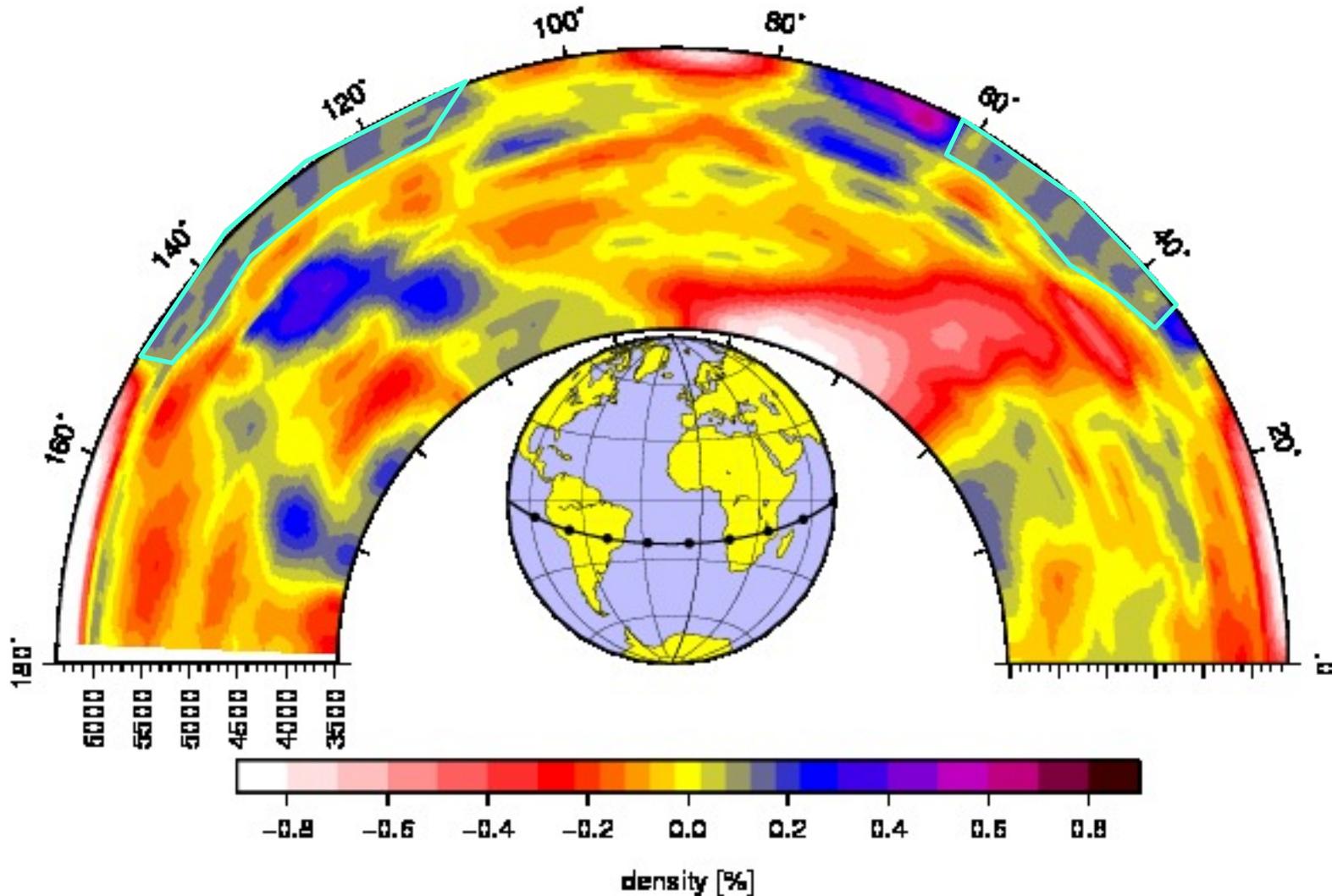
Conversion factor for thermal anomalies inferred from mineral physics (Steinberger and Calderwood, 2006)

Within LLSVPs, probably cannot use "thermal" conversion either

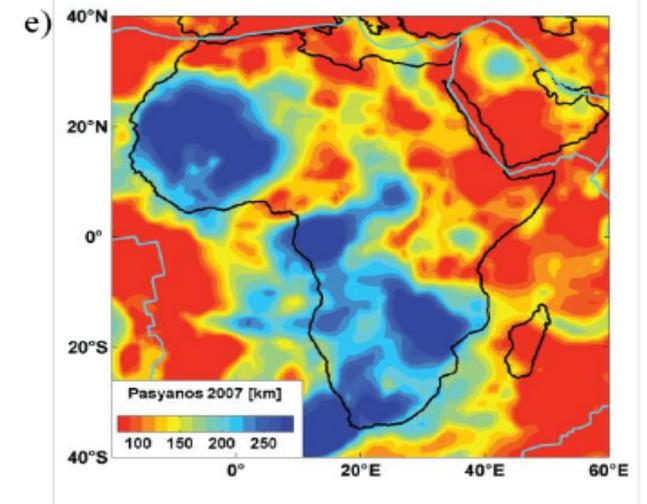
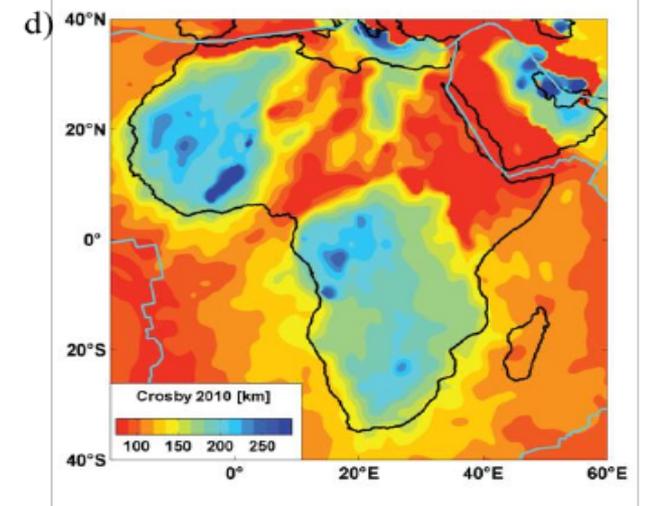
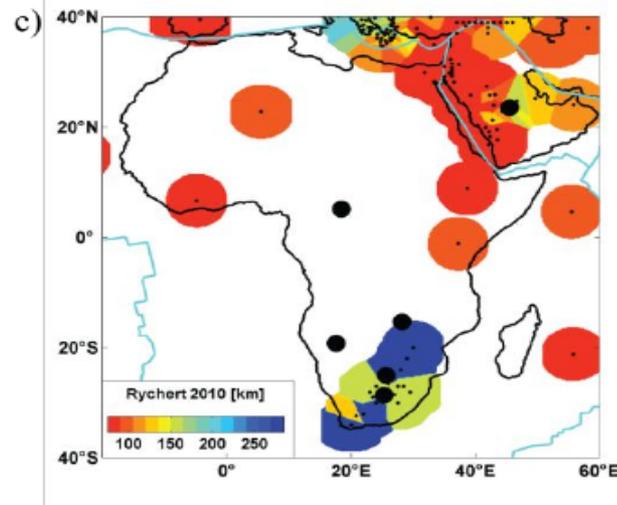
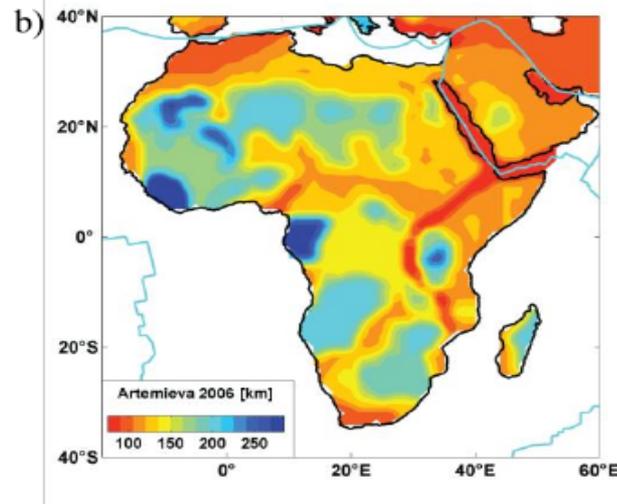
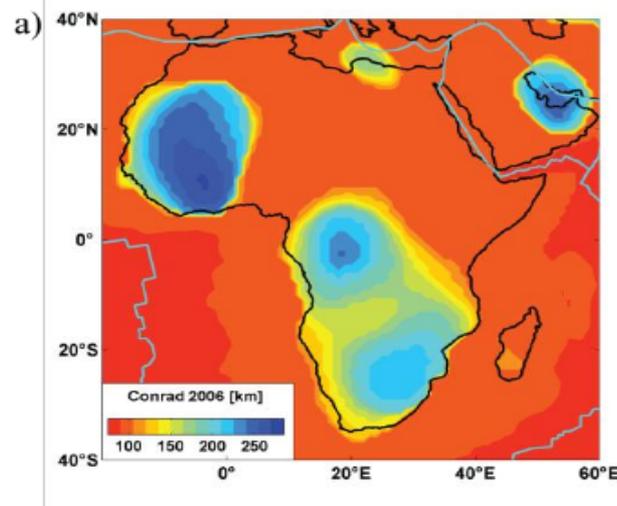


Here: tomography model of Simmons, Forte and Grand (2006)

Attempt to “remove lithosphere” by setting density anomaly to 0.2 % wherever, above 400 km depth and on continents, inferred density anomaly is positive >0.2 % at that depth and everywhere above



→ Lithospheric thickness not well constrained
→ models a, d and e based on tomography
→ b based on heat flow
→ c from receiver functions



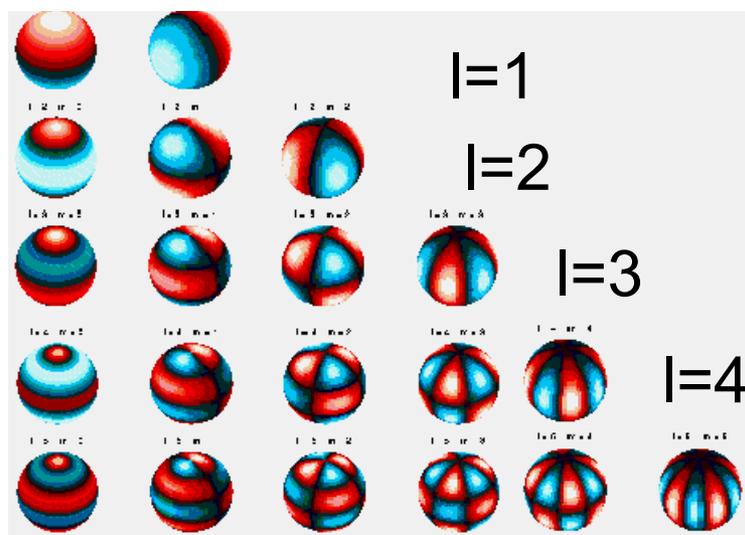
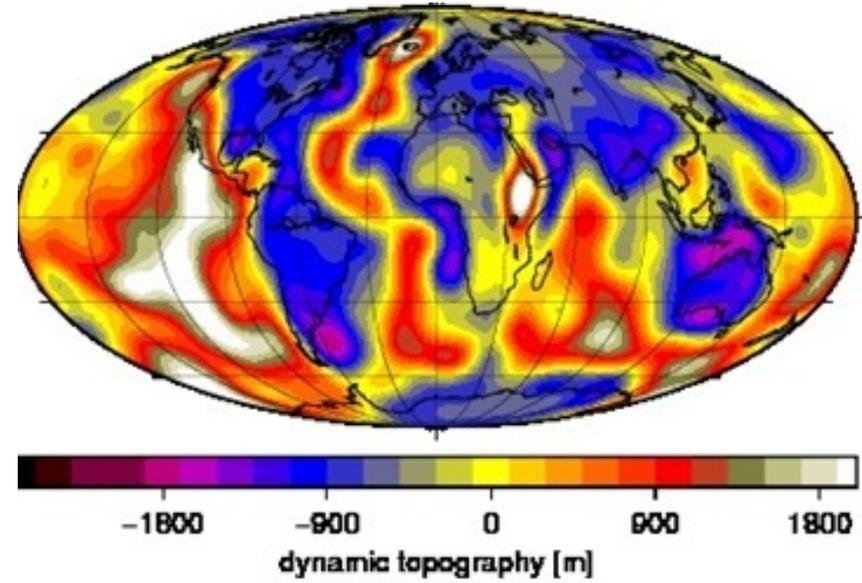
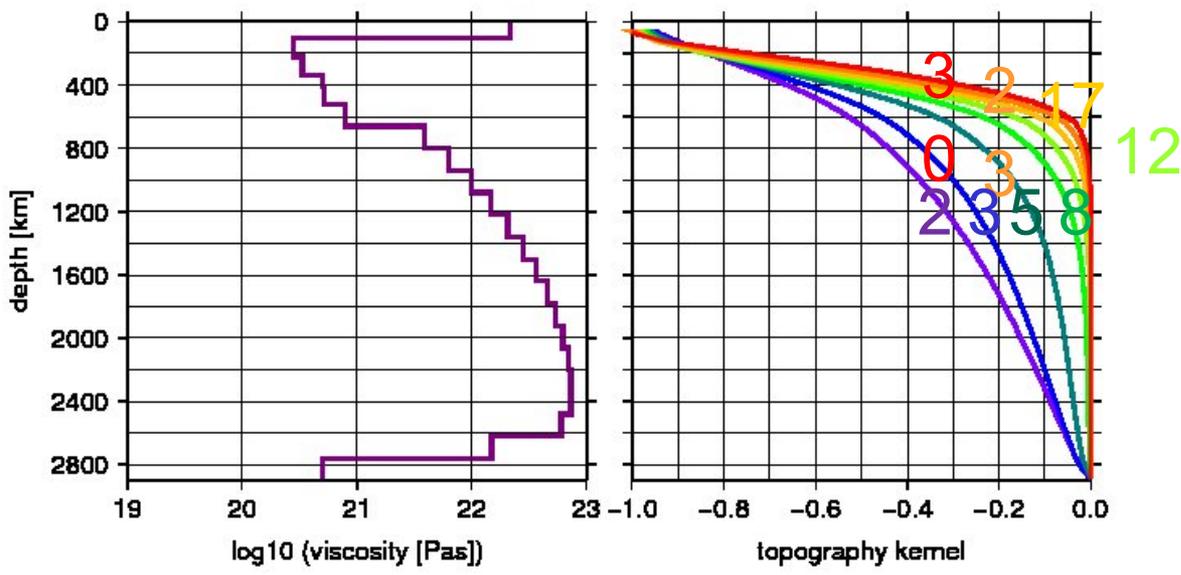
If viscosity only depends on radius:

Effect of density anomalies $\delta\rho_{lm}$ at given depth z and spherical harmonic degree l on topography can be

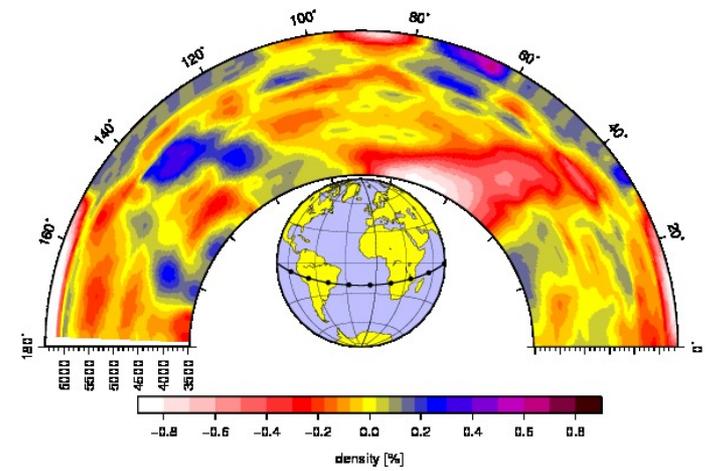
described in terms of topography kernels $K_{r,l}(z)$:

Beneath water : $\Delta\rho_s = 2280 \text{ kg/m}^3$

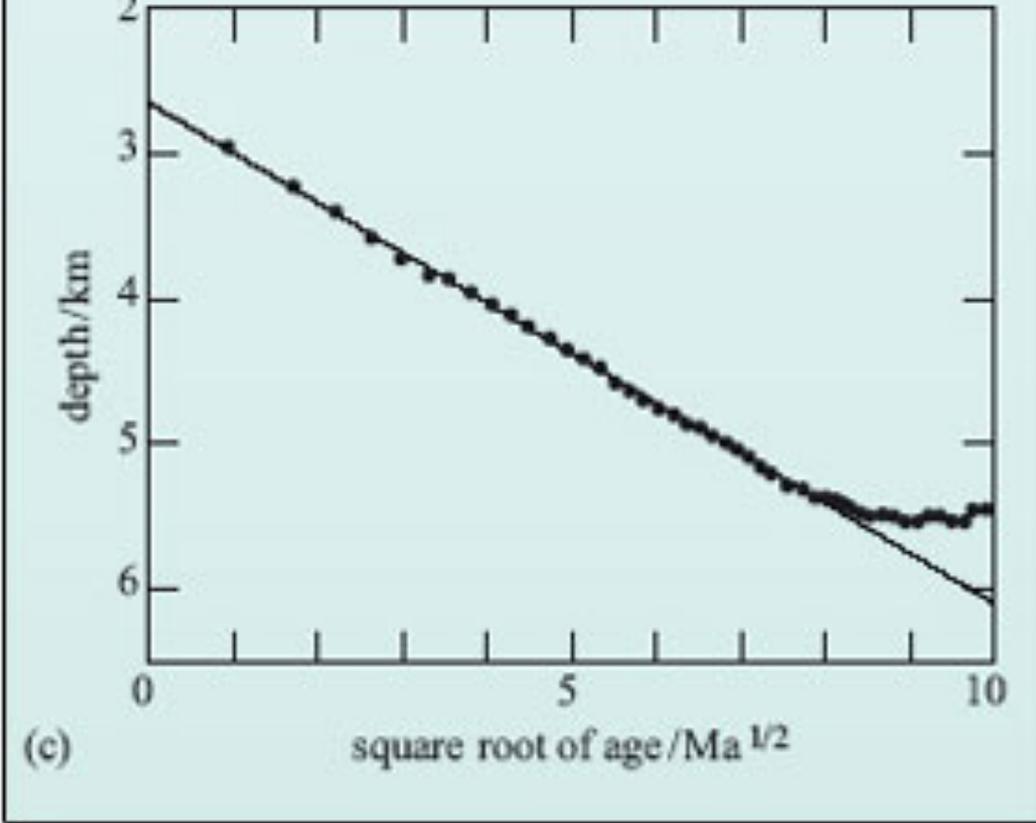
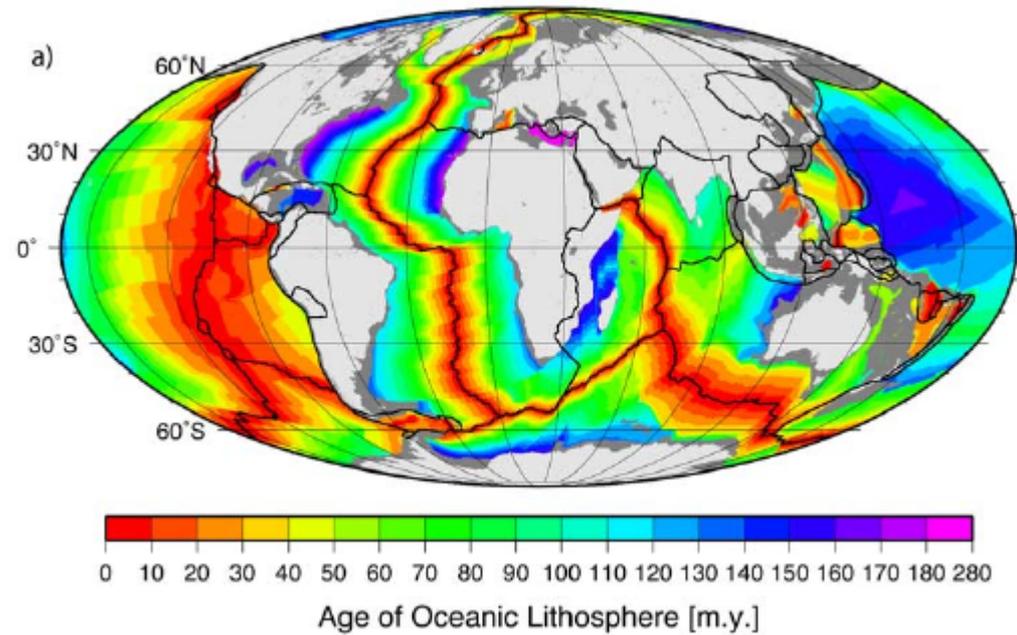
$$h_{lm} = \int \delta\rho_{lm}(z) K_{r,l}(z) dz / \Delta\rho_s$$



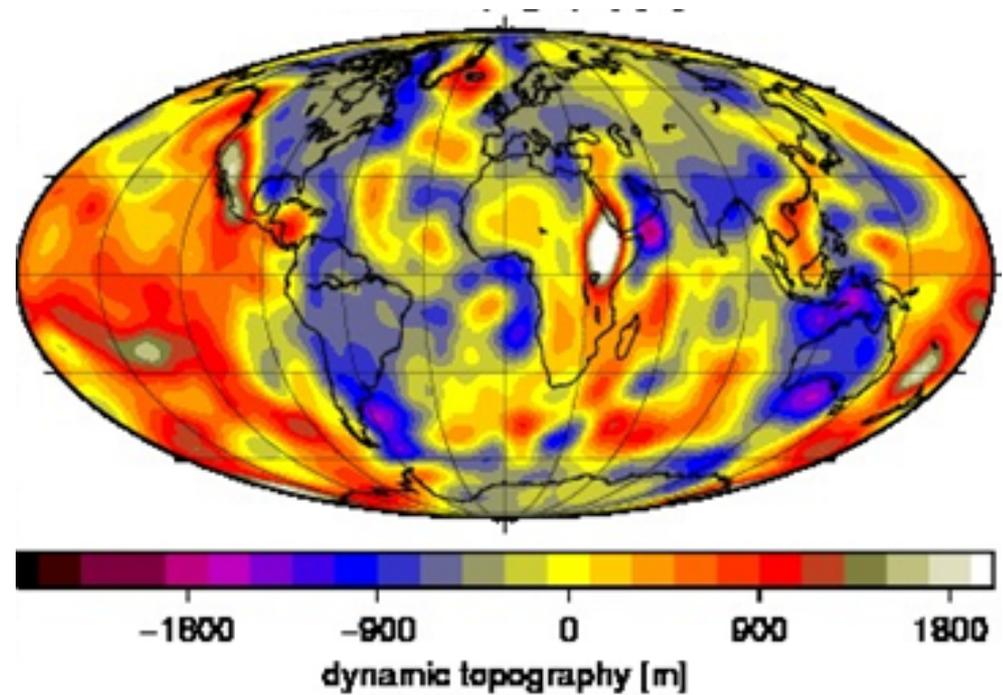
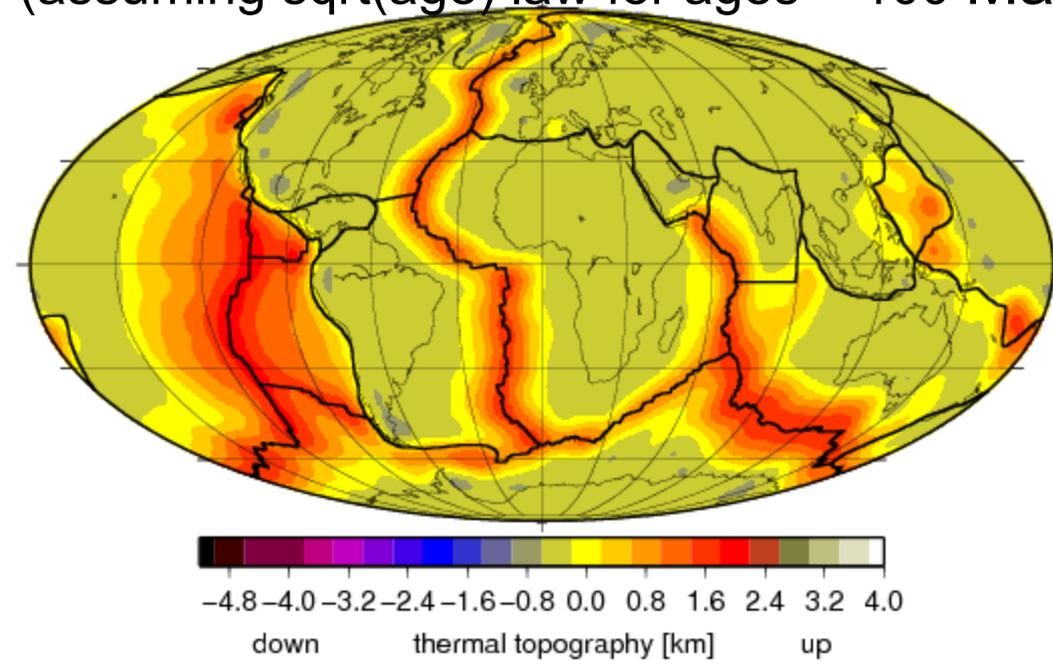
spherical harmonics



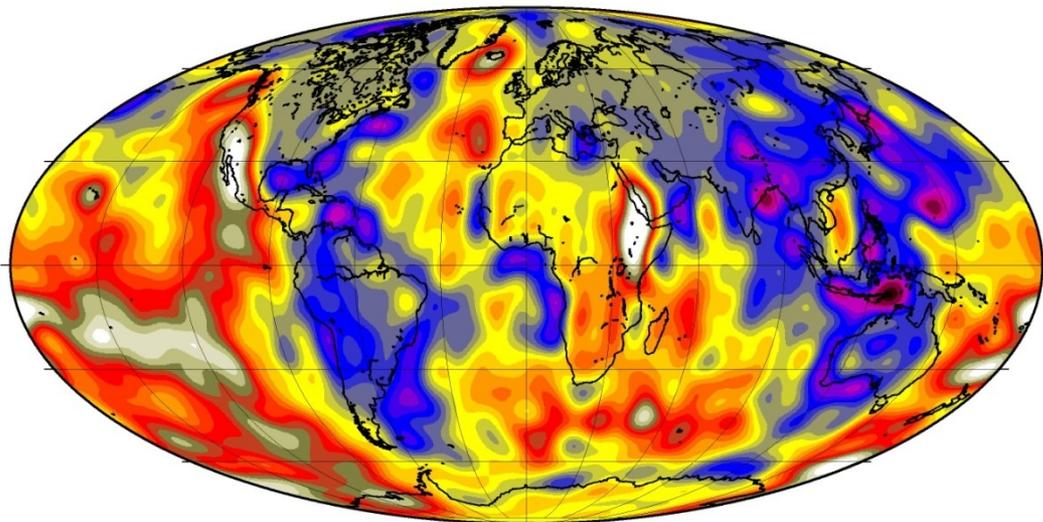
age_3.6 ocean floor age grid (Müller, Sdrolias, Gaina and Roest, G3, 2008)



Computed “ridge” topography
(assuming sqrt(age) law for ages < 100 Ma)

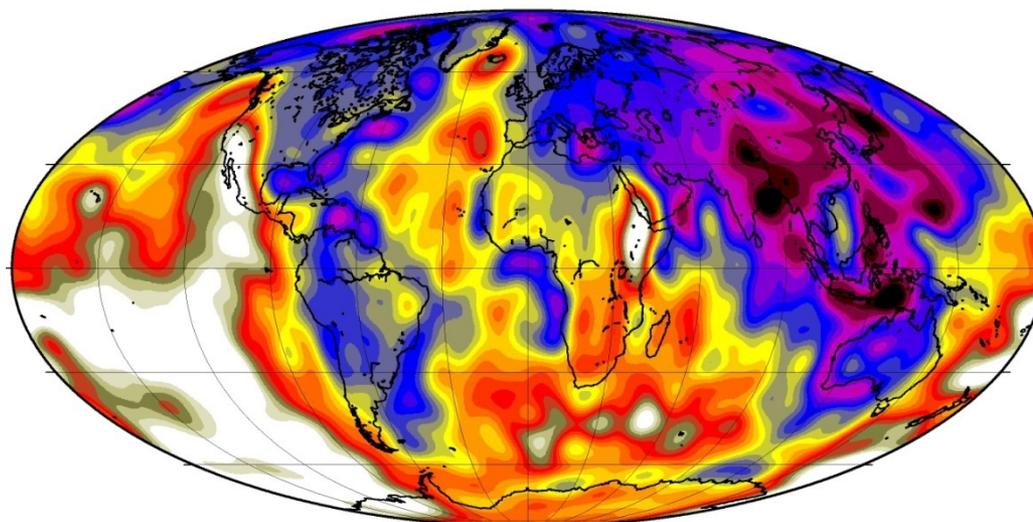


Dynamic topography – dependence on boundary condition



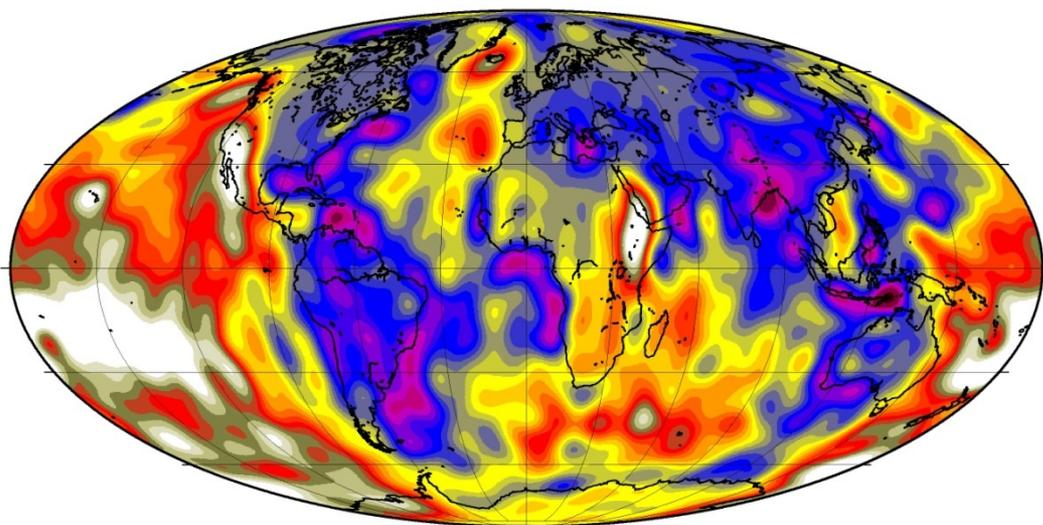
-1800 -900 0 900 1800

dynamic topography (free-slip) [m]



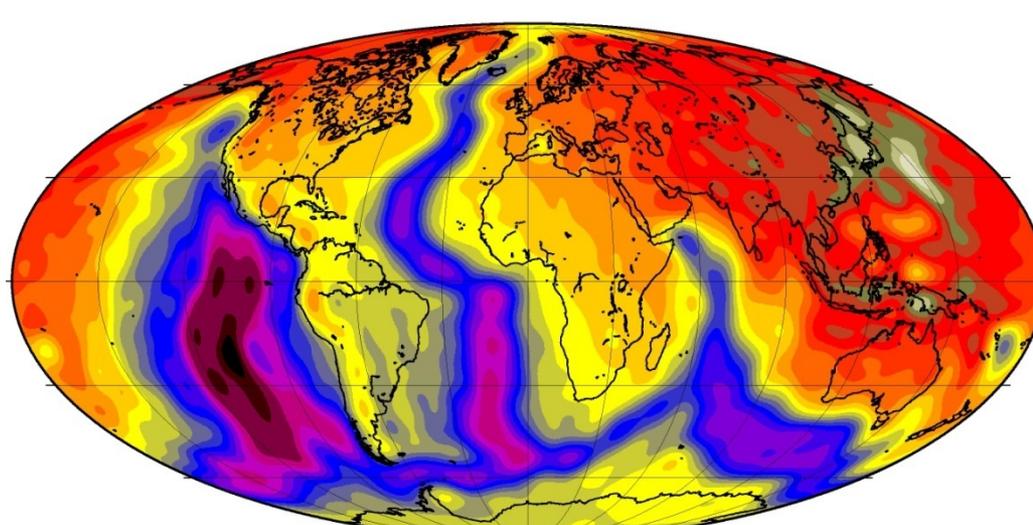
-1800 -900 0 900 1800

dynamic topography (no-slip) [m]



-1800 -900 0 900 1800

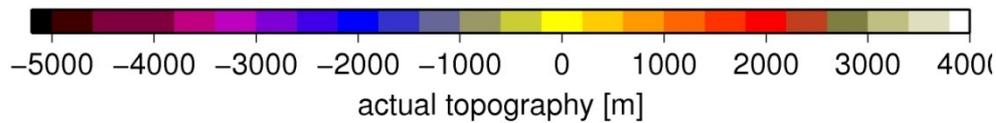
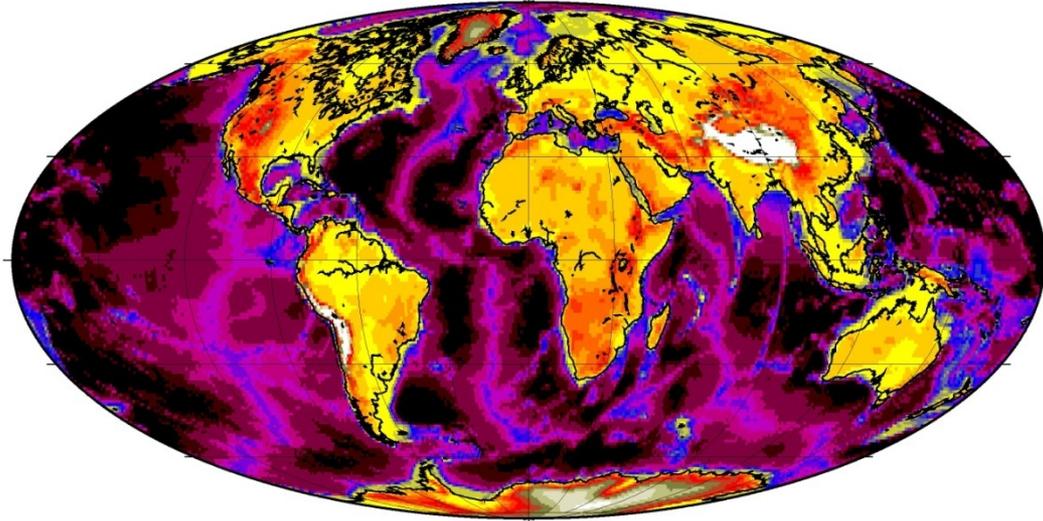
dynamic topography (prescribed plate motions) [m]



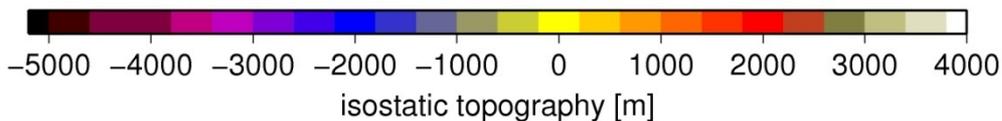
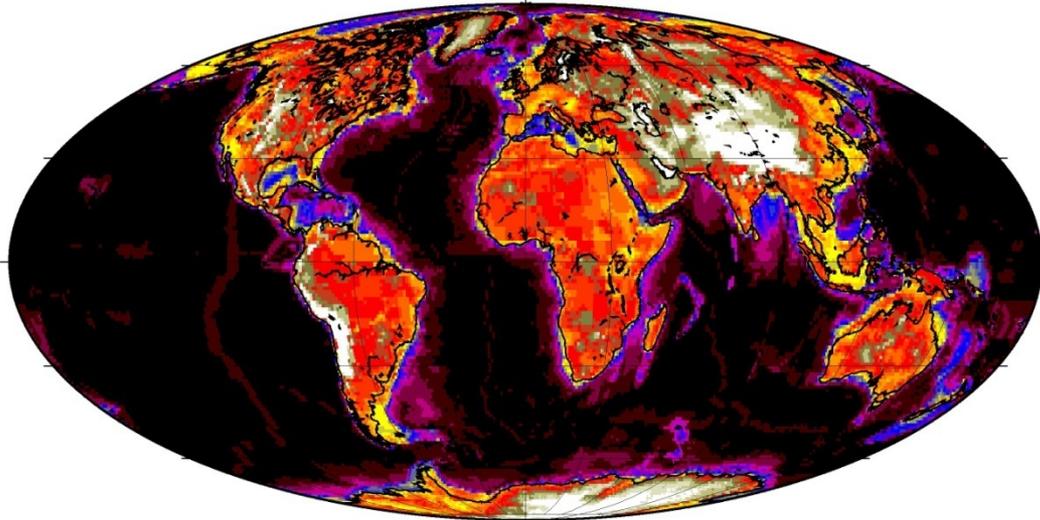
-1800 -900 0 900 1800

kinematic topography [m]

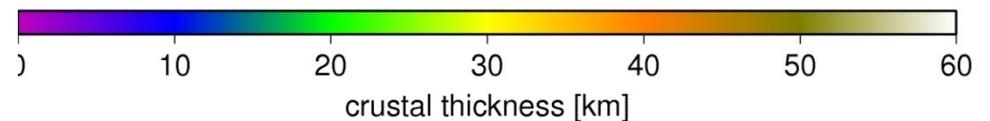
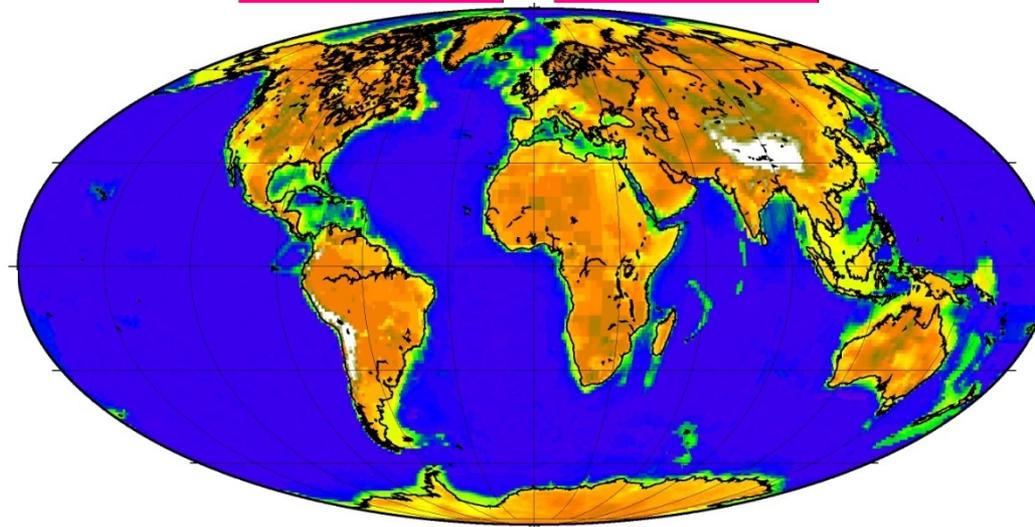
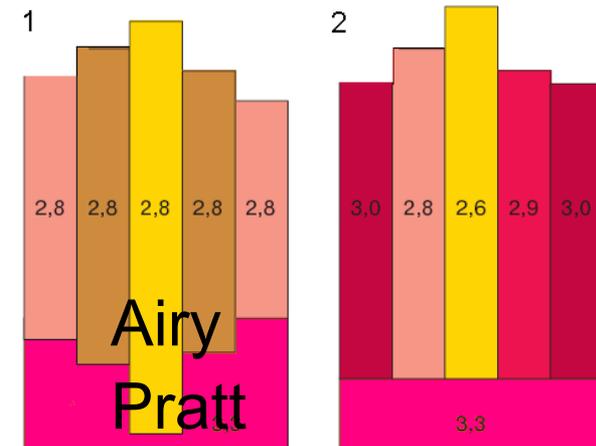
Actual topography



MINUS Isostatic topography

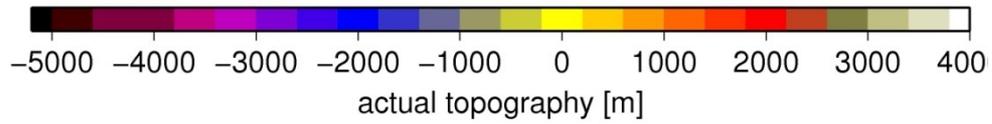
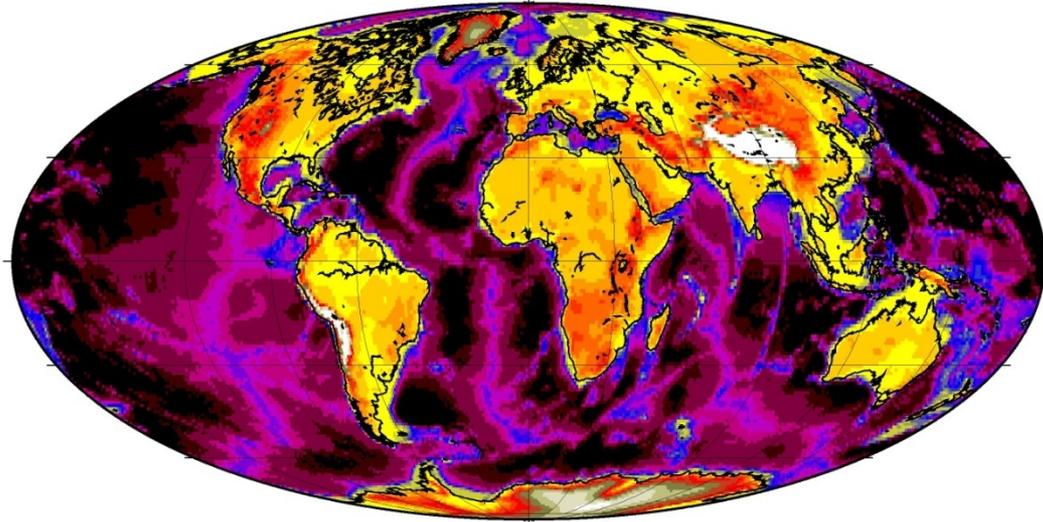


Inferring dynamic topography from observations

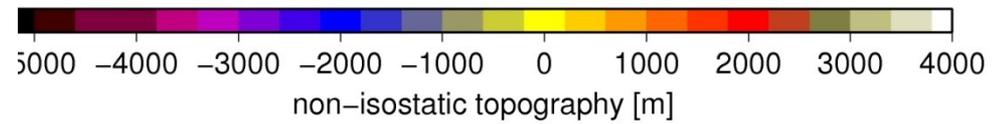
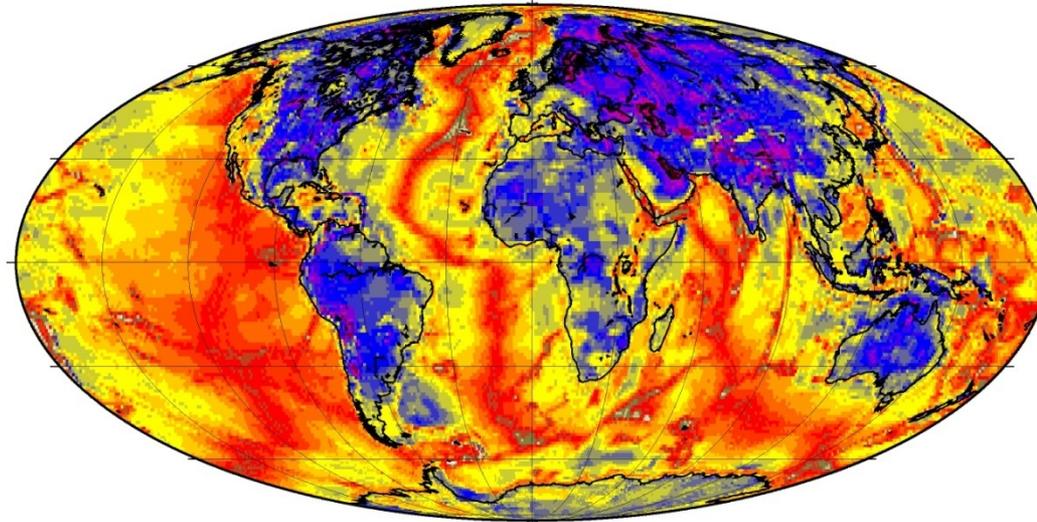


Computed based on densities and thicknesses of crustal layers in CRUST 1.0 model (Laske et al., <http://igppweb.ucsd.edu/~gabi/crust1.html>)

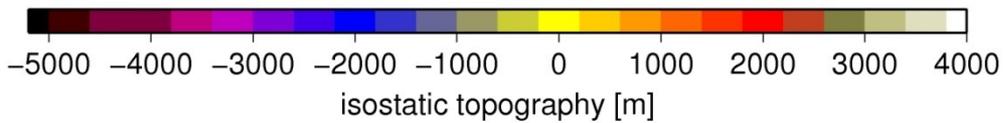
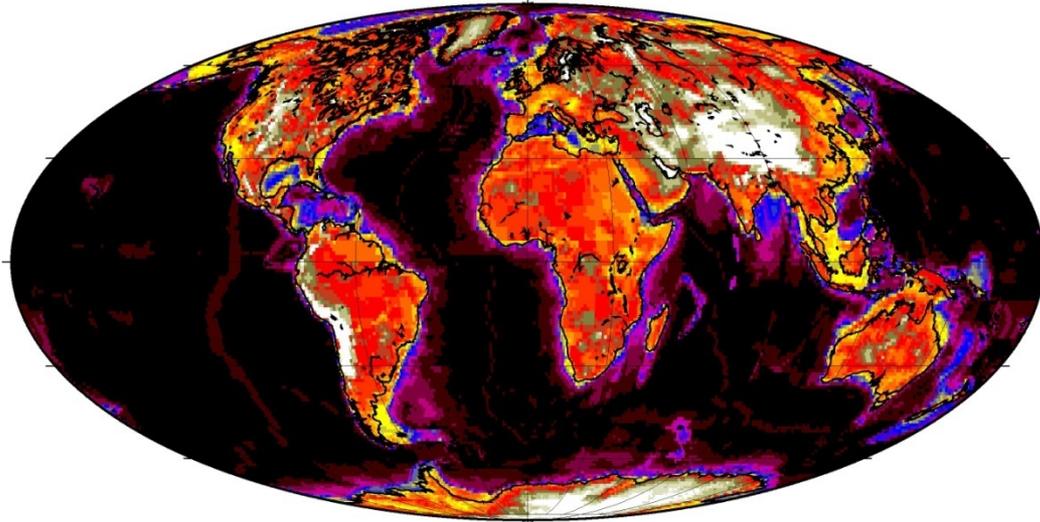
Actual



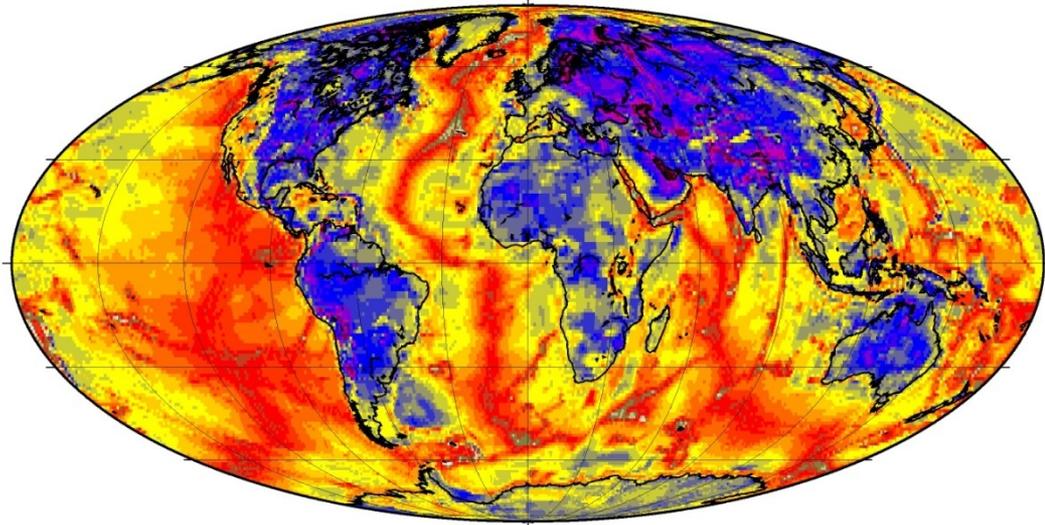
Non-isostatic topography



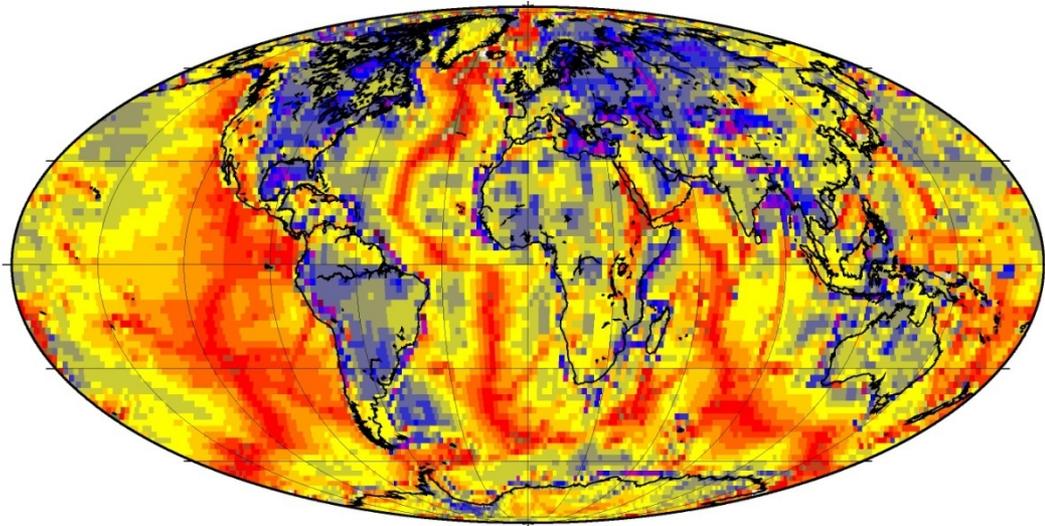
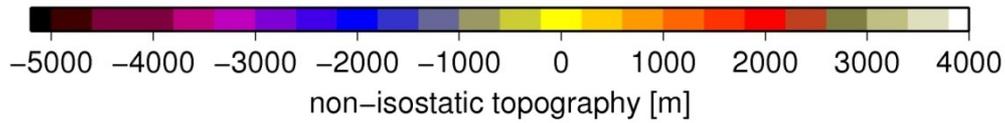
MINUS
Isostatic topography



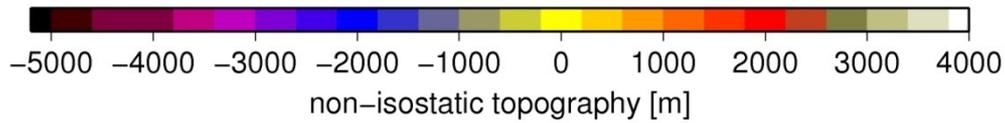
Non-isostatic



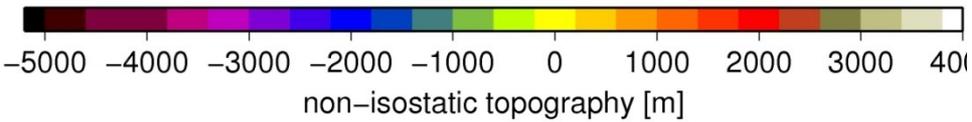
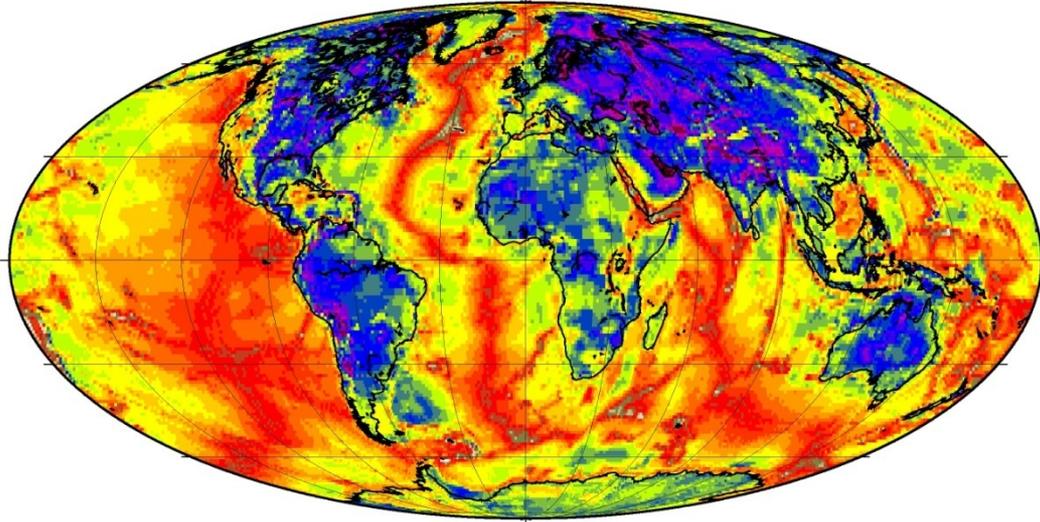
Using Crust 1.0



Using Crust 2.0



Non-isostatic topography

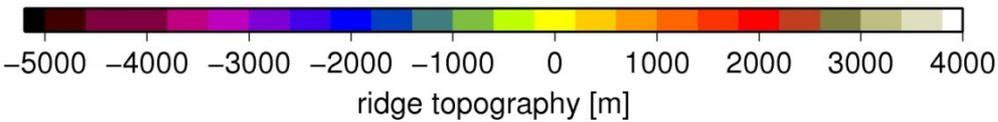
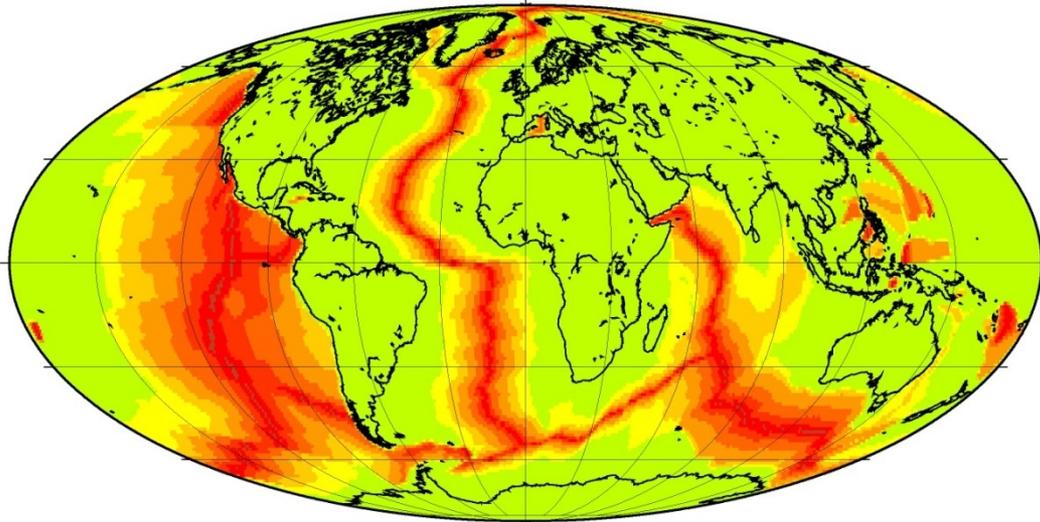
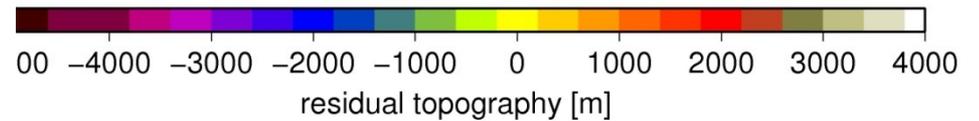
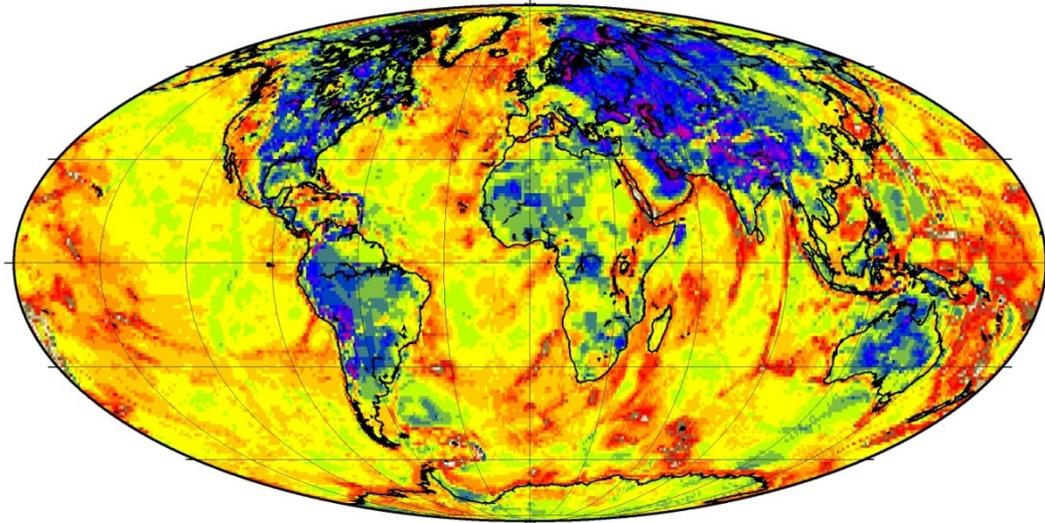


MINUS

ridge topography

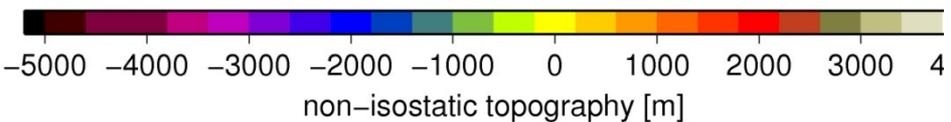
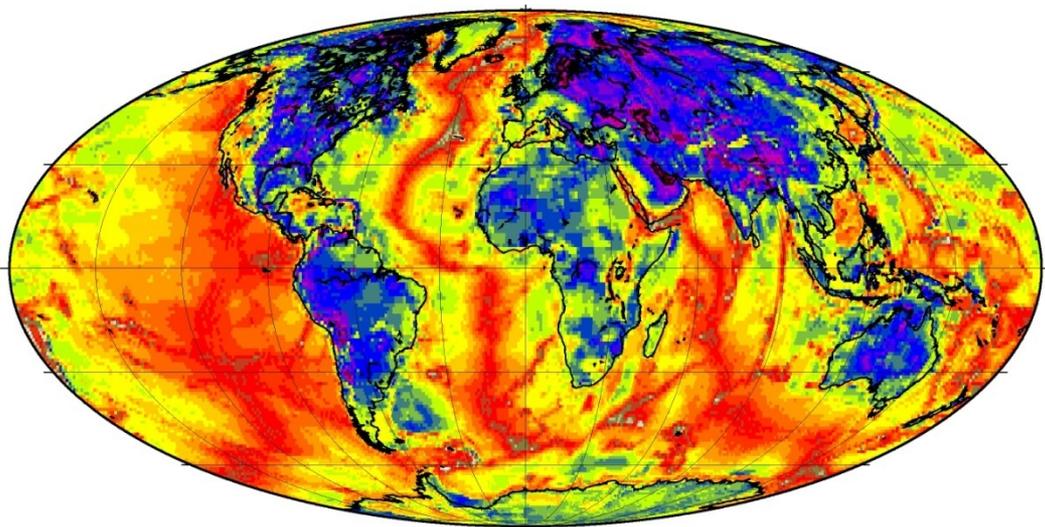


residual topography

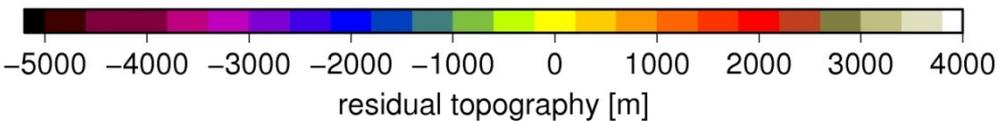
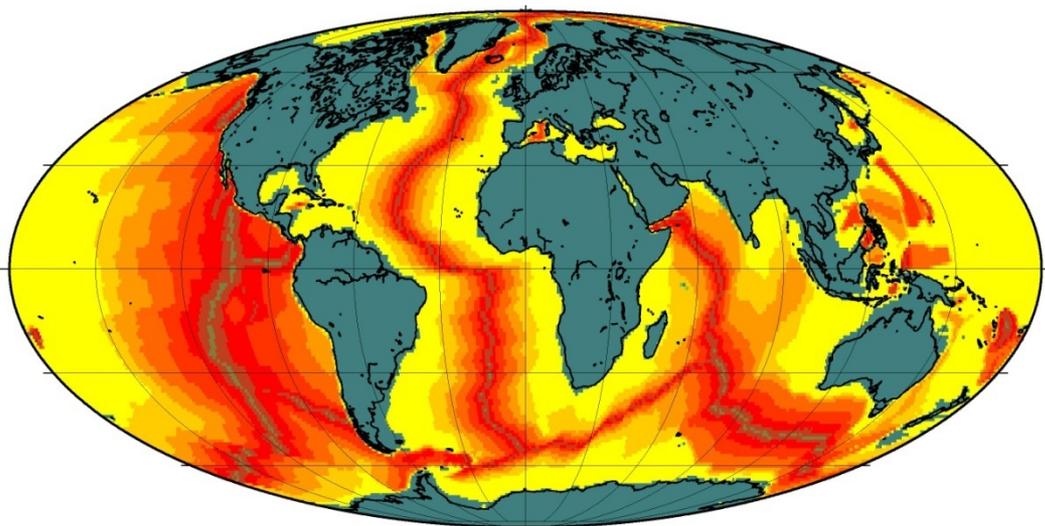


Continents like 100 Ma ocean floor

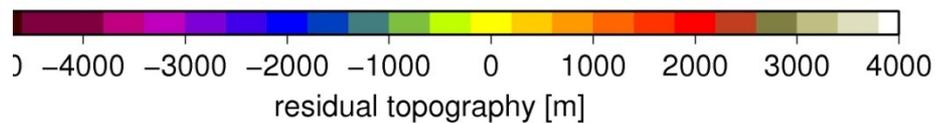
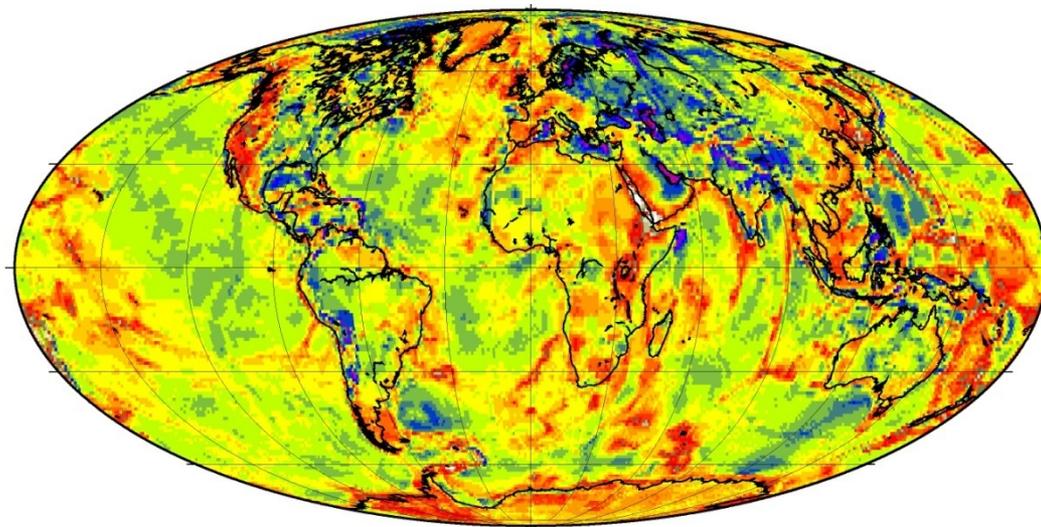
Non-isostatic topography



MINUS
ridge topography

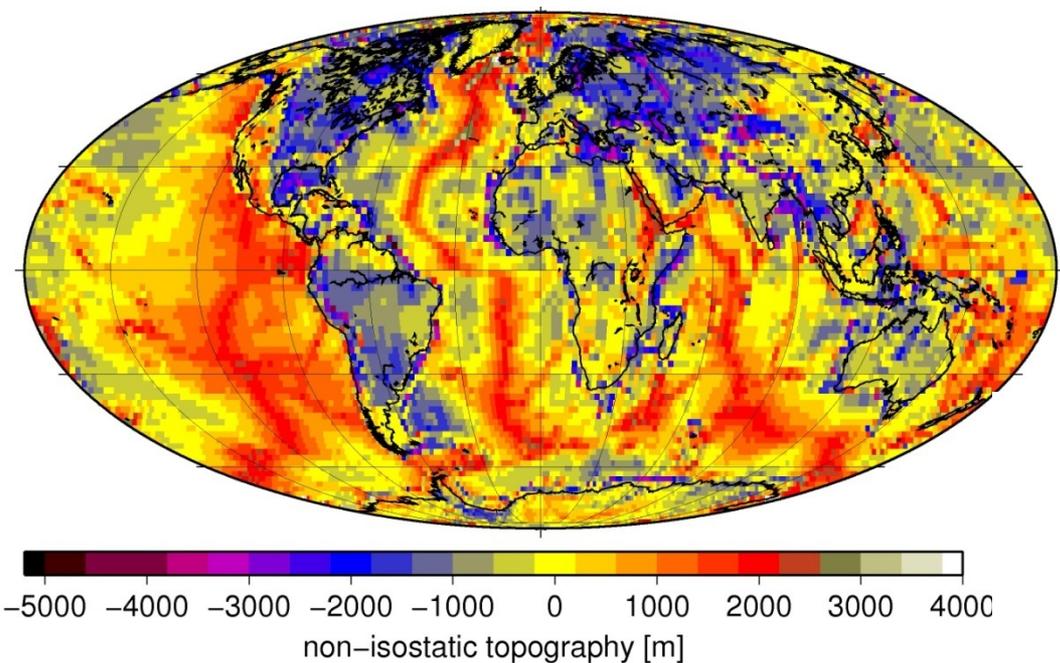


residual topography

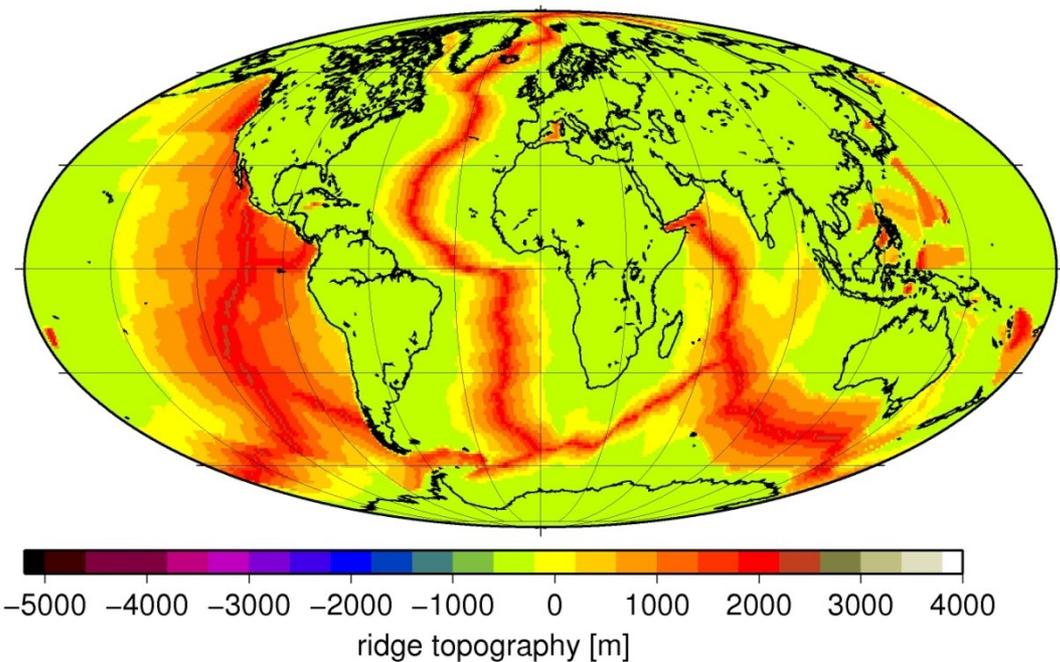


Continents sqrt (200 Ma) topography

Non-isostatic topography

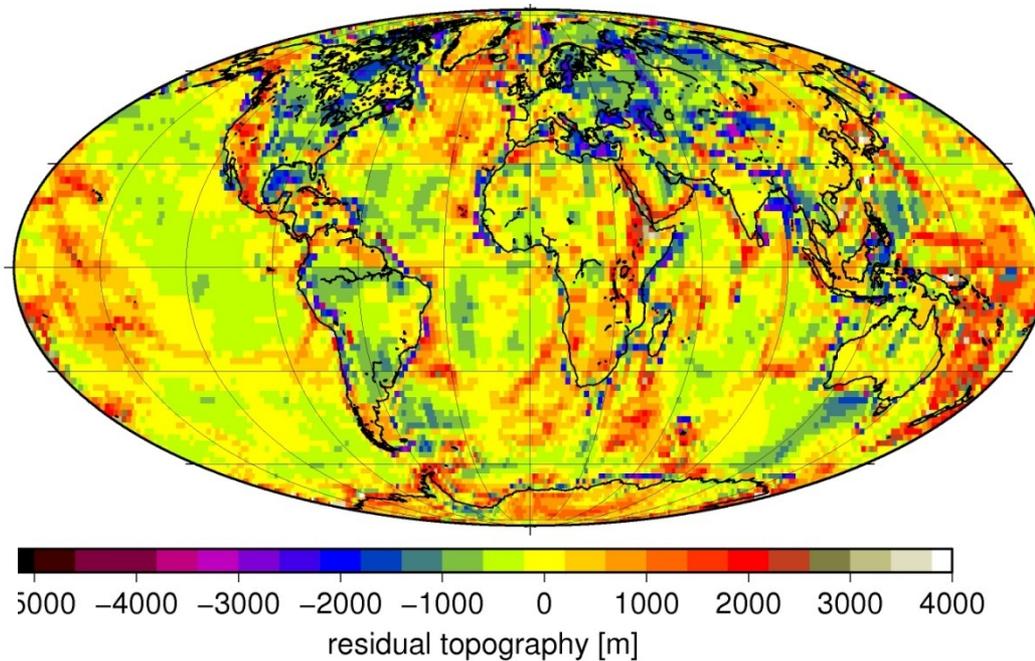


MINUS
ridge topography



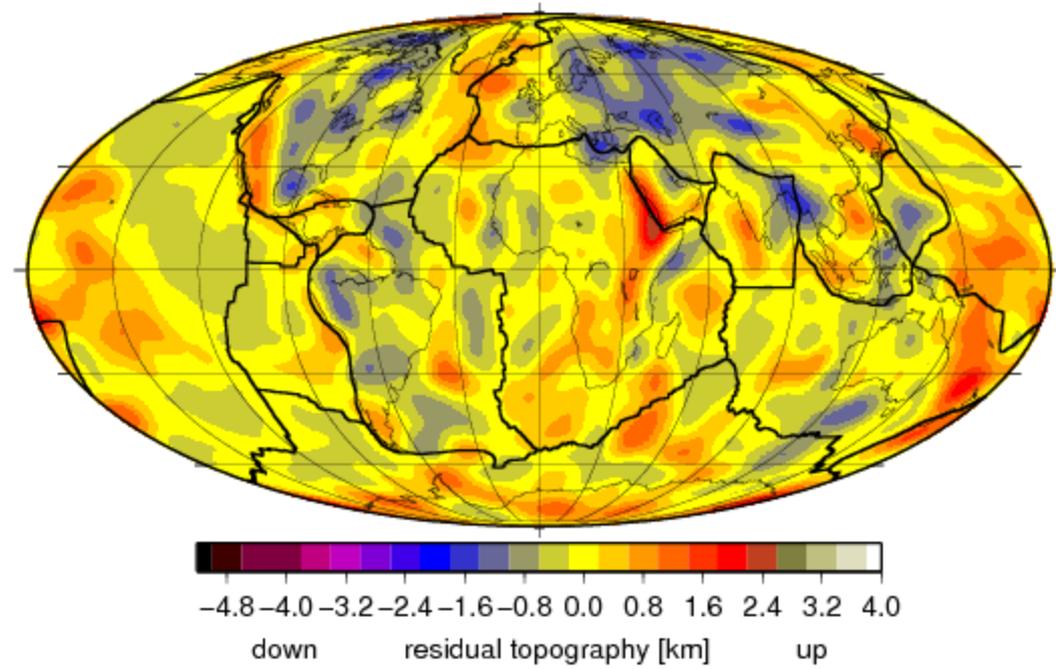
Using Crust 2.0

residual topography



Continents like 100 Ma ocean floor

residual topography, $l=1-31$

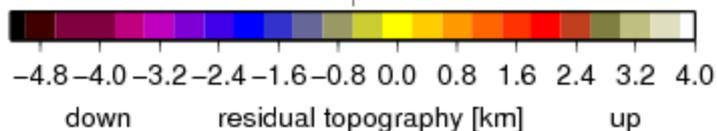
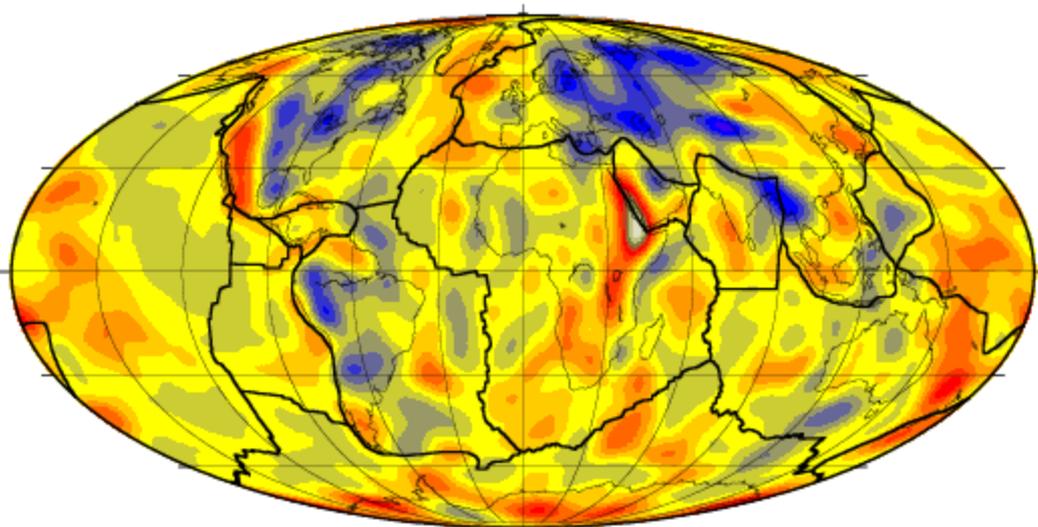
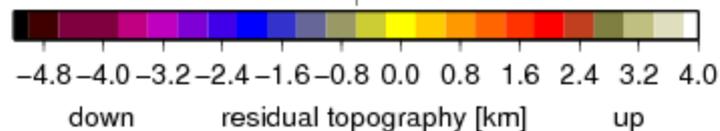
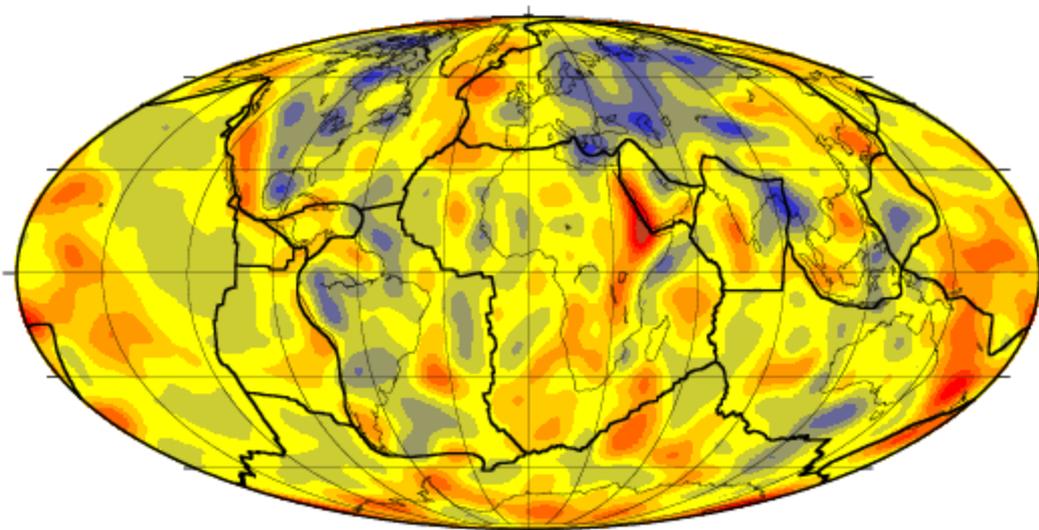


Residual topography is the observation-based quantity to which dynamic topography computed from mantle flow can be compared

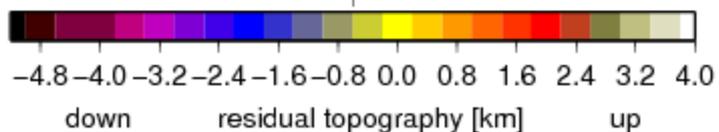
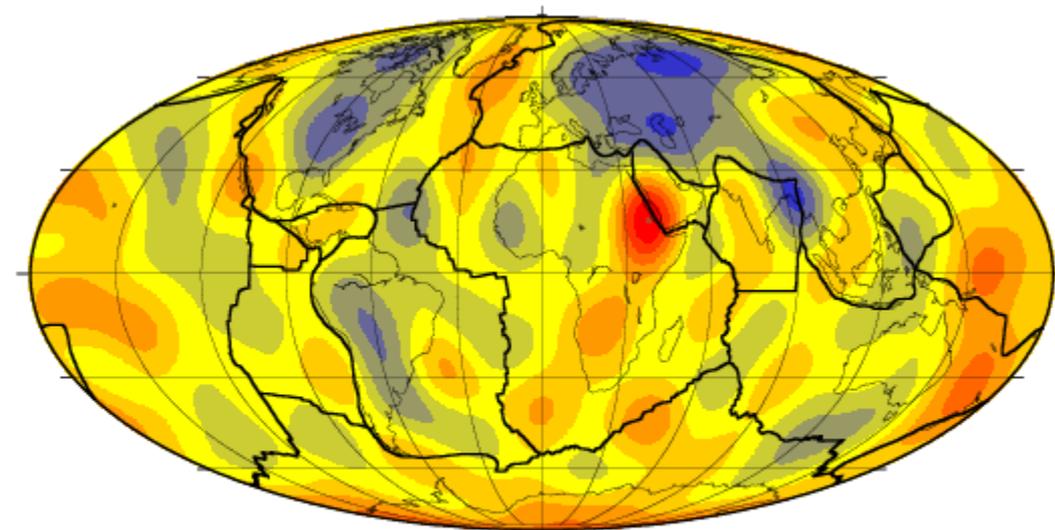
residual topography, $l=1-31$

Values above sea level multiplied with factor 1.45, because dynamic topography is computed for global seawater coverage

residual topography, $l=1-31$

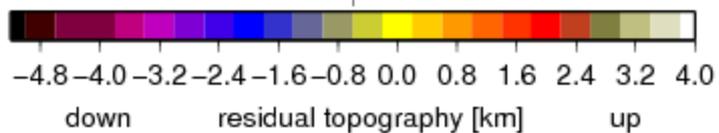
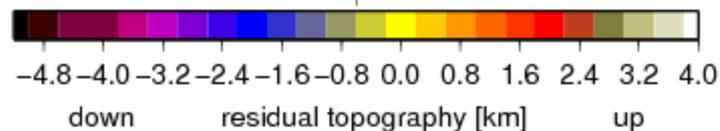
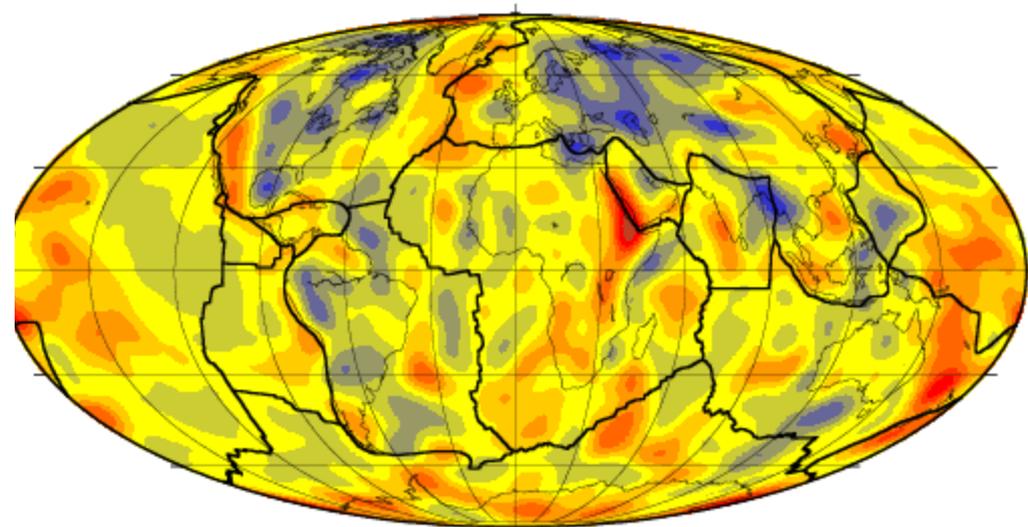


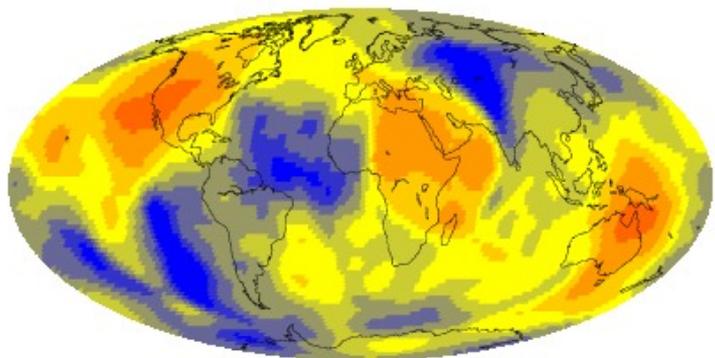
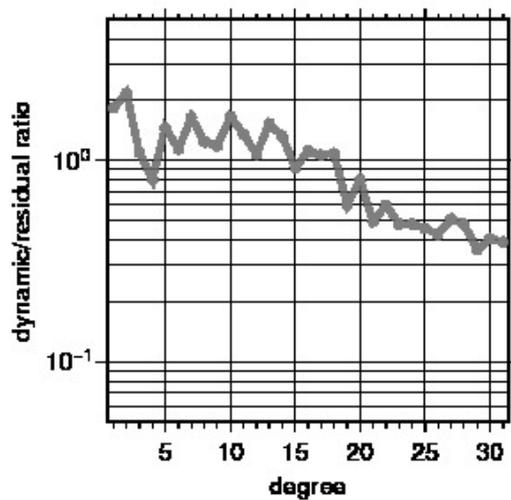
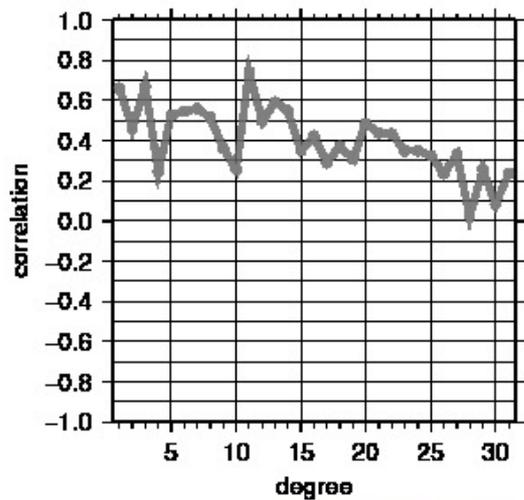
residual topography, $l=1-12$, above sea level multiplied with 1.45



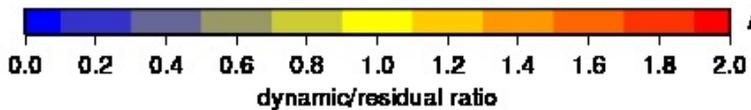
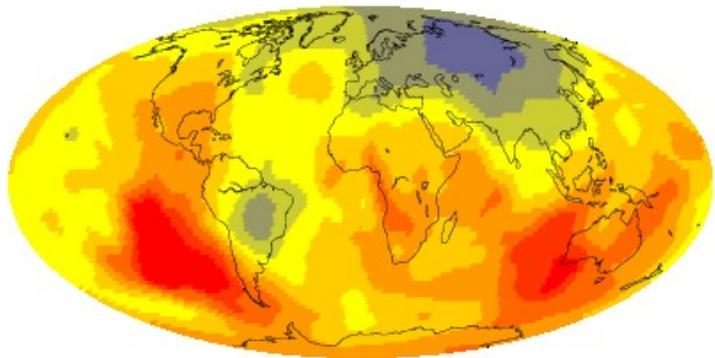
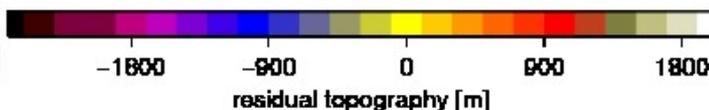
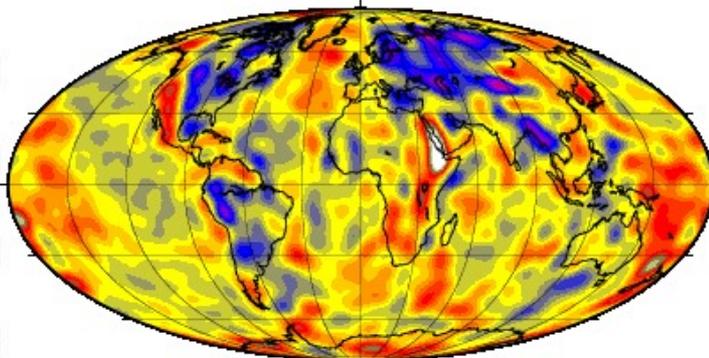
residual topography, $l=1-31$
Above sea level multiplied with 1.45

residual topography, $l=1-31$

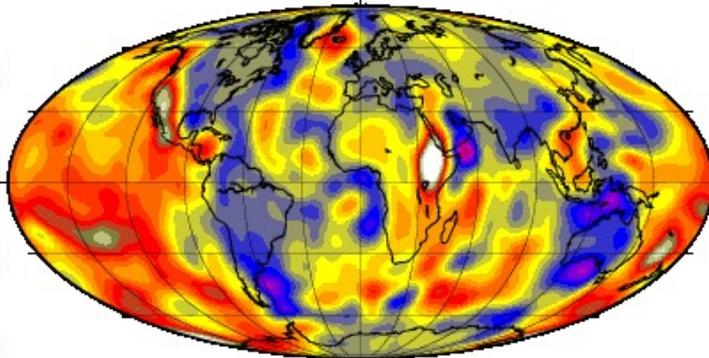




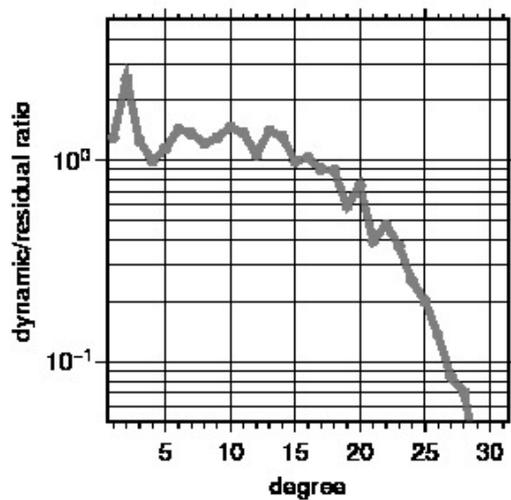
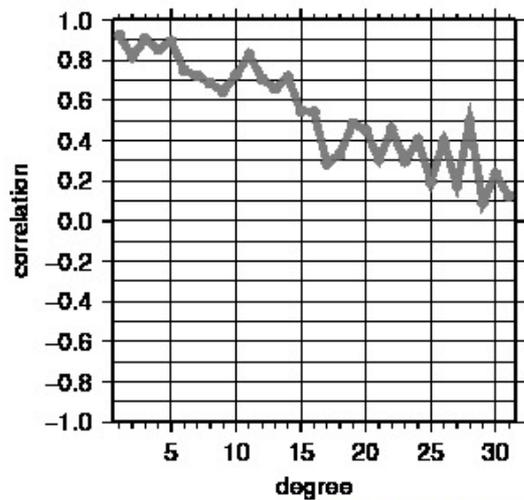
global
 .460 (.510)
 oceans only
 .390 (.429)
 continents only
 .534 (.577)
 AF plate only
 .605 (.683)
 AF continent only
 .694 (.797)



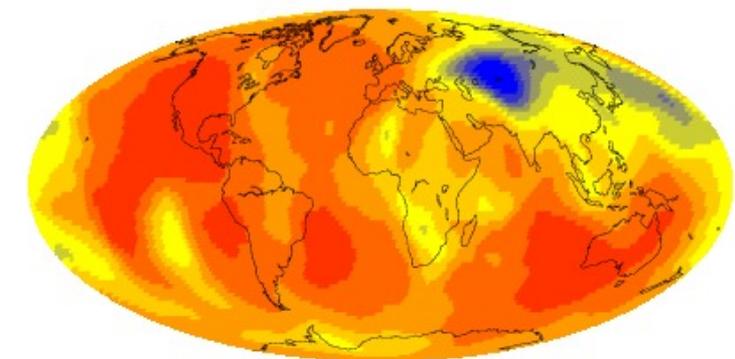
global
 1.272 (1.463)
 oceans only
 1.573 (1.821)
 continents only
 .957 (1.066)
 AF plate only
 1.349 (1.421)
 AF continent only
 1.199 (1.301)
 1<l<31 (1<l<12)



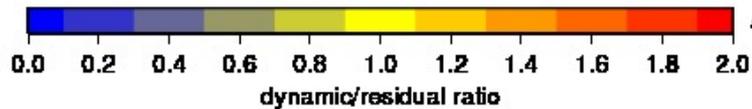
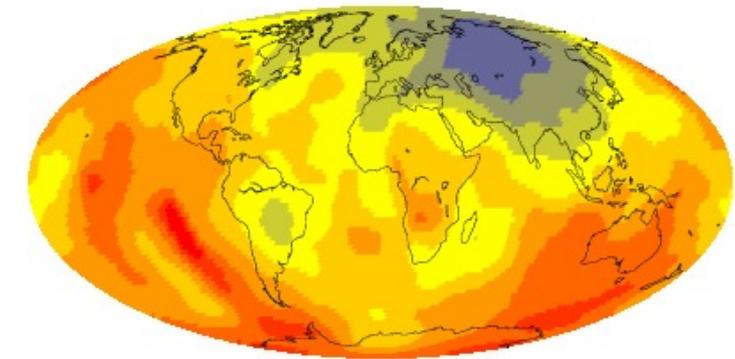
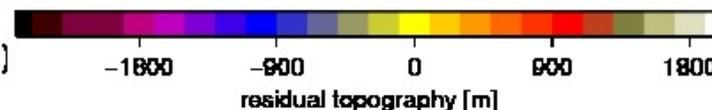
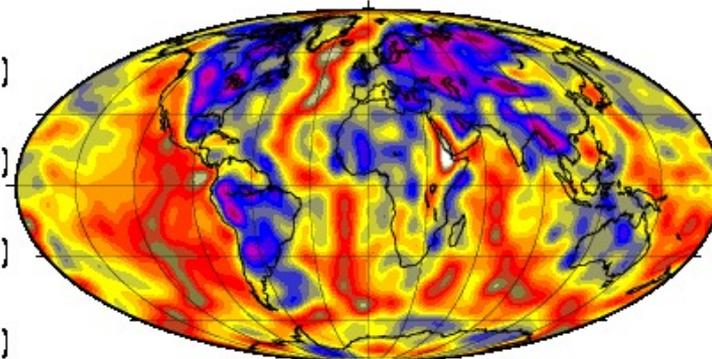
TX2007



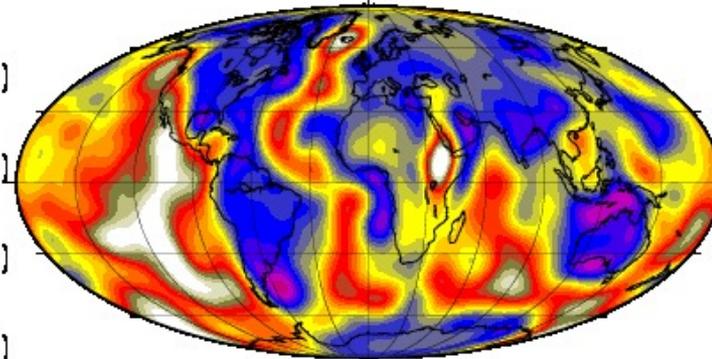
Results without
subtracting ridge
topography



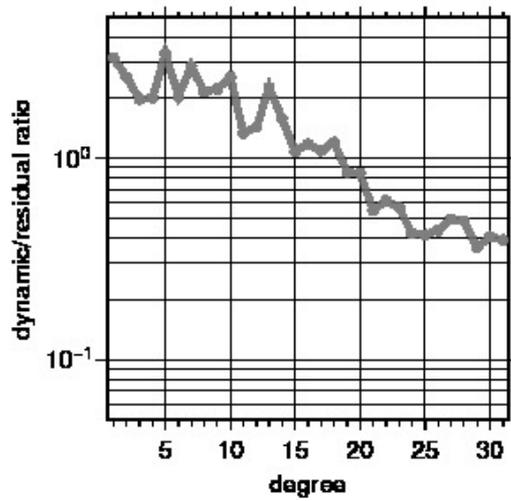
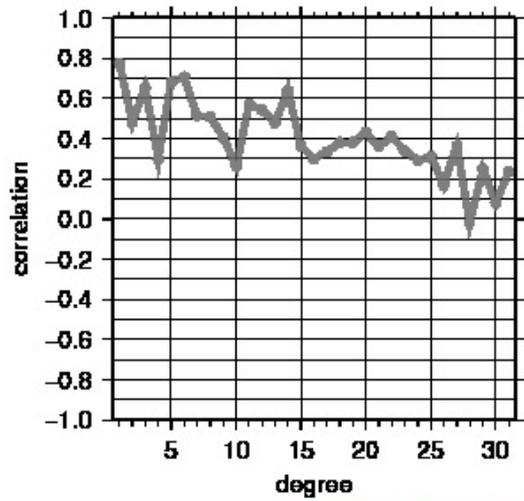
global
.754 (.816)
oceans only
.748 (.817)
continents only
.574 (.657)
AF plate only
.705 (.774)
AF continent only
.686 (.798)



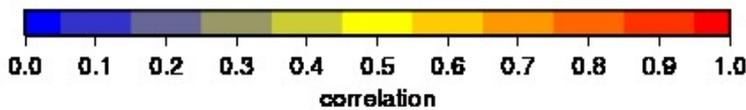
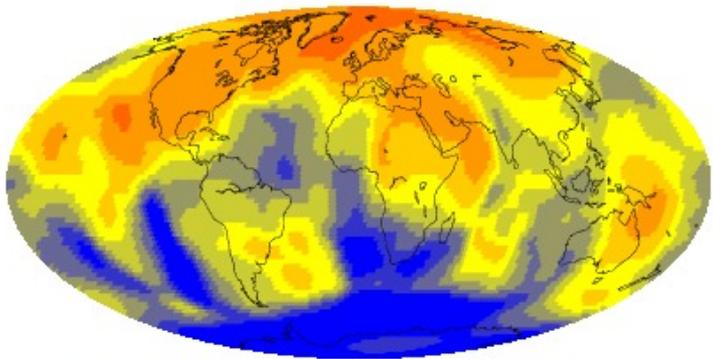
global
1.245 (1.330)
oceans only
1.519 (1.670)
continents only
.950 (1.040)
AF plate only
1.135 (1.130)
AF continent only
1.213 (1.292)
1<l<31 (1<l<12)



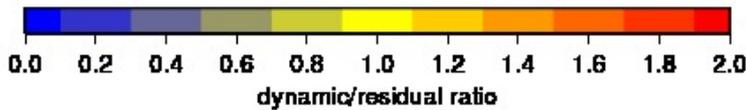
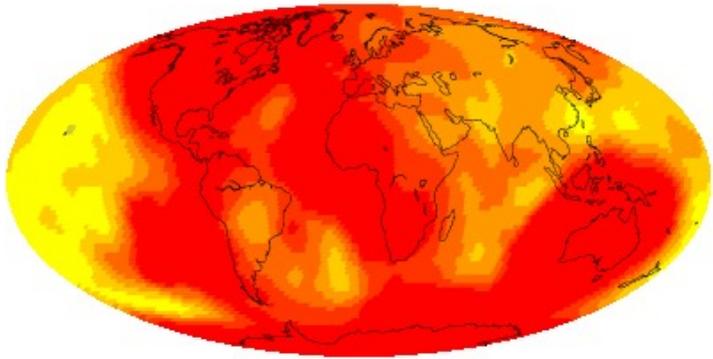
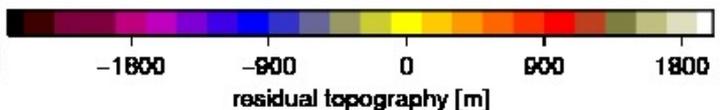
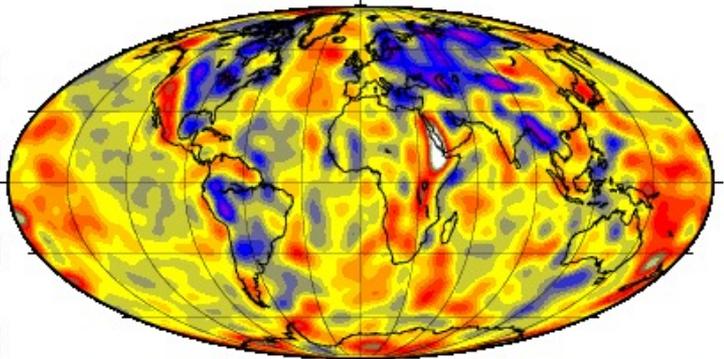
TX2007



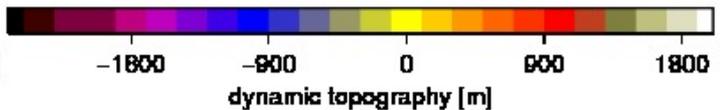
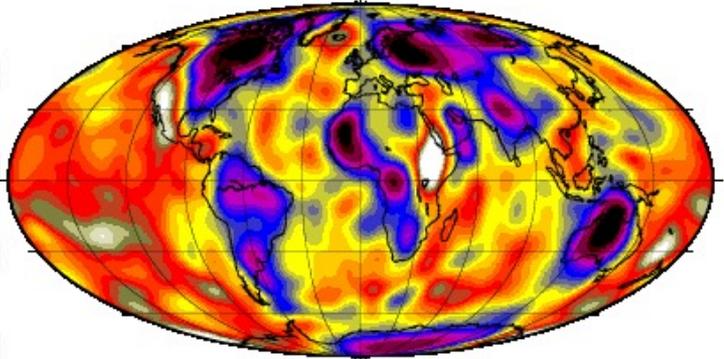
Results with thermal scaling from s-wavespeed to density also within continental lithosphere



global	.489	(.563)
oceans only	.391	(.434)
continents only	.556	(.624)
AF plate only	.485	(.504)
AF continent only	.615	(.659)

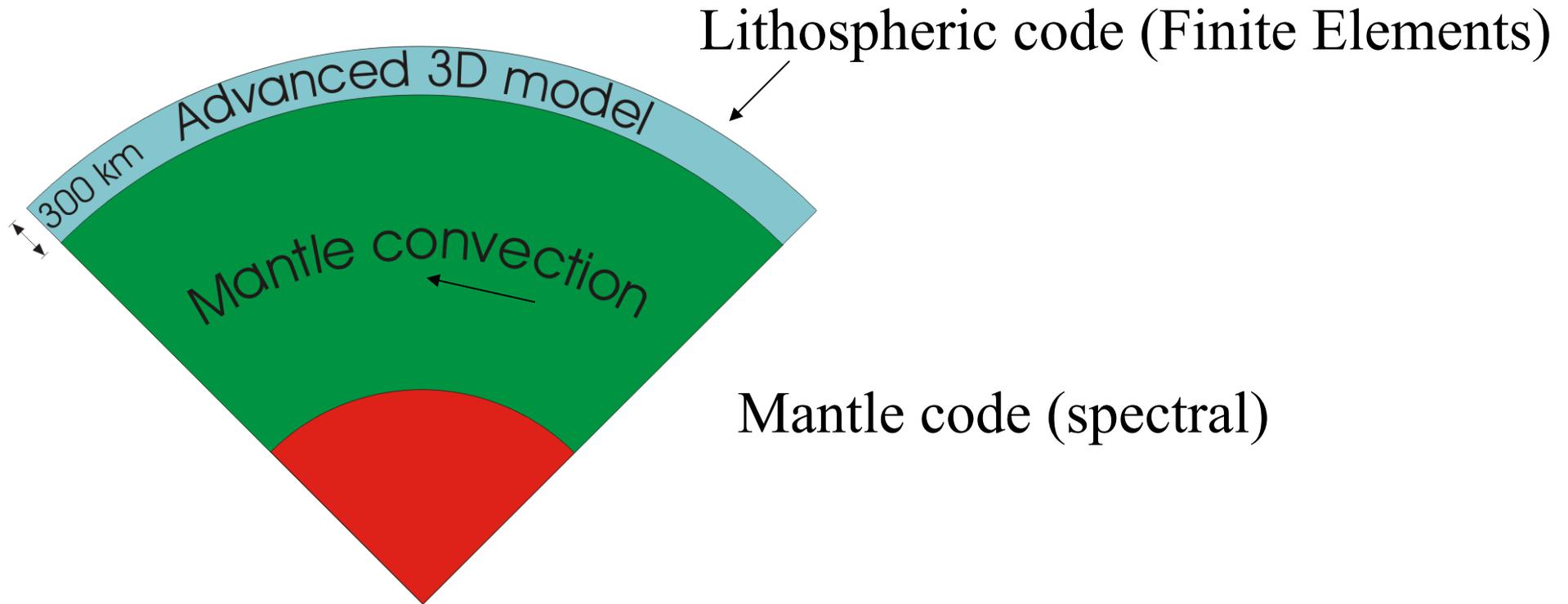


global	2.016	(2.417)
oceans only	1.628	(2.022)
continents only	1.816	(2.113)
AF plate only	2.104	(2.481)
AF continent only	2.187	(2.491)
1<k<31	1<k<31	(1<k<12)



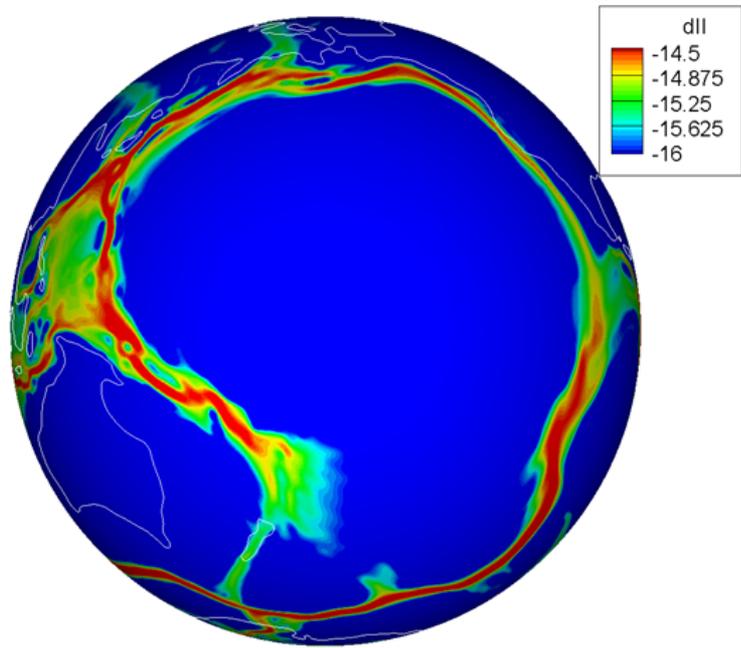
TX2007

With self-generated plate boundaries

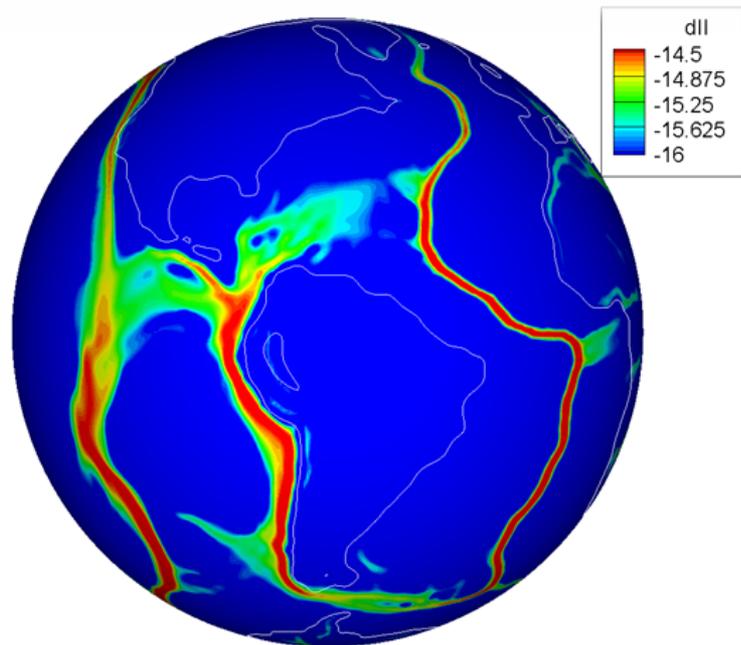
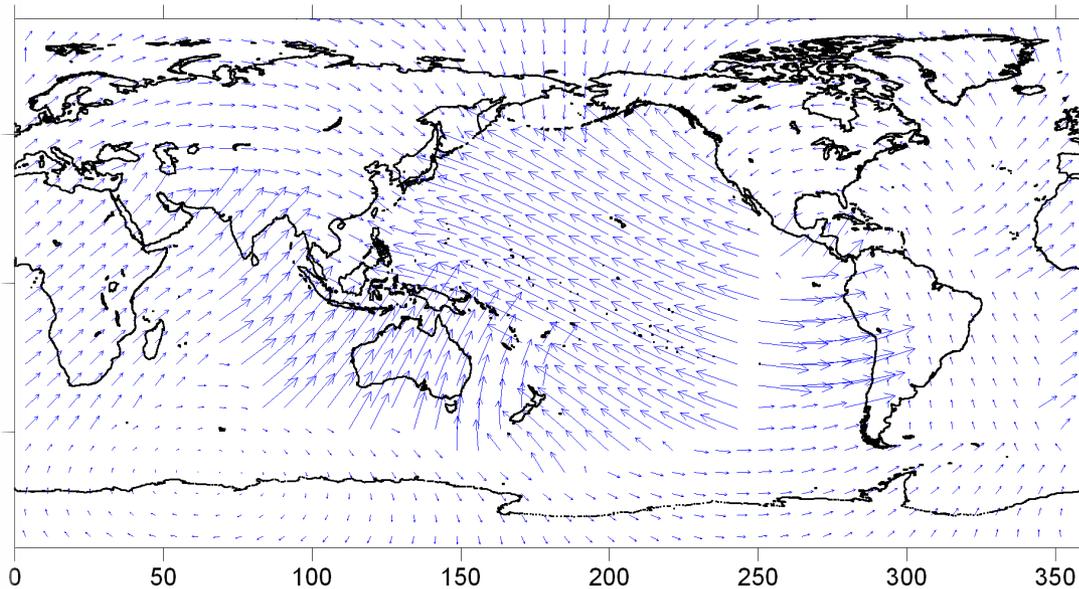


Mantle and lithospheric codes are coupled through continuity of velocities and tractions at 300 km.

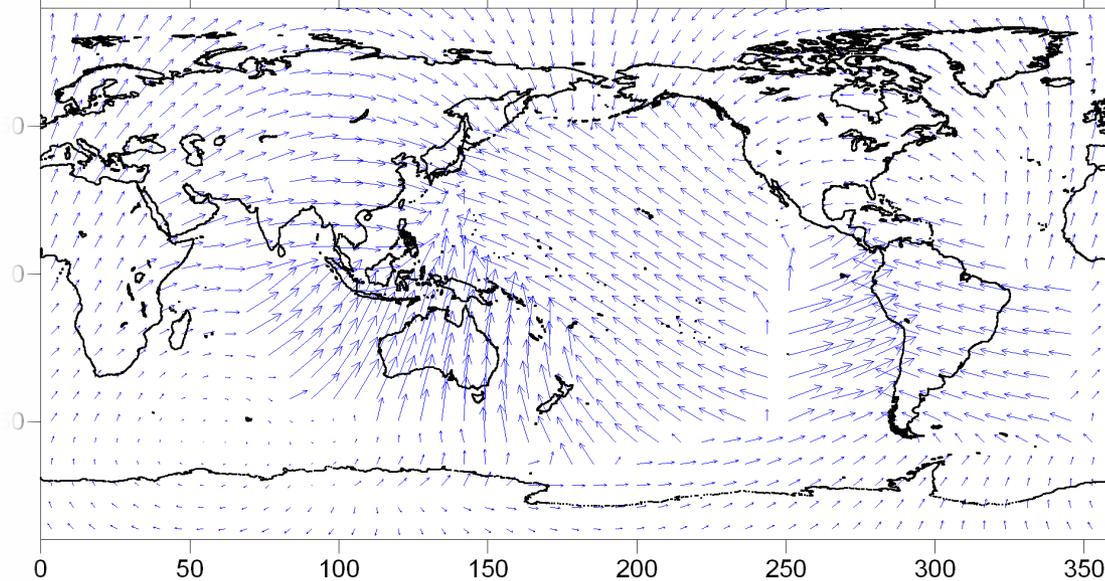
Self-generated plate boundaries



Observed

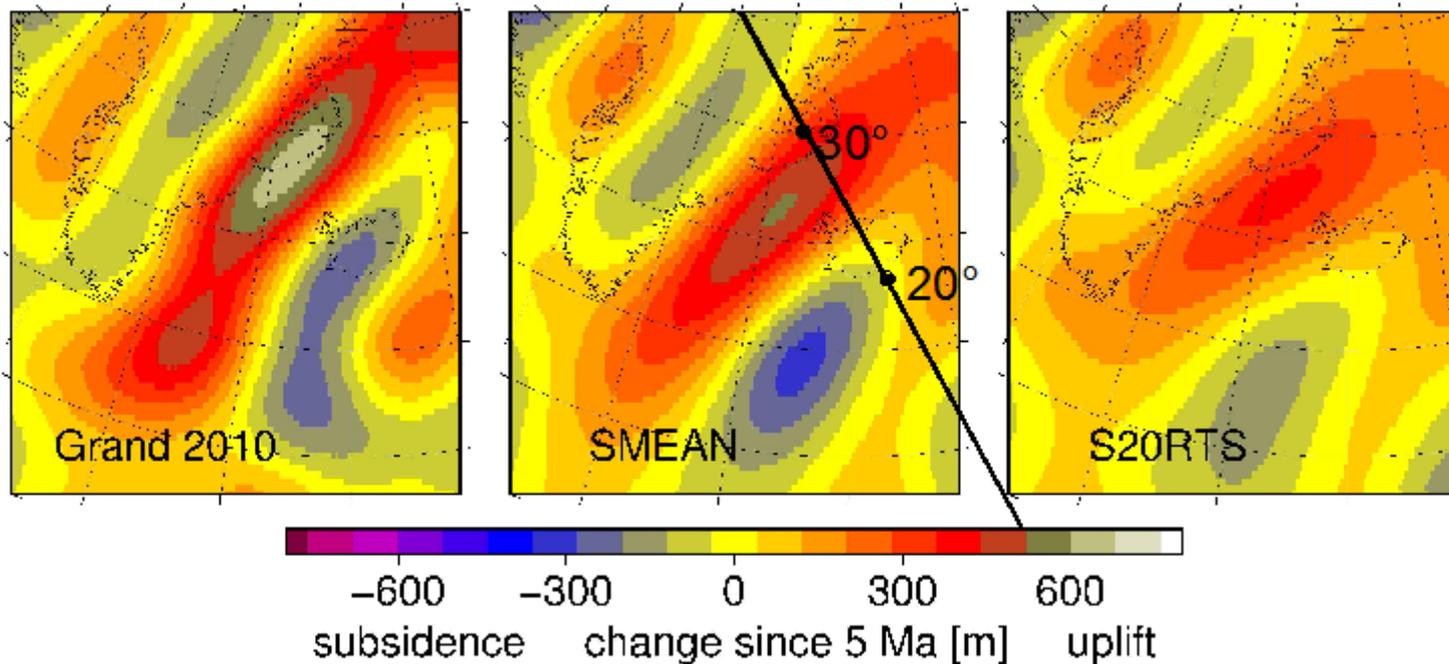
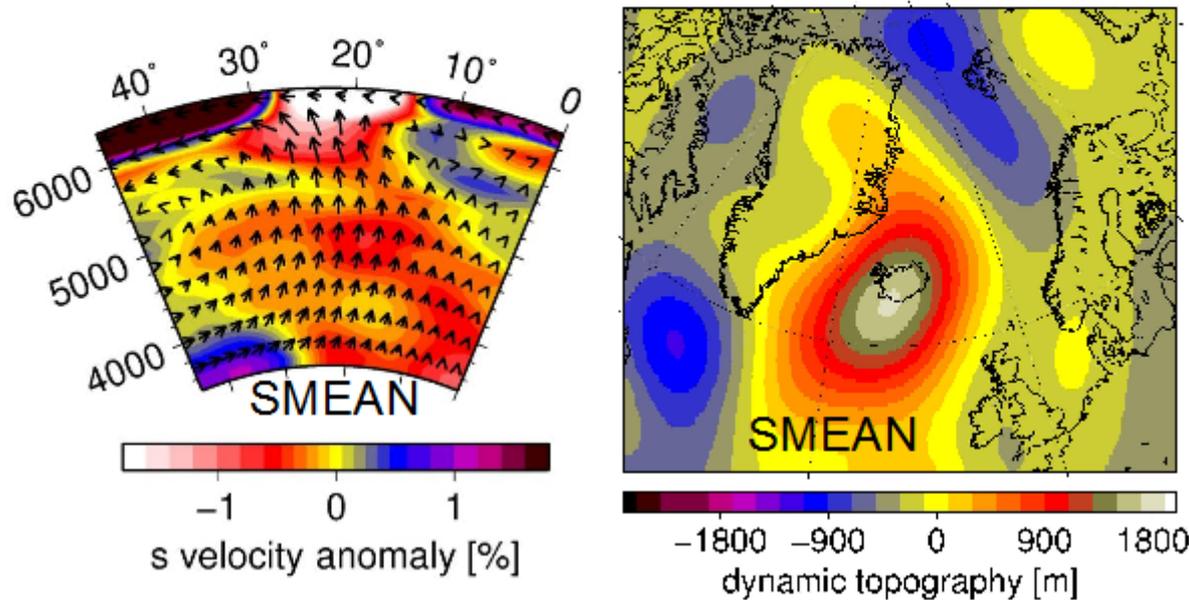


Model



Modelling uplift rates

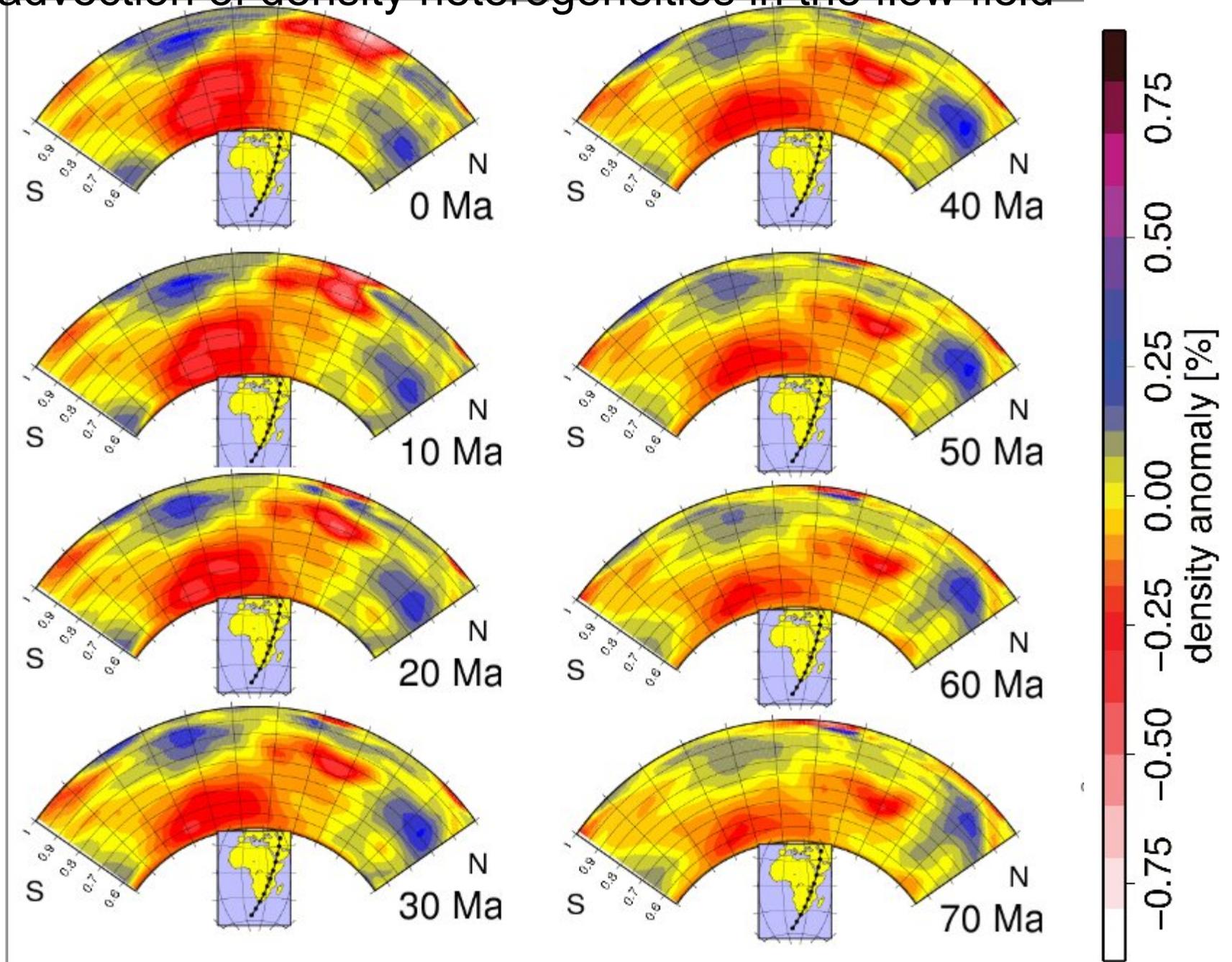
Gurnis, Mitrovica, Ritsema and van Heijst (G-Cubed, 2000)
“conclude that uplift rate, when combined with estimates of present-day dynamic topography, provides a powerful tool to constrain the properties of the deep mantle”



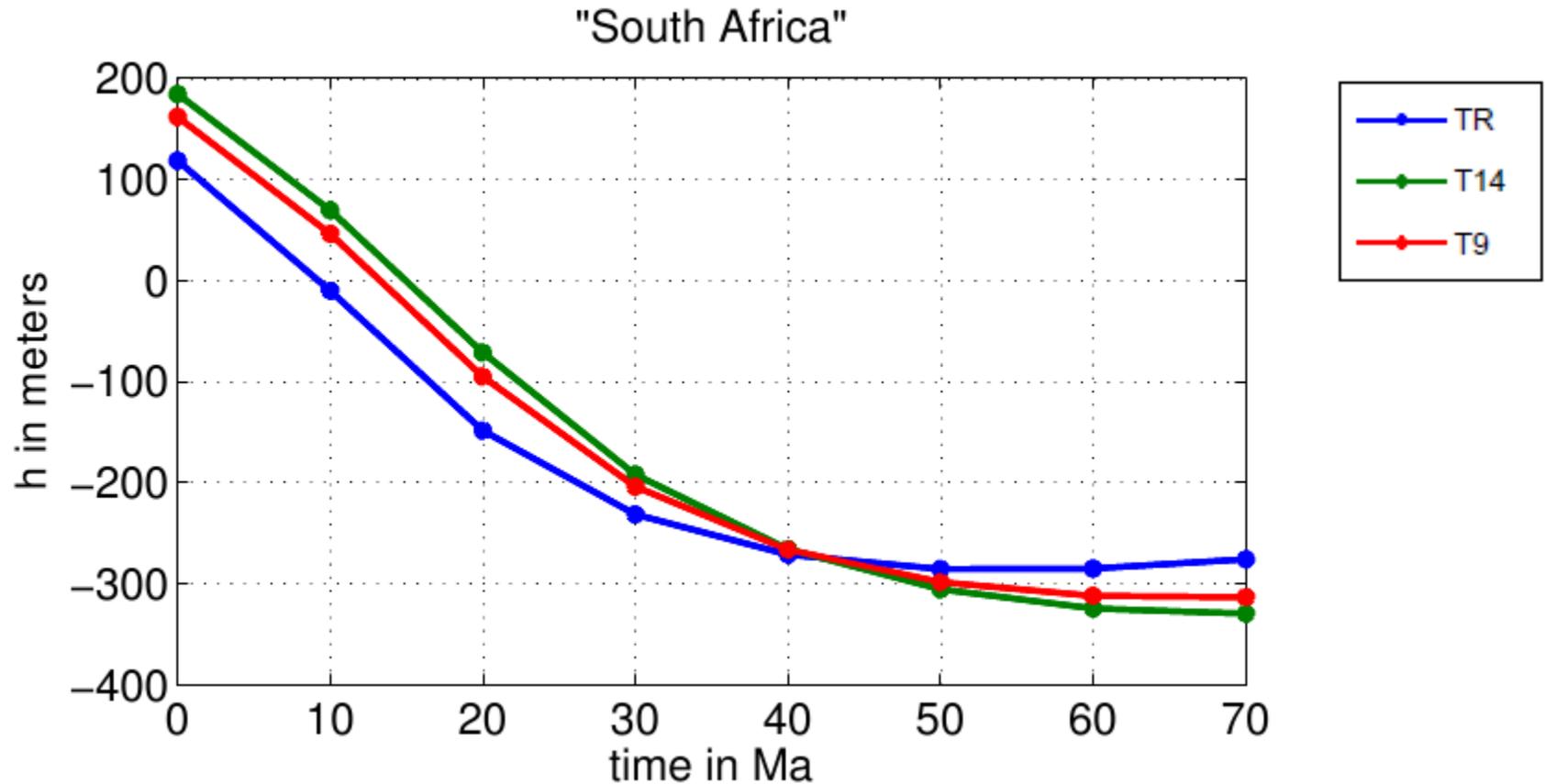
Example
Iceland/
Greenland
(Steinberger
et al., in
preparation)

How is past dynamic topography computed from mantle flow models?

Backward-advection of density heterogeneities in the flow field



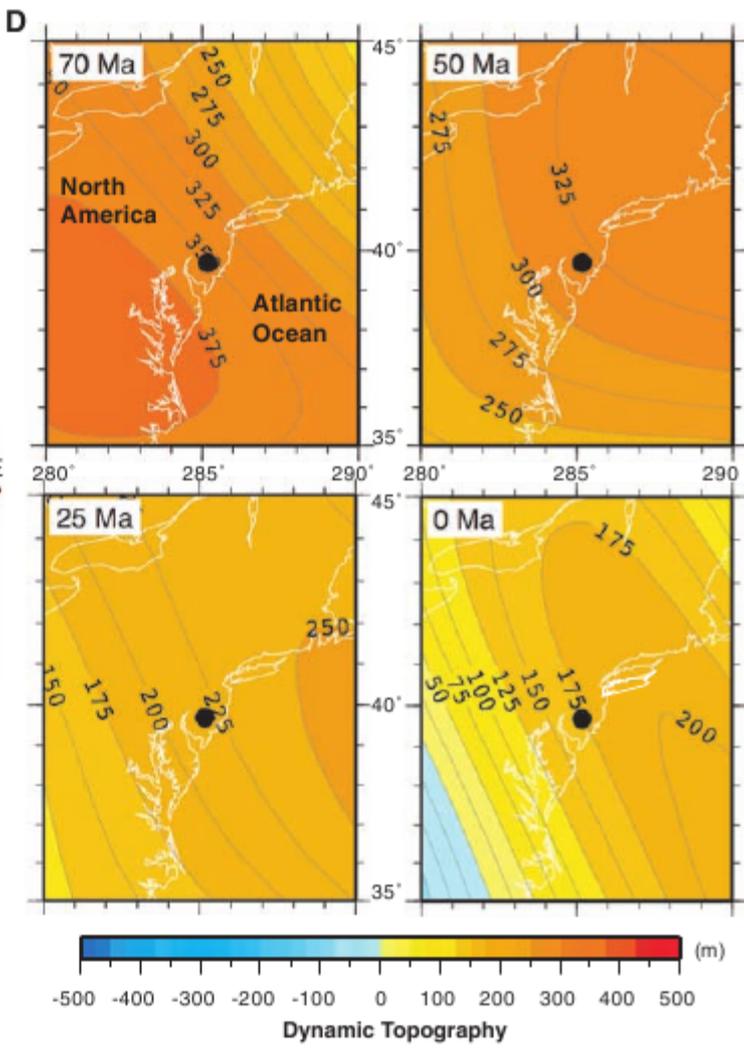
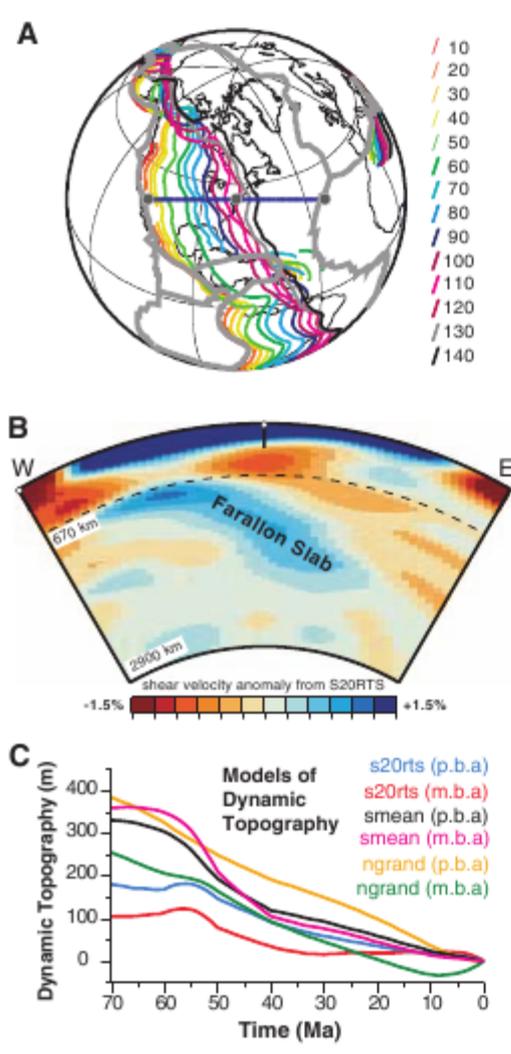
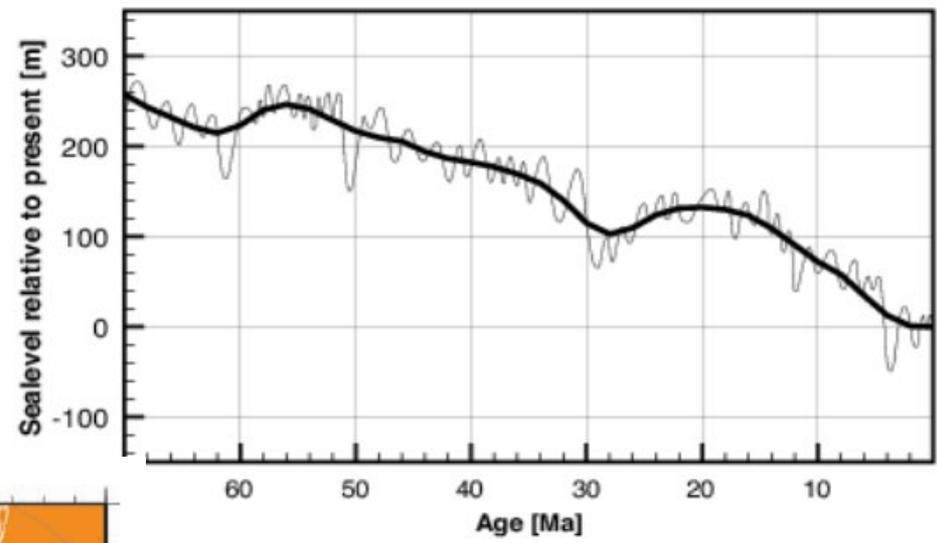
Example 1: Recent uplift of southern Africa



Dependence on lateral viscosity variations

$$\eta = \eta(r) \cdot e^{-E \cdot (T' - 0.5)}$$

Combined with plate reconstructions to compute uplift/subsidence in reference frame of moving plate



Example 2: Explaining Sea Level Curves on the East Coast of North America (Müller et al., 2008)

Use “pure backward advection” vs. “modified backward advection” in which negative density anomalies in upwellings are continued upward to 220 km, and positive density anomalies in downwellings are removed from uppermost 220 km.

Example 3: Explaining marine inundations in Australia (Heine et al., 2009)

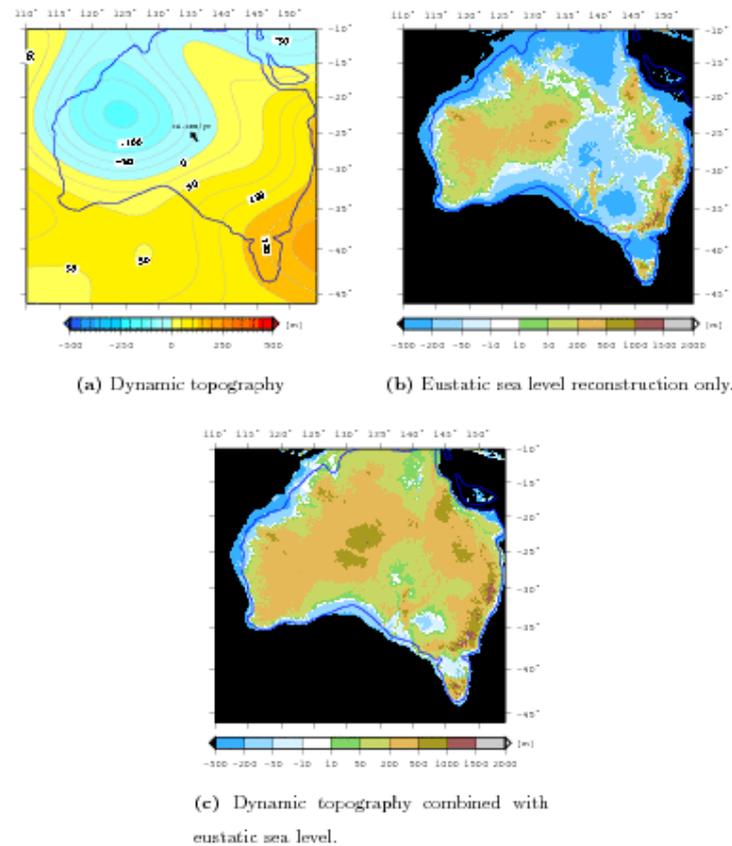


Figure 5: Topographic reconstructions for Australia at ≈ 64 Ma (Early Paleocene) using the interpreted environment for the Cenozoic 1 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at 25° S/ 135° W; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the present day.

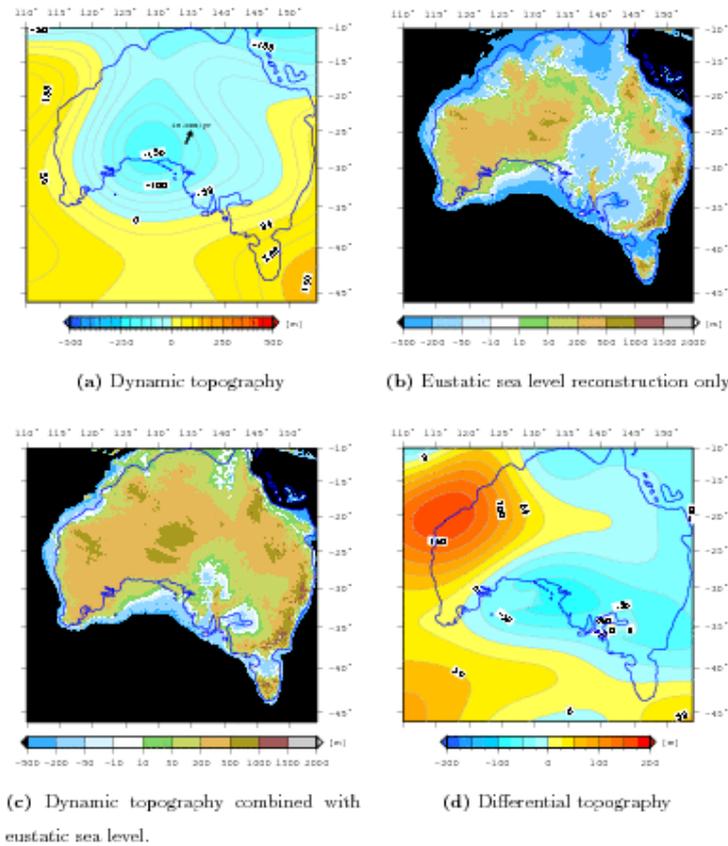


Figure 6: Topographic reconstructions for Australia at ≈ 41 Ma (Mid-Eocene) using the interpreted environment for the Cenozoic 2 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at 25° S/ 135° W; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography (scaled by a factor of 0.3) and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the previous timestep (64 Ma).

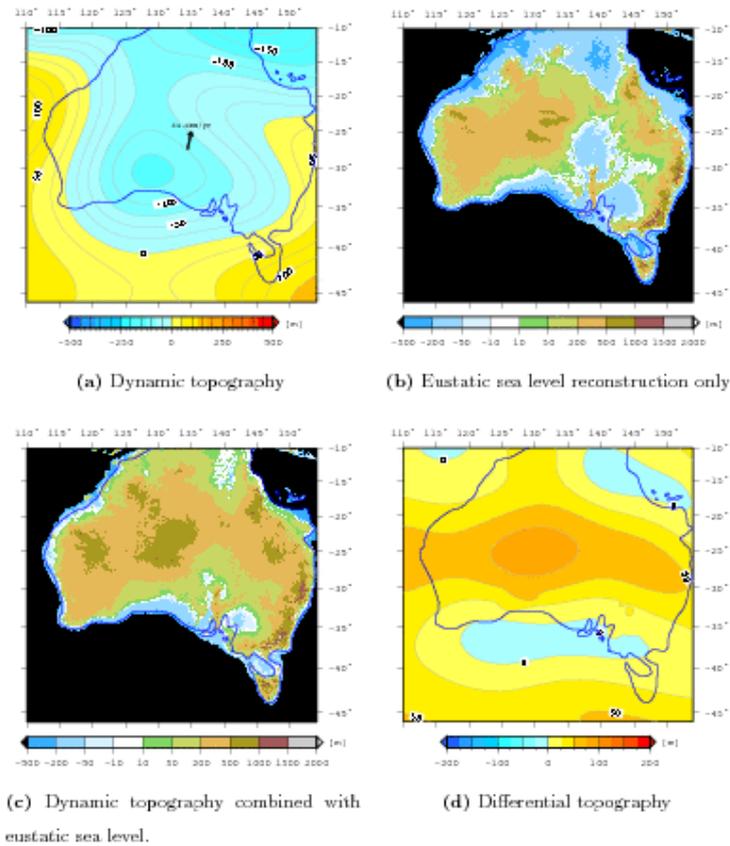


Figure 7: Topographic reconstructions for Australia at ≈ 31 Ma (Early/Mid-Oligocene) using the interpreted environment for the Cenozoic 3 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at 25° S/ 135° W; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography (scaled by a factor of 0.3) and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the previous timestep (41 Ma).

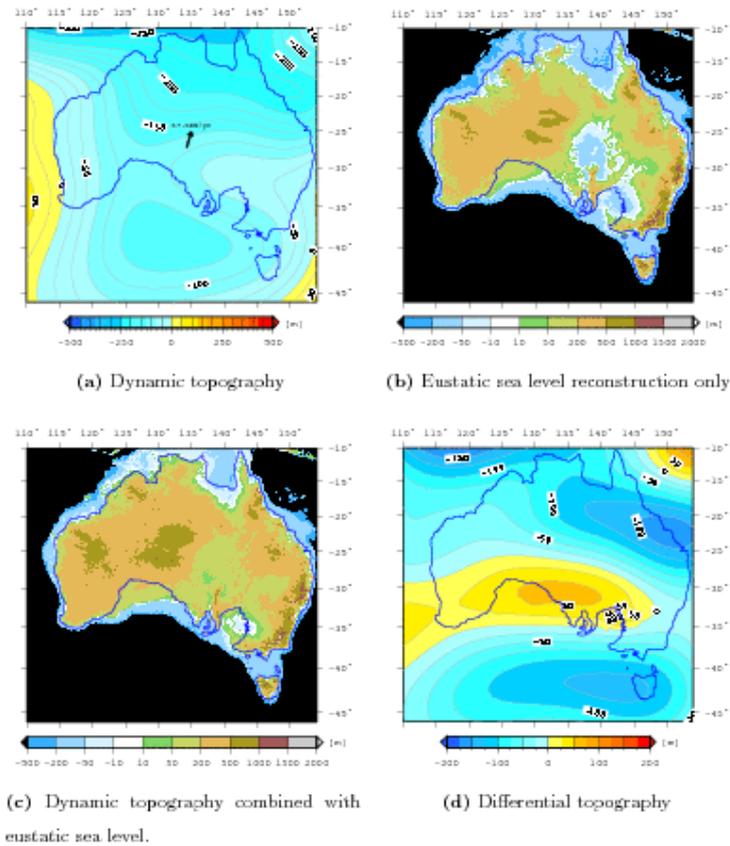


Figure 8: Topographic reconstructions for Australia at ≈ 13 Ma (Early/Mid-Miocene) using the interpreted environment for the Cenozoic 4 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at 25° S/ 135° W; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography (scaled by a factor of 0.3) and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the previous timestep (31 Ma).

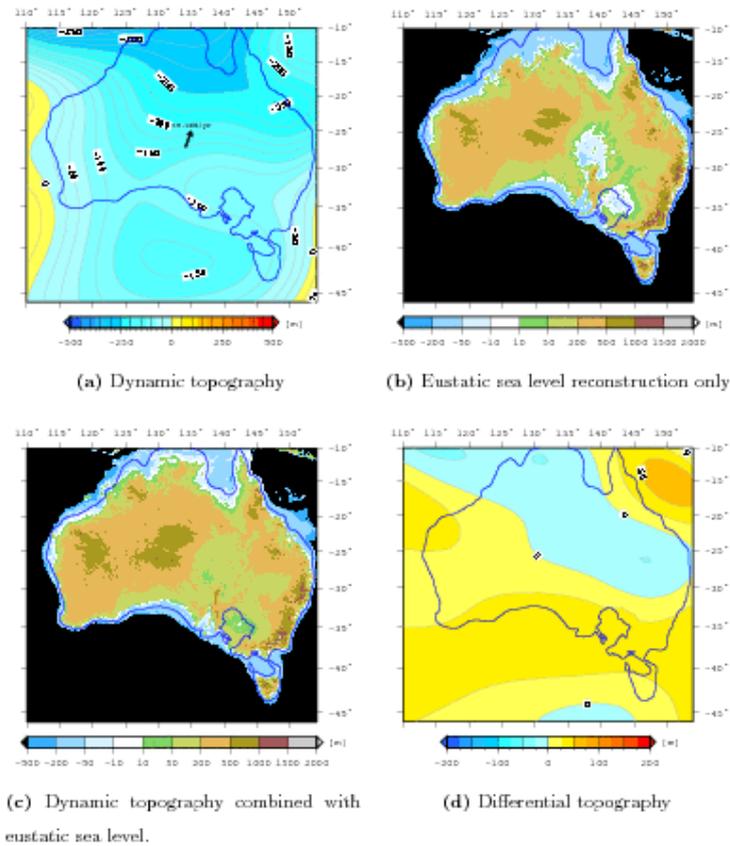


Figure 9: Topographic reconstructions for Australia at ≈ 8 Ma (Late Miocene) using the interpreted environment for the Cenozoic 5 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at 25° S/ 135° W; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography (scaled by a factor of 0.3) and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the previous timestep (13 Ma).

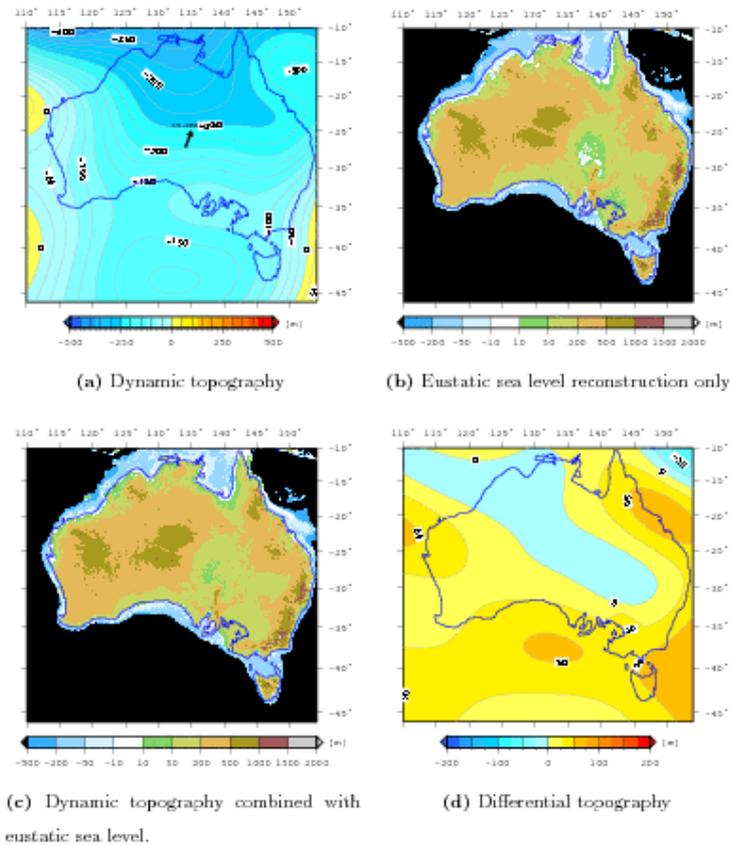


Figure 10: Topographic reconstructions for Australia at ≈ 3 Ma (Pliocene) using the interpreted environment for the Cenozoic 6 timeslice of [Langford et al. \[1995\]](#). a) Isostatically corrected dynamic topography model. Plate motion indicated by vector at $25^\circ\text{S}/135^\circ\text{W}$; b) Reconstruction using eustatic sea level only; c) Combined dynamic topography (scaled by a factor of 0.3) and eustatic sea level reconstruction with backstripped sediments in Eucla and Murray Basins; d) Differential topography relative to the previous timestep (8 Ma).

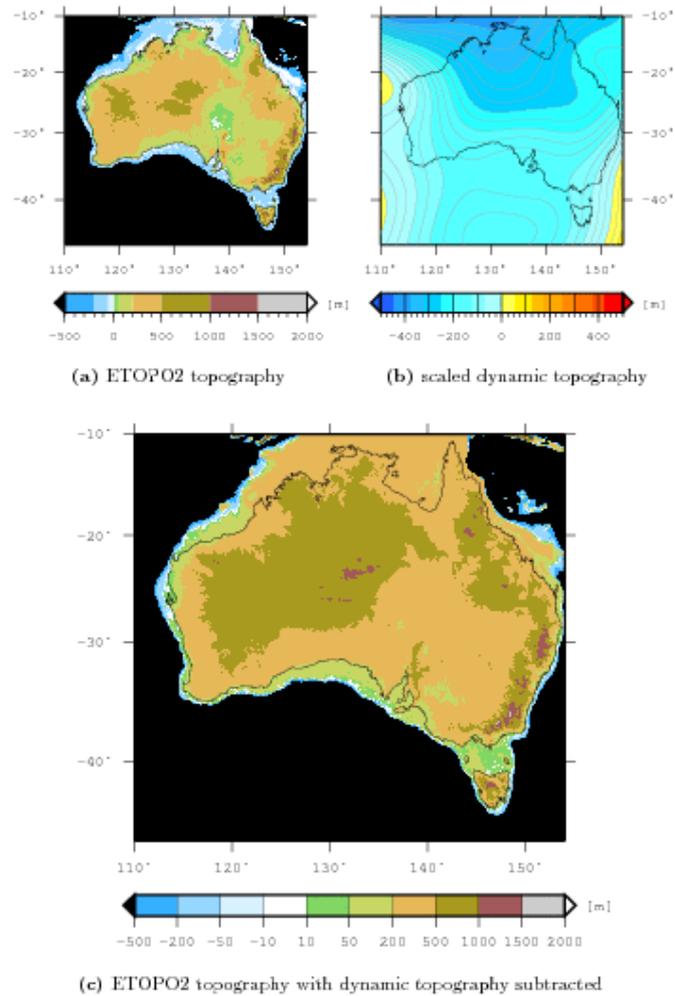
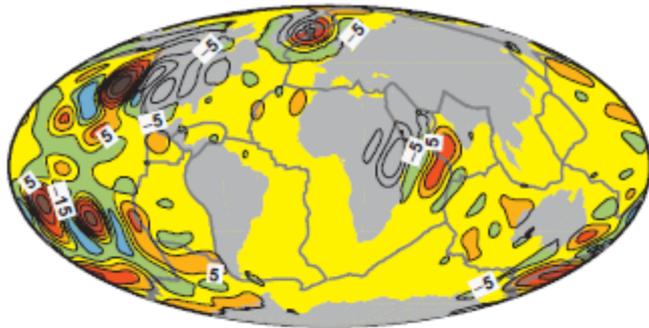
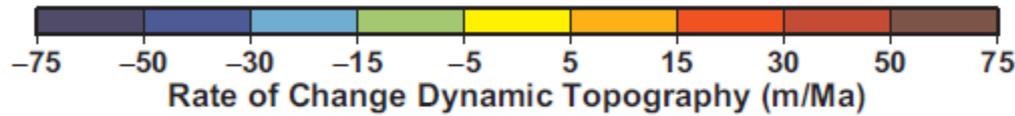
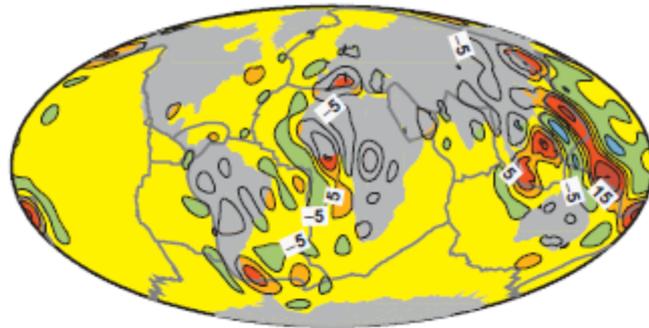


Figure 4: Present-day Australian topography, dynamic topography and difference topography. The topography with the dynamic component subtracted (Fig 4c) is used as base grid for topography reconstructions. (a) ETOPO2 present-day surface topography; (b) Scaled and isostatically corrected dynamic topography; (c) ETOPO2 topography with dynamic topography component subtracted. Thin, black line is present-day coastline, figures (a) and (c) have the same colourscale.

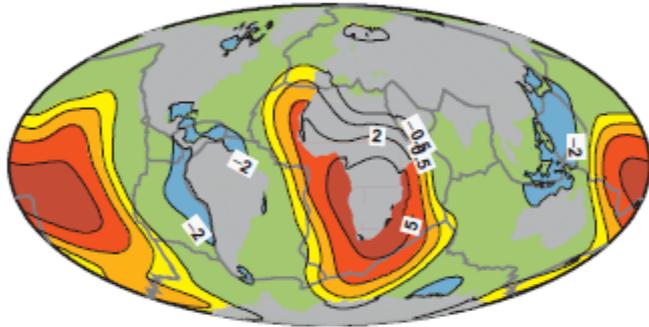
A) Upper-Mantle Upwelling



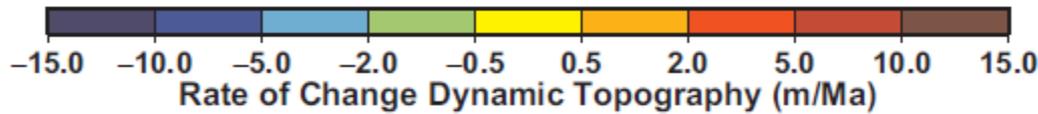
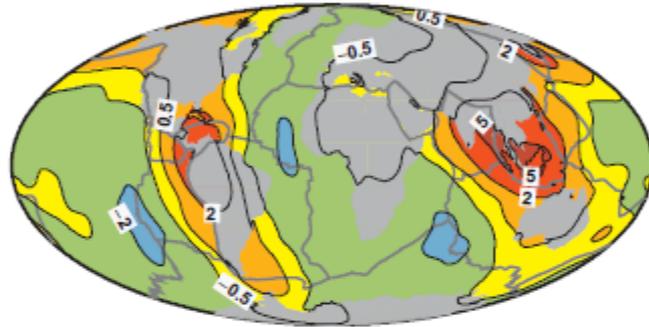
B) Upper-Mantle Downwelling



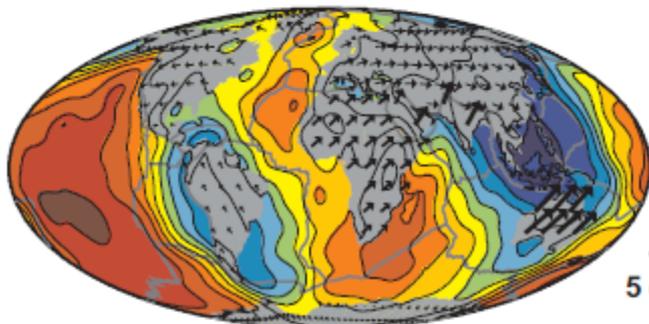
C) Lower-Mantle Upwelling



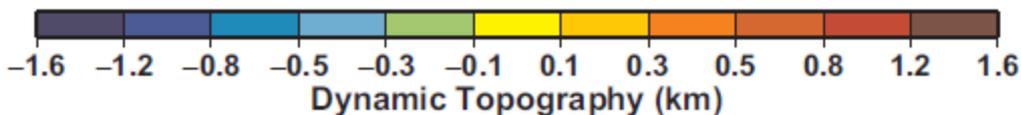
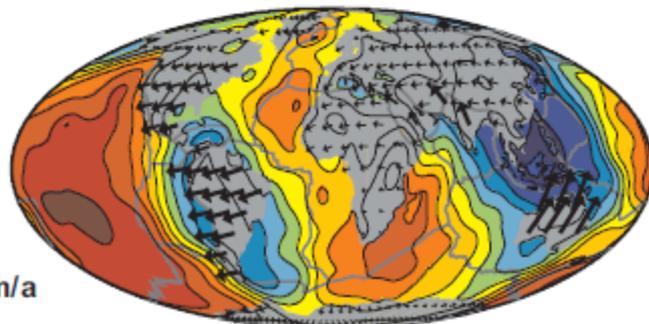
D) Lower-Mantle Downwelling



E) NNR Plate Motions



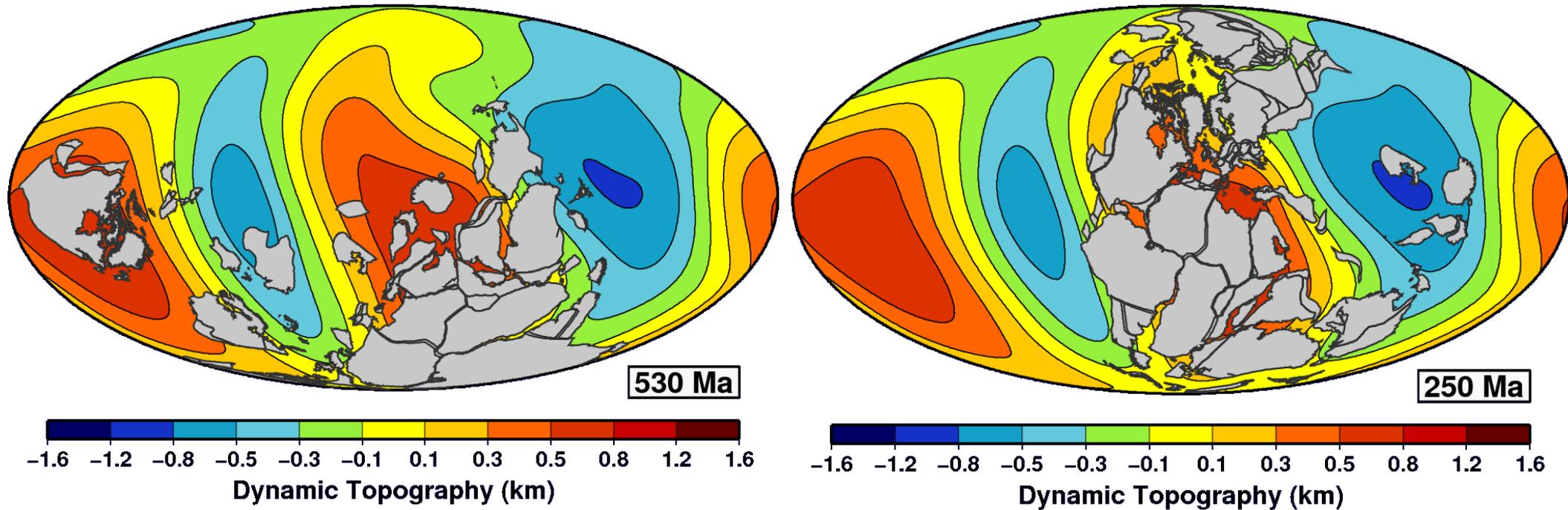
F) HS3 Plate Motions



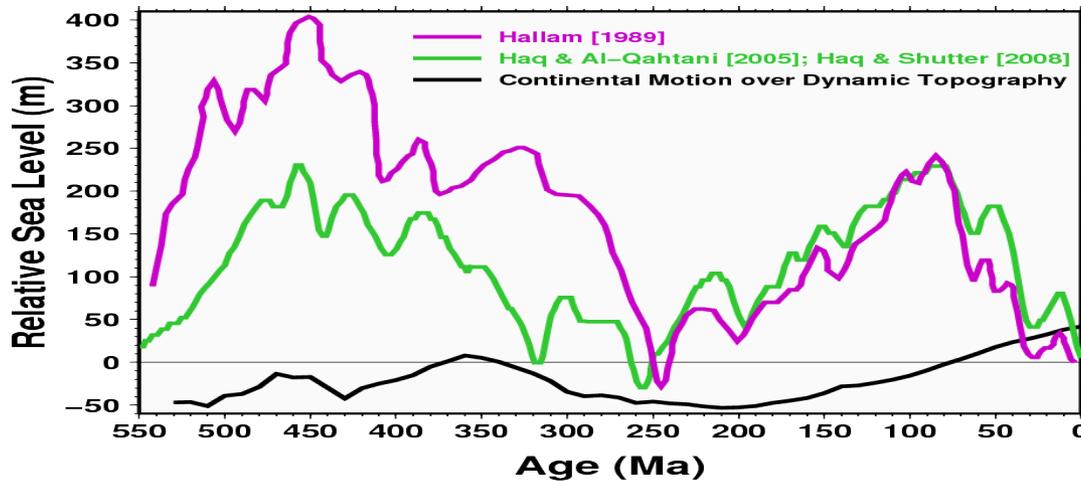
Example 4: Effect on Sea Level (Conrad and Husson, Lithosphere, 2009)

“continents preferentially conceal depressed topography associated with mantle downwelling, leading to net seafl oor uplift and $\sim 90 \pm 20$ m of positive sea-level offset. Upwelling mantle flow is currently amplifying positive dynamic topography and causing up to 1.0 m/Ma of sea-level rise, depending on mantle viscosity.”

Minimum Global Sea Level @ 250 Ma when Pangea was over Tuzo:
 Continents moving laterally toward regions of anomalously low topography will moves the average dynamic deflection of the seafloor toward more positive values (Sea Level Rise)



Old Palaeozoic
 Modell
 Not longitude calibrated or
 corrected for TPW



Hybrid TPW modell 0-250 Ma:
 (Torsvik et al. 2008; Steinberger & Torsvik 2008)

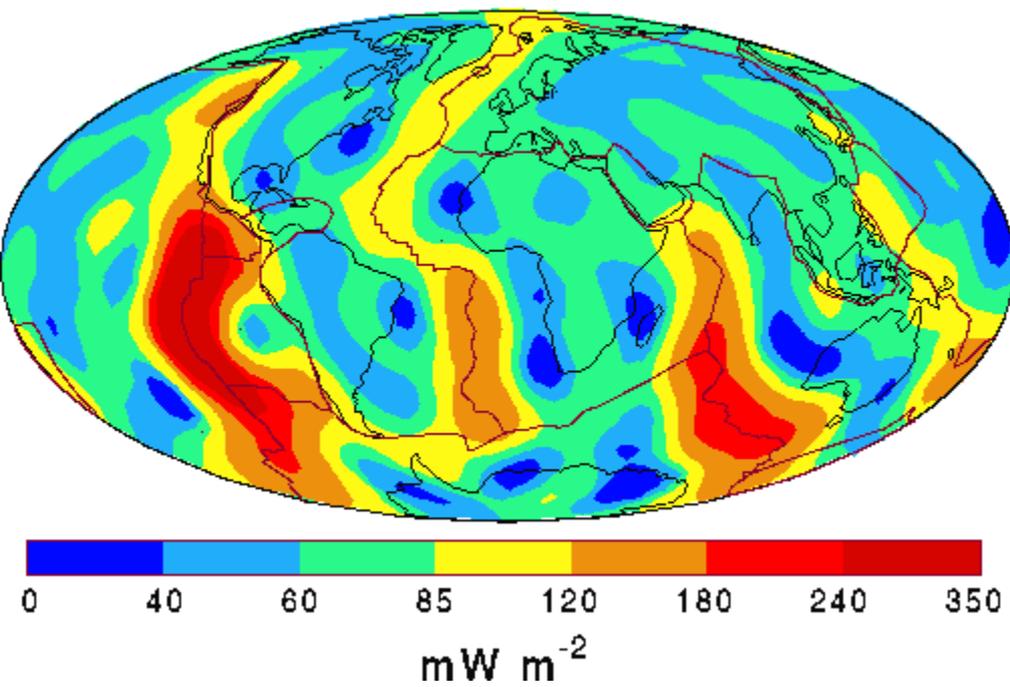
Calculations by C. Conrad (2012)

In reference frame of moving plate, e.g. high heat flow along ridges leads to subsidence as plate moves away,

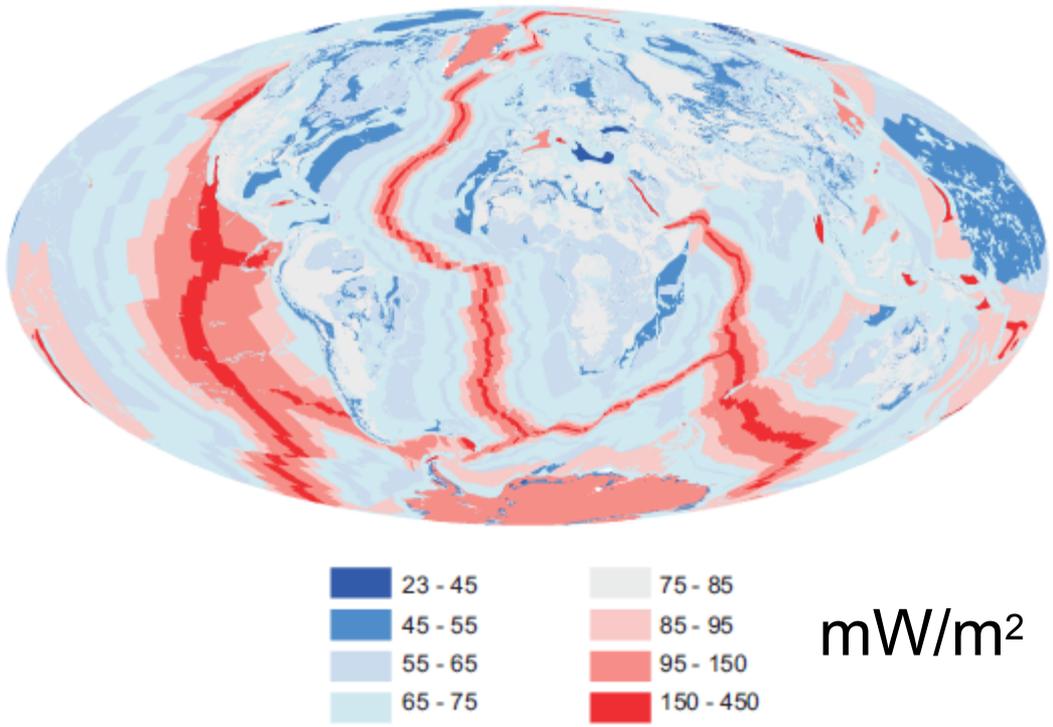


Backward advection does not consider diffusion –
 Possibly correct based on lateral heat flow variations:
 $dh/dt = \alpha / (\rho C_p) \cdot \text{Heat flow}$
 With $\alpha = 4 \cdot 10^{-5} / \text{K}$, $\rho = 3300 \text{ kg/m}^3$, $C_p = 1250 \text{ J/kg/K}$
 a heat flow difference of 100 mW/m^2 corresponds to a relative difference in uplift (subsidence) of 30 m / Myr

Global Heat Flow (Degree 12 Spherical Harmonic)
 From Pollack et al. (1993)
 Heat Flow



From Davies and Davies (2010)

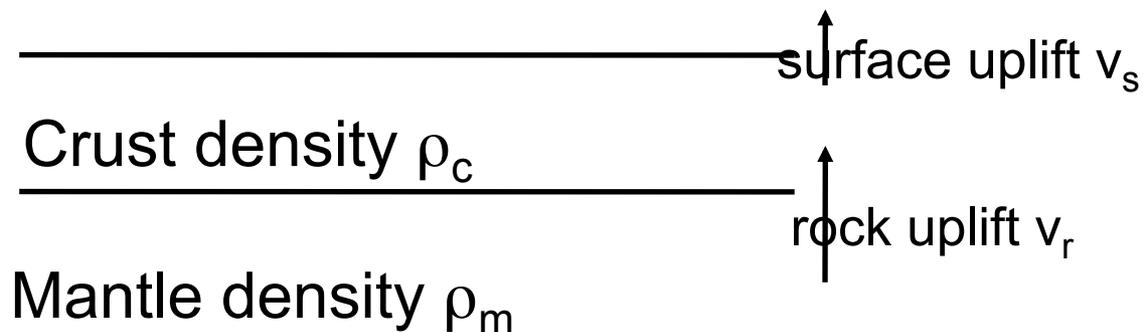


mW/m²

Comparison of computed changes of normal stress T_{rr} (possibly after correction for heat flow variations) to observations, taking erosion into account:

Distinguish rock uplift v_r and surface uplift v_s ;

$v_r - v_s = \text{erosion rate}$



$$\frac{dT_{rr}}{dt} = g (\rho_m \cdot v_r - \rho_c \cdot (v_r - v_s))$$

$$\frac{\partial T_{rr}}{\partial t} + v_x \frac{\partial T_{rr}}{\partial x} + v_y \frac{\partial T_{rr}}{\partial y}$$

in reference frame of moving plate

Forward models based on subduction history

Example 1: Gurnis, Nature, 364, 1993

ARTICLES

Phanerozoic marine inundation of continents driven by dynamic topography above subducting slabs

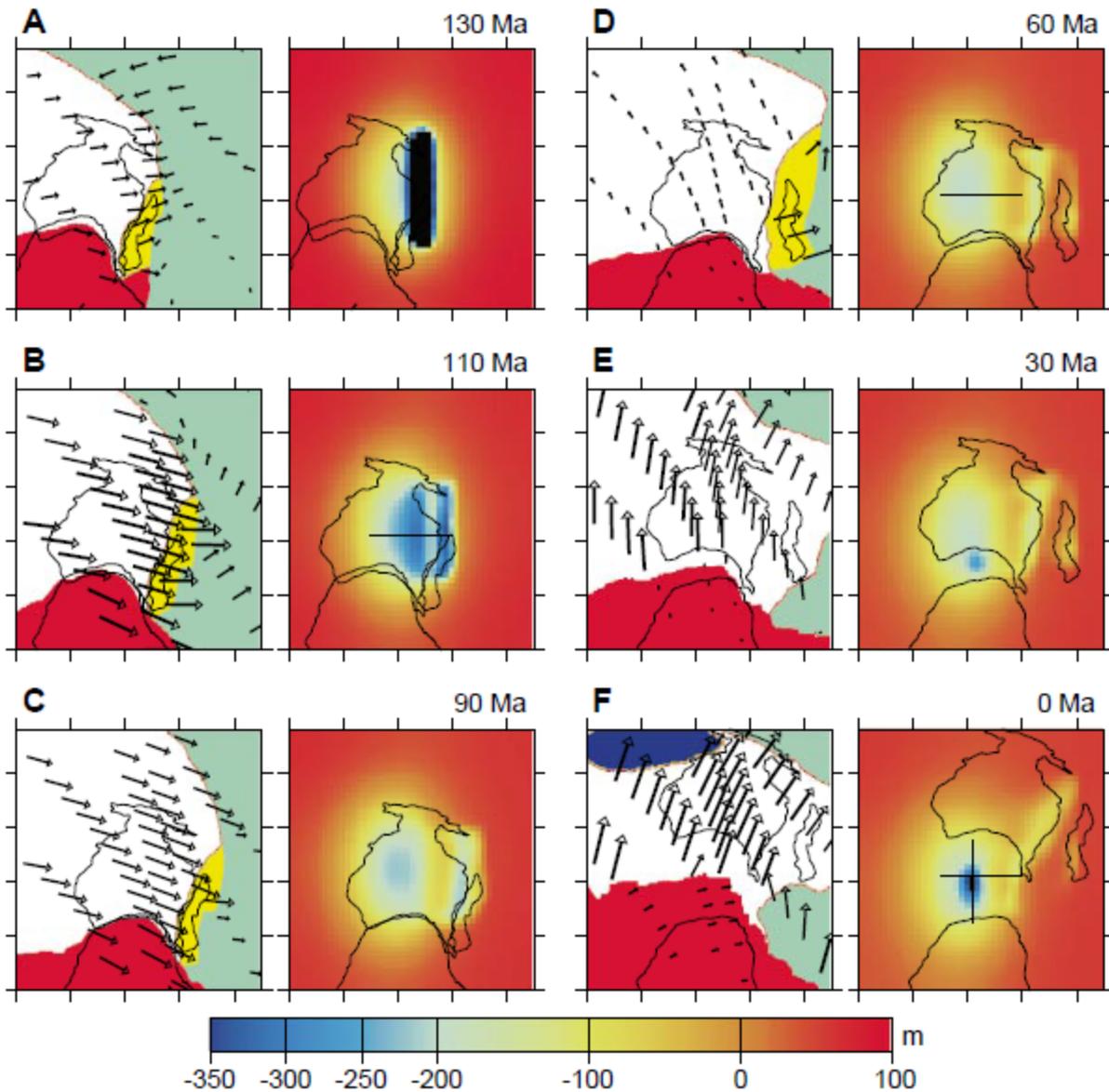
Michael Gurnis

Department of Geological Sciences, University of Michigan, Ann Arbor, Michigan 48109-1063, USA

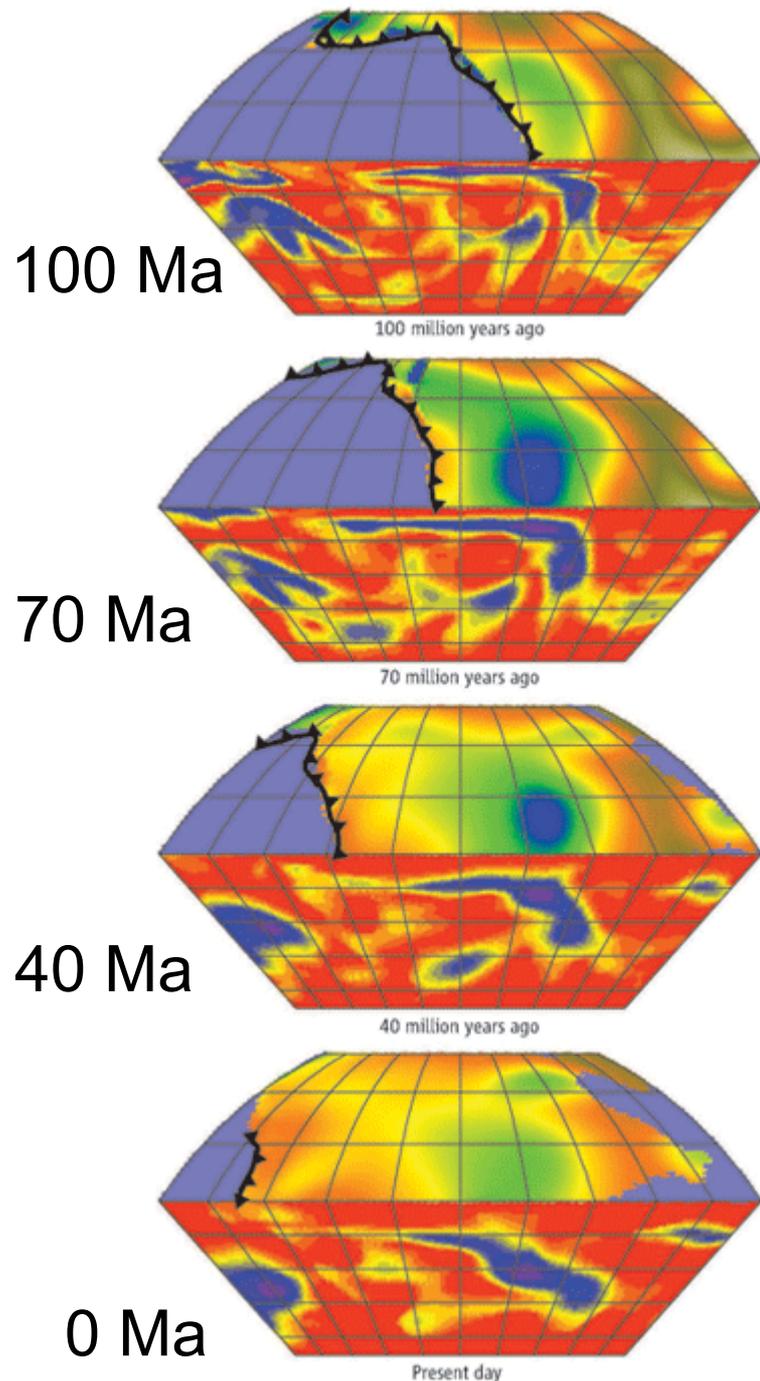
A spherical model of mantle flow constrained by the locations of trenches can be used to predict the dynamic topography of the Earth's surface, and hence the marine inundation of continents. For past periods of high sea level, the predicted geographical pattern of flooding correlates well with the geological record. The high spatial correlation may result from increased plate velocities at these times, leading to increased rates of subduction, subsidence and inundation at convergent margins.

Forward models based on subduction history

Example 2: Gurnis, Müller and Moresi, Science, 279, 1998



“The dynamic models infer that a subducted slab associated with the long-lived Gondwanaland-Pacific converging margin passed beneath Australia during the Cretaceous, partially stagnated in the mantle transition zone, and is presently being drawn up by the Southeast Indian Ridge.”



Adjoint models:
Finding the initial model
that matches present-day
structure (inferred from
tomography) with surface
plate motion boundary
conditions through time.
Example: Model of Liu,
Spasojević and Gurnis
(Science, 2008)

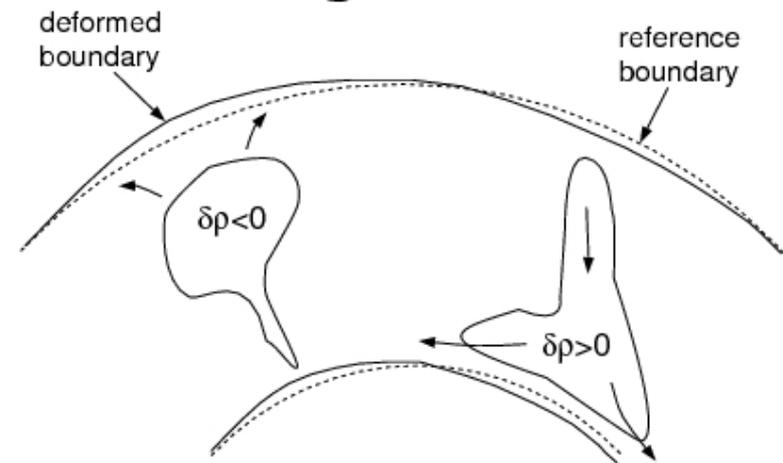
TABLE 2. STRENGTHS AND LIMITATIONS OF DYNAMIC TOPOGRAPHY MODELS

Strengths	Limitations
<i>Forward models</i>	
Computationally relatively cheap	Dependent on synthetic initial condition
Can achieve large dimensional time	Kinematically driven
Can achieve high resolution	May result in unrealistic slab advection
Can be compared to mantle tomography	Upwelling is usually passive
<i>Backward advection</i>	
Computationally relatively cheap	Thermal diffusivity is neglected
Can achieve high resolution	Limited to a few tens of millions of years
Consistent with the present-day density structure of the mantle	Thermal boundary layers require special treatment
	Usually kinematically driven
<i>Adjoint models</i>	
Consistent with the present-day density structure of the mantle	Computationally expensive
Thermal diffusivity is accounted for	Usually kinematically driven

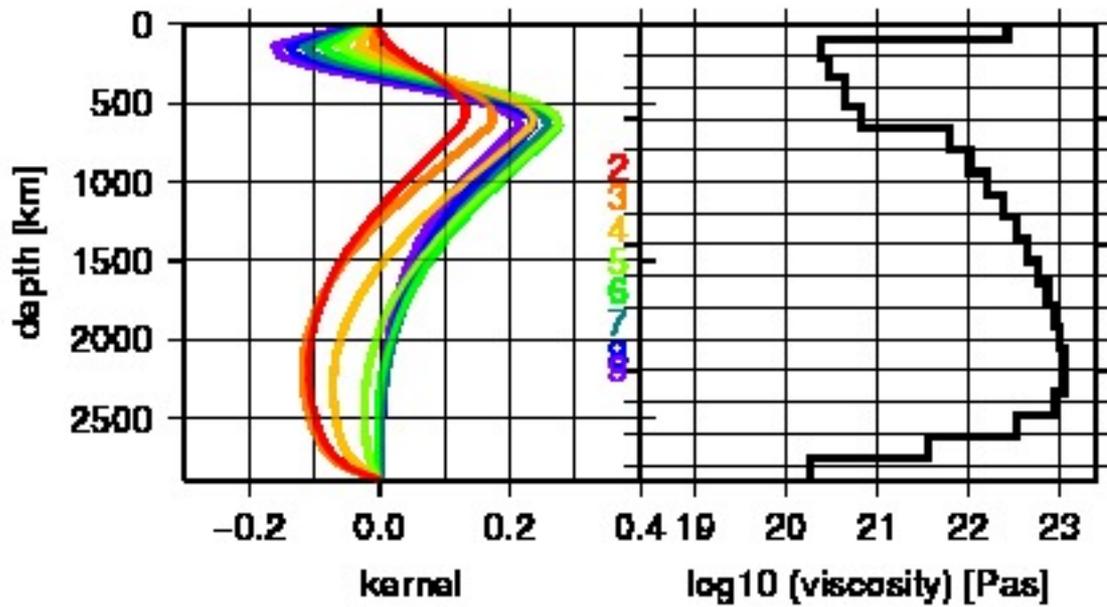
From Flament, Gurnis and Mueller, Invited Review in *Lithosphere*, **5**, 2013.

Models of topography and equipotential surface on other planets

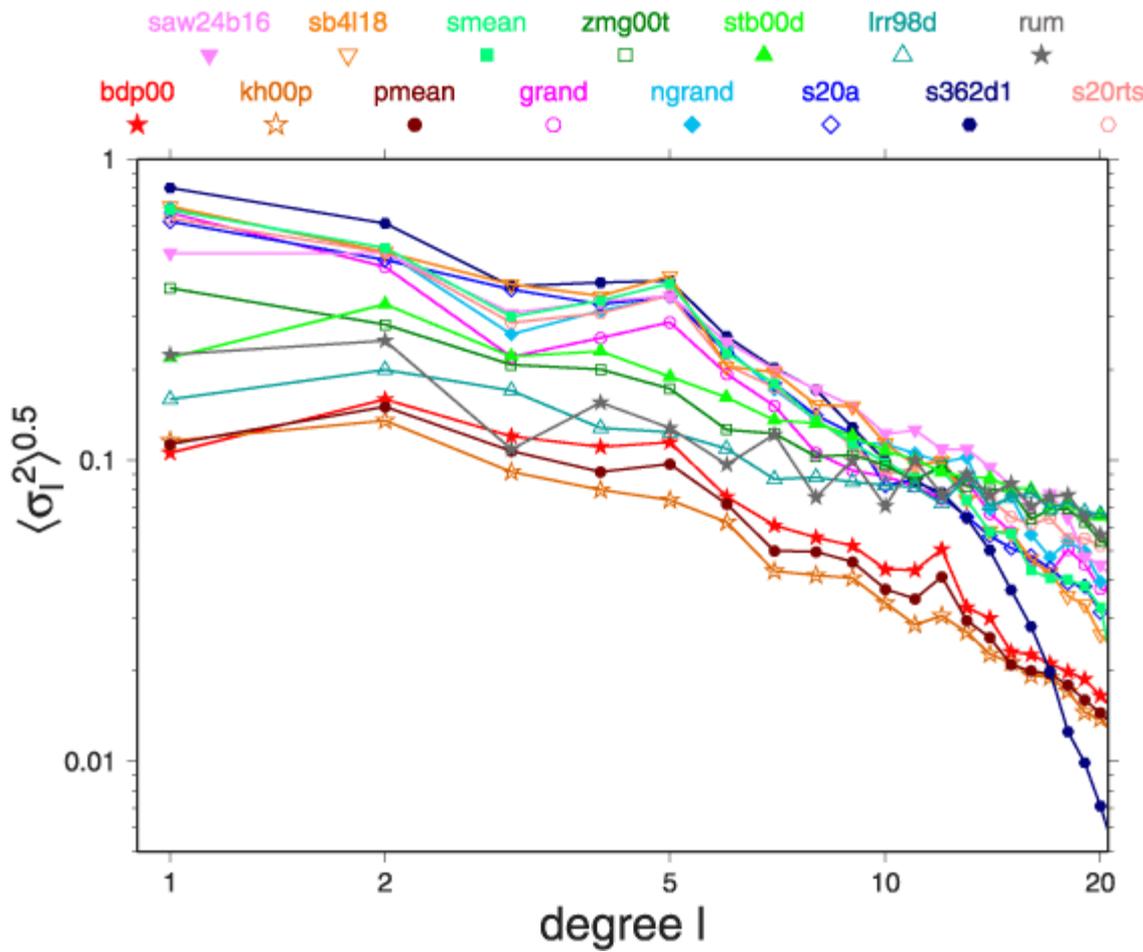
Computation of mantle flow field, boundary deformations and geoid



- Density anomalies (inferred from tomography models) drive flow, computed with spectral method (Hager and O'Connell, 1981)
- Flow deforms boundaries
- Density anomalies and deformed boundaries contribute to geoid anomalies



Geoid kernels $K_l(r)$ describe effect of density anomalies at spherical harmonic degree l and radius r on geoid



Compute expected gravity power vs. l

Assume radial correlation $\sim 1/l$

Power spectrum of tomography models (depth averaged)

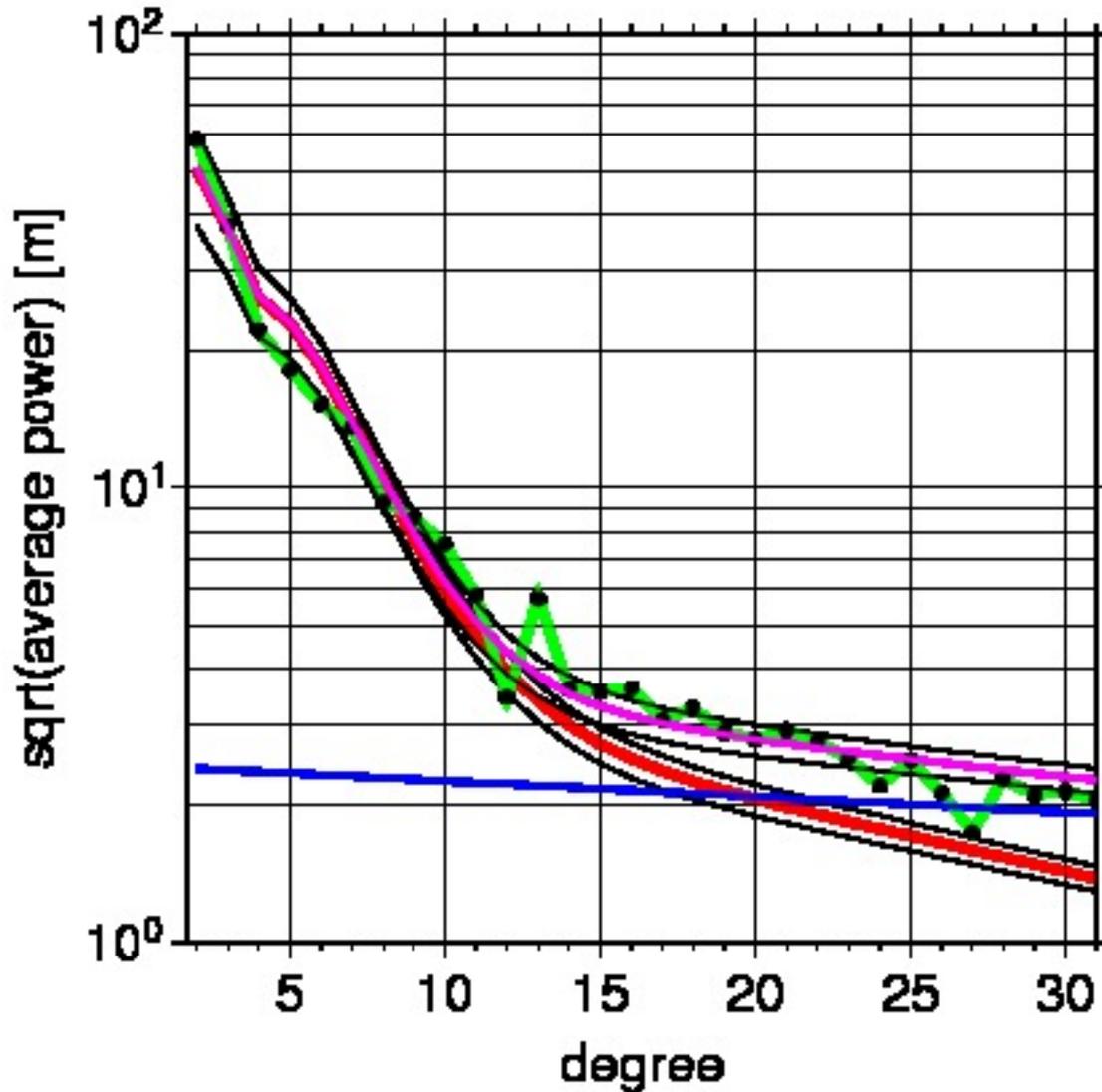
$$\sigma_{1,l} \cong \sigma_0 / \sqrt{(l+1)(2l+1)}$$

(Becker and Boschi, 2002)



Earth **observed** – modelled
mantle / **lithosphere** / **total**
contribution

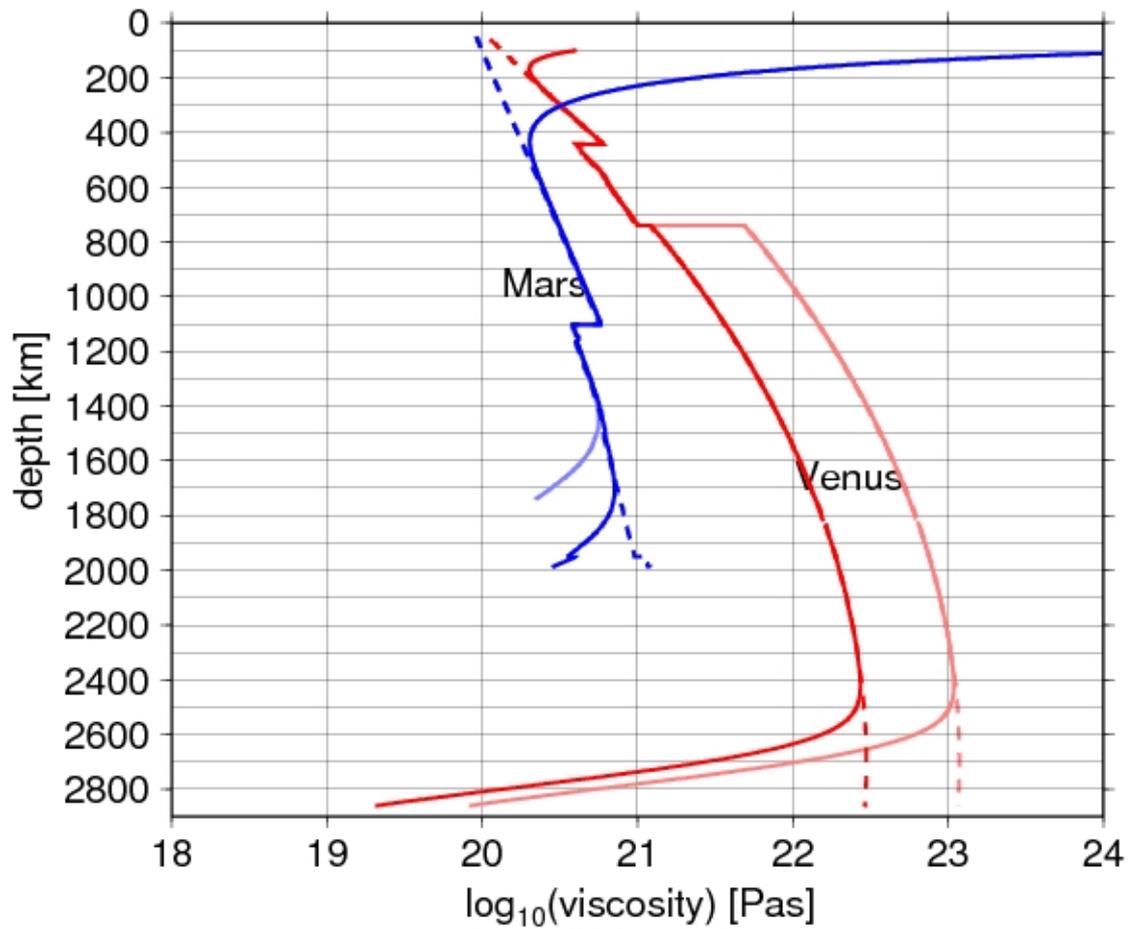
(Steinberger and Holme, GRL,
2002)

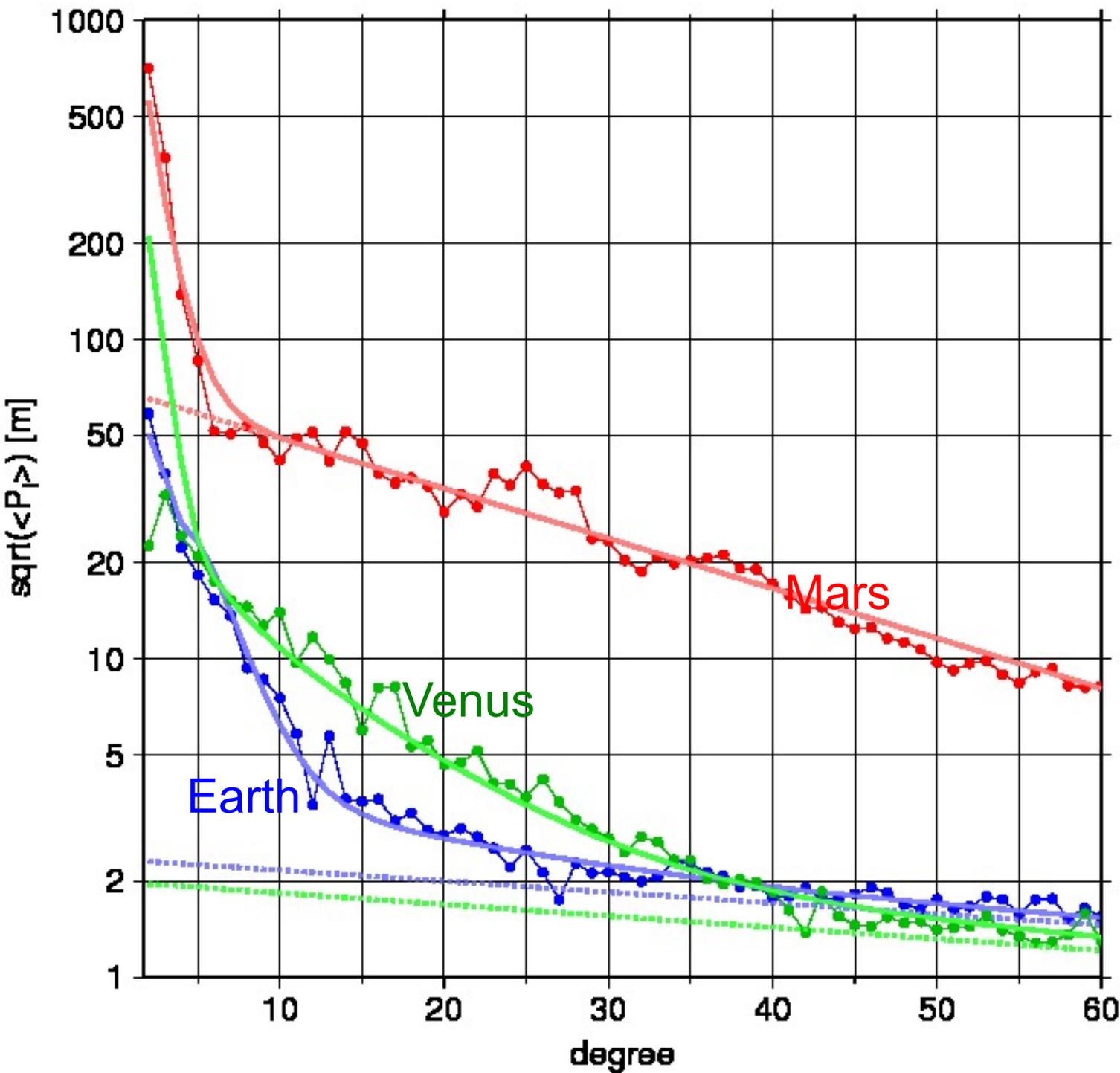


→ Sublithospheric mantle
contribution important
up to ~ degree 25-30

→ Lithosphere contribution
with “white” power spectrum;
observation-based magnitude

Pressure and temperature, and hence viscosity
Increase less strongly with depth in other planets



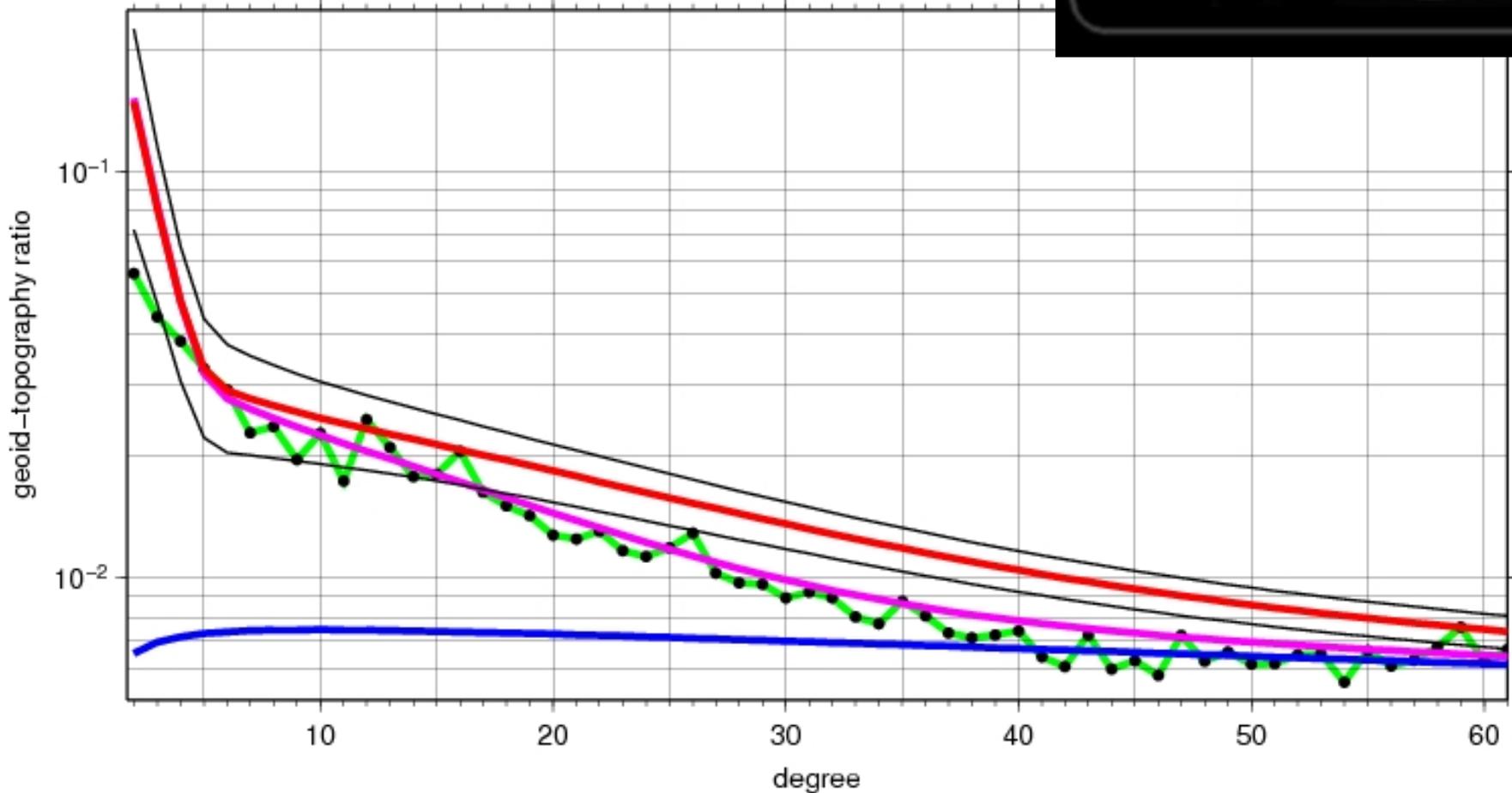


→ With suitable modifications (viscosity profile from temperature and pressure vs. depth; elastic lithosphere) match spectra for Venus and Mars (Steinberger, Werner and Torsvik, Icarus, 2010)

→ For mantle-dominated part can infer depth averaged mantle density (and compare with distribution of volcanics)

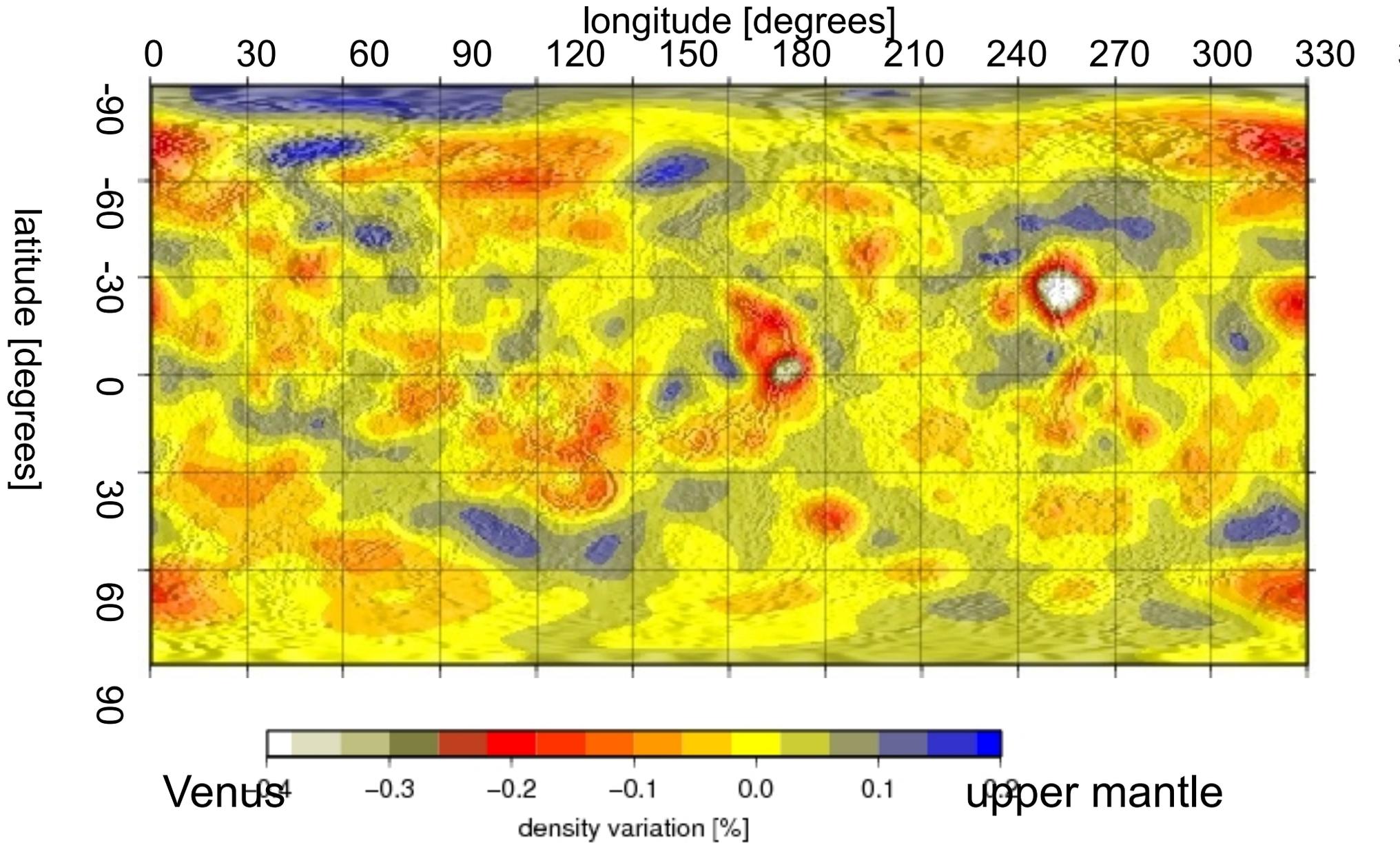
Venus **observed** – modelled
mantle / **lithosphere** / **total**
contribution

In contrast to Earth and Mars our model
also fits the topography spectrum for
Venus for degrees $>\sim 5$ –
Venus topography appears to be
largely dynamically supported



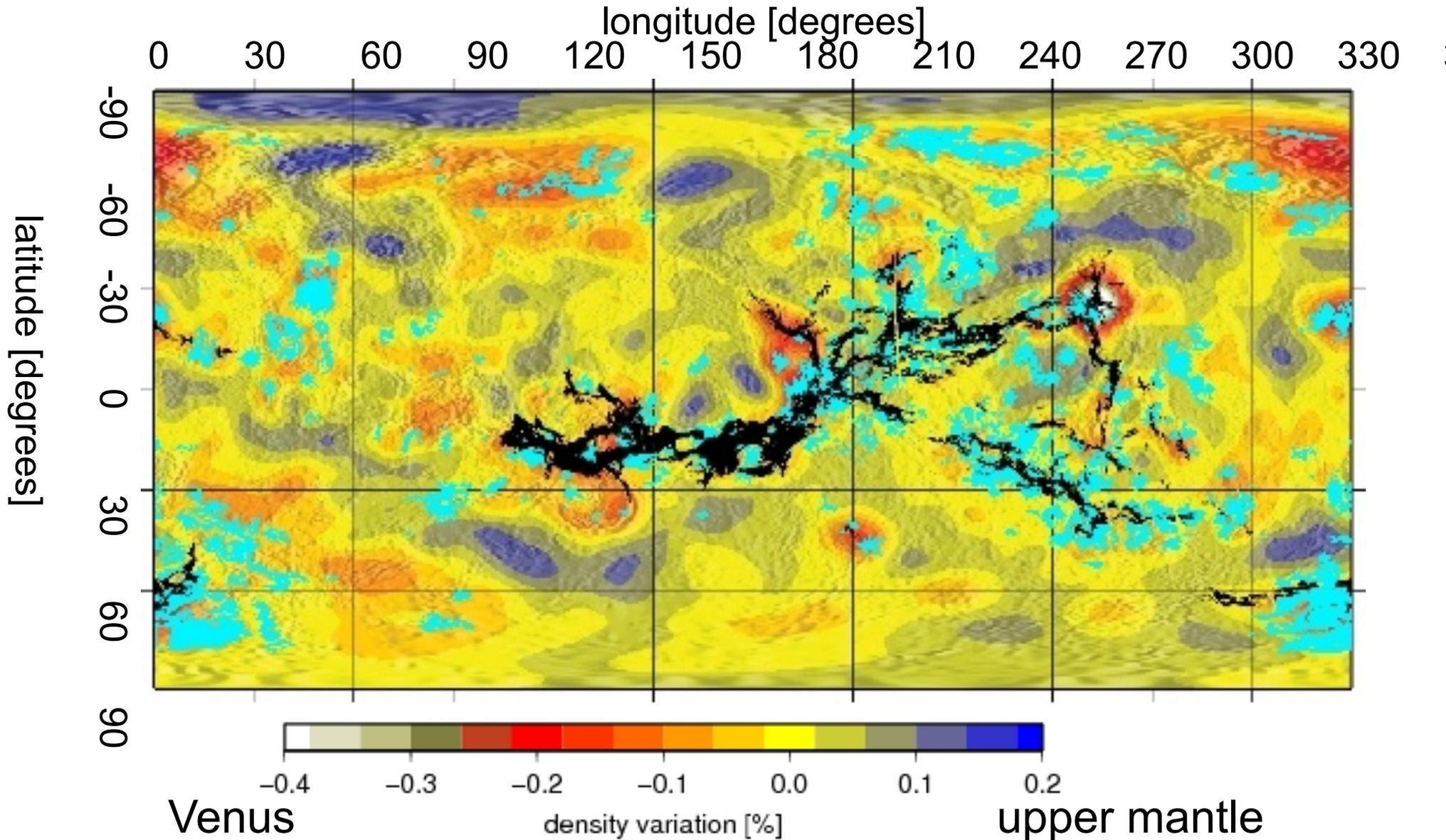
Venus

inferred upper mantle density variation



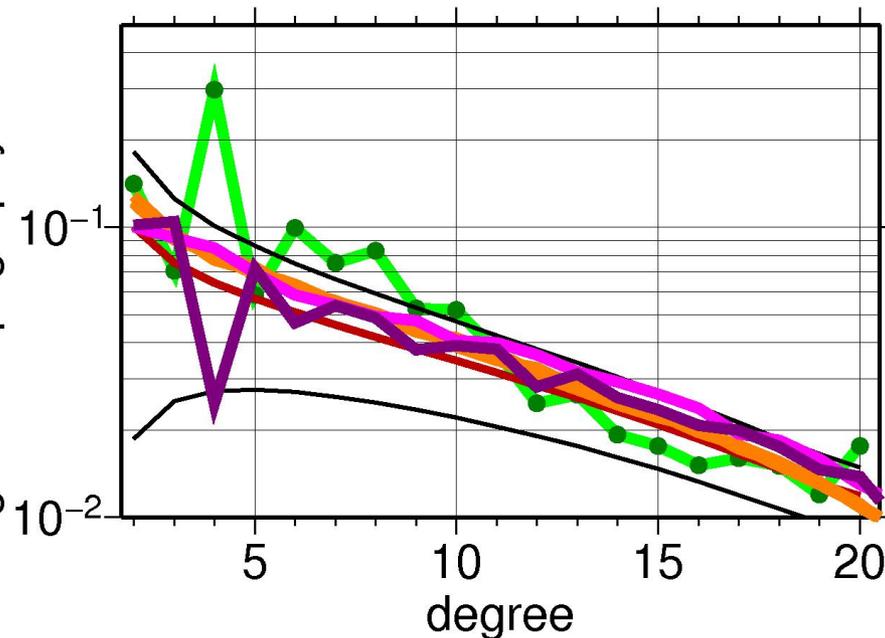
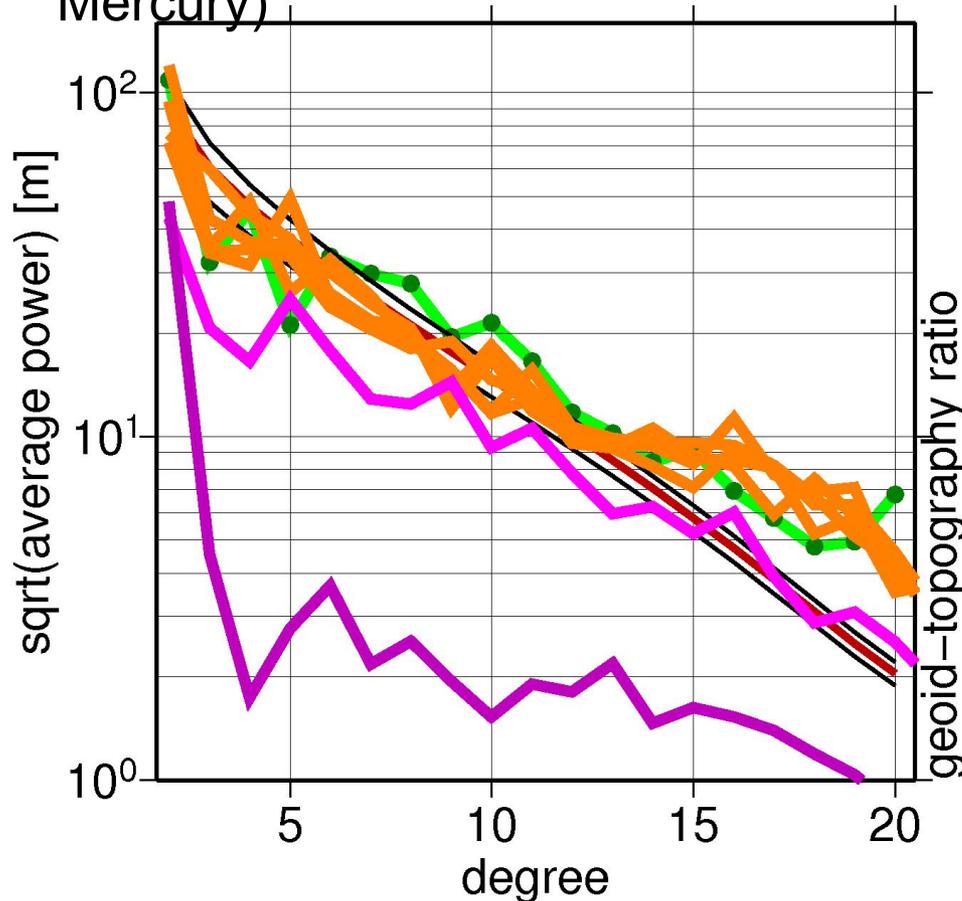
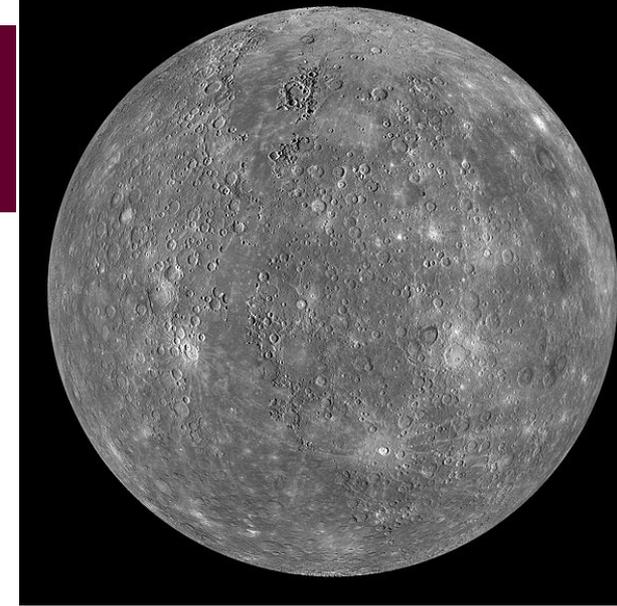
Venus

inferred upper mantle density variation
distribution of rift zones (in black)
and lobate plains (Ivanov, 2008)



But:

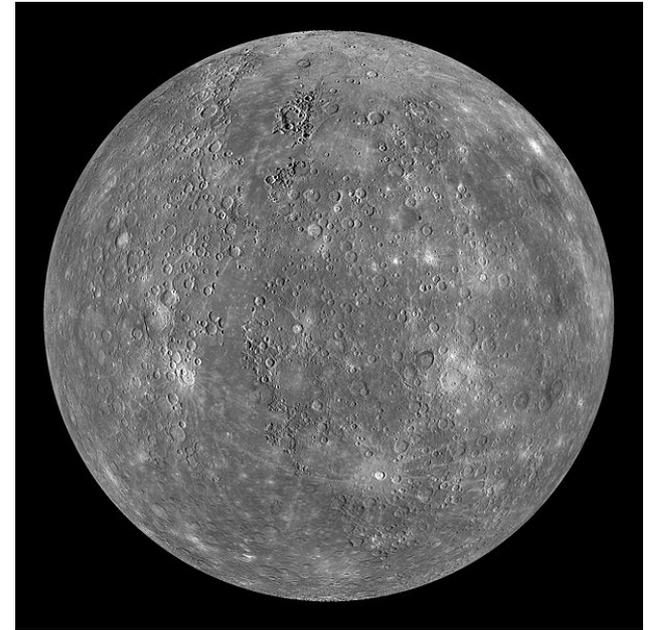
Assumption that density spectrum is same as on Earth not necessarily correct. Therefore also evaluate results from forward convection model (here: Mercury)



Red lines: Assuming same density spectrum as Earth

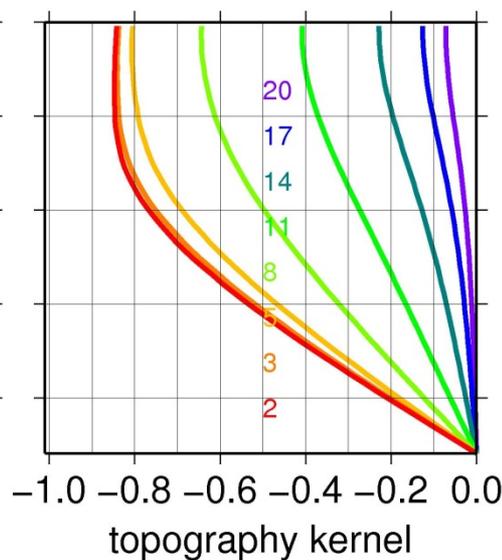
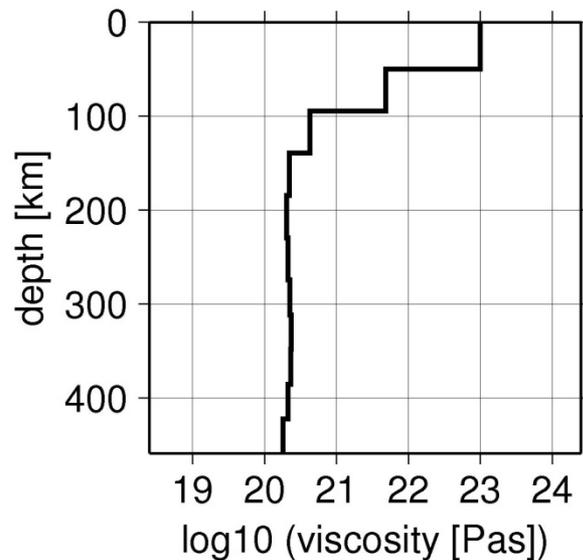
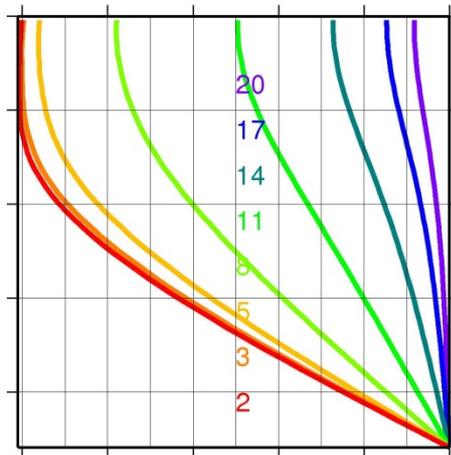
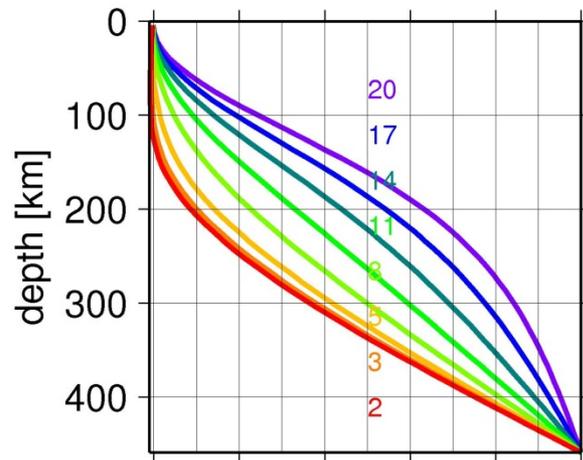
Green lines: Observed

Orange / light purple / dark purple: Forward convection models, degrees 31, 63, 127



No-slip;
viscous rheology

No-slip;
elastic lithosphere
174 km thickness;
membrane stresses
not considered



No-slip; elastic lithosphere 174 km thickness;
membrane stresses considered
(Turcotte et al., JGR 86, 3951-59, 1981)

$$p = g[\rho_c h - \rho_m h_g - (\rho_m - \rho_c)w] \quad (3)$$

In writing the term $-(\rho_m - \rho_c)w$ it is implicitly assumed that crust with density ρ_c fills the region between 0 and w . The

(modified; we do not consider crustal fill)



Thank you for your attention