

Supplemental Material for: Fluctuation profiles in inhomogeneous fluids

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Appendix A: Ideal Gas

To illustrate the fluctuation profiles, we consider the ideal gas ($u(\mathbf{r}^N) \equiv 0$) in the presence of an external potential $V_{\text{ext}}(\mathbf{r})$. The resulting density profile follows a generalized barometric law [1],

$$\rho_{\text{id}}(\mathbf{r}) = \Lambda^{-D} \exp(-\beta(V_{\text{ext}}(\mathbf{r}) - \mu)), \quad (\text{A1})$$

where $\Lambda(T) = (2\pi\beta\hbar^2/m)^{1/2}$ indicates the thermal de Broglie wavelength; here $\hbar$ denotes the reduced Planck constant and m is the particle mass. One readily obtains the fluctuation profiles from (1)–(3) [in the main text] as

$$\chi_{\mu}^{\text{id}}(\mathbf{r}) = \beta\rho(\mathbf{r}), \quad (\text{A2})$$

$$\chi_T^{\text{id}}(\mathbf{r}) = \frac{\rho(\mathbf{r})}{T} (\beta V_{\text{ext}}(\mathbf{r}) - \beta\mu + D/2), \quad (\text{A3})$$

$$\chi_{\star}^{\text{id}}(\mathbf{r}) = \rho(\mathbf{r}) - (\beta V_{\text{ext}}(\mathbf{r}) + D/2)\rho(\mathbf{r}). \quad (\text{A4})$$

In the context of the OZ relations, (A2) and (A3) are relevant beyond the ideal gas, when $\rho(\mathbf{r})$ is taken to be general. For the ideal case,

$$\rho_{\text{id}}(\mathbf{r}) = \Lambda^{-D} \exp\left(\frac{D}{2} - \frac{\chi_T(\mathbf{r})}{k_B\chi_{\mu}(\mathbf{r})}\right), \quad (\text{A5})$$

$$\chi_{\star}^{\text{id}}(\mathbf{r}) = \frac{1 - D/2}{\Lambda^D} \exp\left(\frac{D}{2} - \frac{\chi_T(\mathbf{r})}{k_B\chi_{\mu}(\mathbf{r})}\right) - V_{\text{ext}}\chi_{\mu}(\mathbf{r}), \quad (\text{A6})$$

where (A6) follows from (A4) upon using (A2), (A3) and (A5) for $D \neq 2$. The Legendre transform (3) is hence complete, including the replacement of the original variables T, μ by the new variables χ_T, χ_{μ} . For completeness, the temperature of the inhomogeneous ideal gas can be expressed for $D \neq 2$ via the local susceptibilities as

$$T = (k_B\chi_{\mu})^{2/(D-2)} \left(\frac{mk_B}{2\pi\hbar^2}\right)^{D/(2-D)} \exp\left(\frac{2}{2-D}\left(\frac{D}{2} - \frac{\chi_T}{k_B\chi_{\mu}}\right)\right), \quad (\text{A7})$$

where position arguments of $\chi_{\mu}(\mathbf{r})$ and $\chi_T(\mathbf{r})$ are omitted for clarity. The thermal wavelength satisfies

$$\Lambda^{D-2} = \frac{m}{2\pi\hbar^2\chi_{\mu}} \exp\left(\frac{D}{2} - \frac{\chi_T}{k_B\chi_{\mu}}\right). \quad (\text{A8})$$

Throughout, we have retained the temperature dependence of $\Lambda(T)$. Corresponding results for the fluctuation profiles within the frequently used convention of measuring lengths against a molecular size σ as the fundamental scale and fixing $\Lambda = \sigma$ from the outset are obtained by formally setting $D = 0$ in (A3)–(A6).